

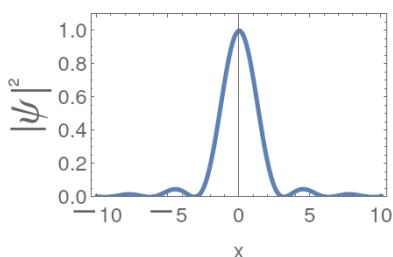
## Physics 309 Final

I pledge that I have given nor received unauthorized assistance during the completion of this work.

Name \_\_\_\_\_ Signature \_\_\_\_\_

Questions (4 pts. apiece) Answer questions in complete, well-written sentences WITHIN the spaces provided.

1. What is an absorption spectrum?
2. In our study of the CO molecule what is the classical energy?
3. Why do we use energy eigenstates?
4. What is Rutherford scattering?
5. Consider the spectral distribution shown below. What is the mathematical definition of the width of the distribution? Sketch on the plot what this width represents. Explain.

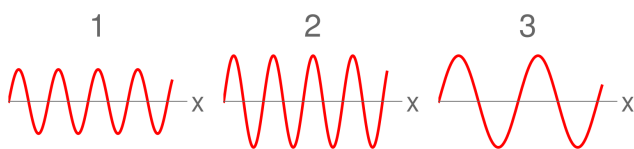


6. Cite at least two experimental measurements that violated classical physics and required quantum mechanics to explain. Discuss how they violated classical physics.

7. Why does the Sun shine? What are the important features of the process? Your answer should be descriptive and qualitative - not quantitative.

8. Recall that for a bound, one-dimensional system at  $t = 0$  that  $|\psi\rangle = \sum_{n=0}^{\infty} b_n |\phi_n\rangle$ . What is  $b_n$  and how is it related to a measurement of the momentum?

9. The figure below shows the de Broglie waves of three equal-mass particles as a function of position  $x$ . Rank them according to their speed. Explain your reasoning. The range of  $x$  and  $y$  are the same in each plot.



10. The components of the angular momentum vector  $\vec{L}$  are orthogonal to each other. Can a measurement of the  $x$  component of the angular momentum,  $L_x$ , of a state with  $m_z = \ell$  produce a nonzero result? Explain.

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**Problems.** Clearly show all work for full credit on a separate piece of paper.

1. (10 pts.) A mass  $m_0 = 0.910 \text{ kg}$  is oscillating freely on a vertical spring. The period for  $m_0$  is  $T_0 = 1.10 \text{ s}$ . An unknown mass  $m_1$  replaces  $m_0$  on the same spring and has a period of  $T_1 = 1.32 \text{ s}$ . What is the spring constant  $k$  and the unknown mass  $m_1$ ?

2. (10 pts.) Show that the frequencies of photons due to energy decays between successive levels of a rotator with momentum of inertia  $I$  are given by the following.

$$\hbar\omega = \frac{\hbar^2}{I}(\ell + 1) \quad \text{or} \quad \frac{\hbar^2}{I}\ell$$

3. (12 pts.) Find  $\psi(x)$  and  $P(E_n)$  at  $t = 0$  relevant to a one-dimensional box with walls at  $(0, a)$  for the following initial state.

$$\psi(x, 0) = A_3 (e^{i\pi(x-a)/a} - 1)$$

Make sure you get an expression for  $\psi(x)$  valid for all eigenstates. The eigenfunctions and eigenvalues of the one-dimensional particle in a box of width  $a$  are

$$|\phi(x)\rangle = \sqrt{\frac{2}{a}} \sin\left(\frac{n\pi x}{a}\right) \quad E_n = n^2 E_1 = n^2 \frac{\hbar^2 \pi^2}{2ma^2} \quad .$$

4. (13 pts.) A molecule behaves like a one-dimensional harmonic oscillator with a spring constant  $k = 2.42 \text{ eV/\AA}^2$ . It emits a photon of energy  $E_\gamma = h\nu = \hbar\omega_0 = 0.1 \text{ eV}$  going from the third excited state to the second excited state. The oscillating part is a proton ( $m_p = 938 \text{ MeV}/c^2$ ).

- a) What is the value of the classical turning point  $x_t$  if the molecule is in the  $n = 2$  state?
- b) What is the numerical value of  $\beta = \sqrt{m\omega_0/\hbar}$ ?
- c) What is the probability that a proton in this state is at a distance from the origin forbidden to it by classical mechanics? Get your answer in terms of  $x_t$  and  $\beta$  where  $\beta$  is defined below with the wave function for the  $n = 2$  state.

$$|\phi_2\rangle = \frac{1}{\sqrt{8\sqrt{\pi}}} (4\xi^2 - 2) e^{-\xi^2/2} \quad \xi = \beta x \quad \beta = \sqrt{\frac{m\omega_0}{\hbar}}$$

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5. (15 pts.) Suppose a rigid rotator is in the eigenstate of  $\hat{L}^2$  with  $\ell = 1$  and  $m_z = -1$  ( $Y_1^{-1}(\theta, \phi)$ ). We want to find the probability of obtaining the values of  $m_x = 0, \pm 1$  from a measurement of  $\hat{L}_x$ . We always measure eigenvalues so to get the results of a measurement of  $L_x$  we need to construct the appropriate operator  $\hat{L}_x$  which satisfies  $\hat{L}_x X = \alpha \hbar X$  where  $X$  is an eigenfunction of the  $\hat{L}_x$  operator and  $\alpha$  is the eigenvalue. To do that we can assume

$$X = aY_1^1 + bY_1^0 + cY_1^{-1}$$

since the spherical harmonics form a complete set and we know what they are. We restrict our attention to only  $\ell = 1$  states as a consequence of angular momentum conservation. Generate the set of simultaneous equations that the coefficients ( $a$ ,  $b$ ,  $c$ ) and  $\alpha$  must satisfy.

### Physics 309 Equations

$$R_T(\nu) = \frac{\text{Energy}}{\text{time} \times \text{area}} \quad E = h\nu = \hbar\omega \quad v_{\text{wave}} = \lambda\nu \quad I \propto |\vec{E}|^2 \quad \lambda = \frac{h}{p} \quad p = \hbar k \quad K = \frac{p^2}{2m} \quad K_{\text{max}} = h\nu - \Phi$$

$$-\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} \Psi(x, t) + V(x) \Psi(x, t) = i\hbar \frac{\partial}{\partial t} \Psi(x, t) \quad \hat{p}_x = -i\hbar \frac{\partial}{\partial x} \quad \hat{A} |\phi\rangle = a |\phi\rangle \quad \langle \hat{A} \rangle = \int_{-\infty}^{\infty} \psi^* \hat{A} \psi dx$$

$$\langle \phi_{n'} | \phi_n \rangle = \int_{-\infty}^{\infty} \phi_{n'}^* \phi_n dx = \delta_{n', n} \quad \langle \phi(k') | \phi(k) \rangle = \int_{-\infty}^{\infty} \phi_{k'}^* \phi_k dx = \delta(k - k') \quad e^{i\phi} = \cos \phi + i \sin \phi$$

$$|\psi\rangle = \sum b_n |\phi_n\rangle \rightarrow b_n = \langle \phi_n | \psi \rangle \quad |\phi\rangle = e^{\pm ikx} \quad |\psi\rangle = \int b(k) |\phi(k)\rangle dk \rightarrow b(k) = \langle \phi(k) | \psi \rangle$$

$$|\psi(x, t)\rangle = \sum b_n |\phi_n\rangle e^{-i\omega_n t} \quad |\psi(x, t)\rangle = \int b(k) |\phi(k)\rangle e^{-i\omega(k)t} dk \quad \Delta p \Delta x \geq \frac{\hbar}{2} \quad (\Delta x)^2 = \langle x^2 \rangle - \langle x \rangle^2$$

The wave function,  $\Psi(\vec{r}, t)$ , contains all we know of a system and its square is the probability of finding the system in the region  $\vec{r}$  to  $\vec{r} + d\vec{r}$ . The wave function and its derivative are (1) *finite*, (2) *continuous*, and (3) *single-valued* ( $\psi_1(a) = \psi_2(a)$  and  $\psi'_1(a) = \psi'_2(a)$ ).

$$V_{HO} = \frac{\kappa x^2}{2} \quad \omega = 2\pi\nu = \sqrt{\frac{\kappa}{m}} \quad E_n = (n + \frac{1}{2})\hbar\omega_0 = \hbar\omega \quad |\phi_n\rangle = e^{-\xi^2/2} H_n(\xi) \quad \xi = \beta x \quad \beta^2 = \frac{m\omega_0}{\hbar}$$

$$\psi_1 = \mathbf{t}\psi_3 = \mathbf{d}_{12}\mathbf{p}_2\mathbf{d}_{21}\mathbf{p}_1^{-1}\psi_3 \quad T = \frac{1}{|t_{11}|^2} \quad R + T = 1$$

$$\mathbf{d}_{ij} = \frac{1}{2} \begin{pmatrix} 1 + \frac{k_j}{k_i} & 1 - \frac{k_j}{k_i} \\ 1 - \frac{k_j}{k_i} & 1 + \frac{k_j}{k_i} \end{pmatrix} \quad \mathbf{p}_i = \begin{pmatrix} e^{-ik_i 2a} & 0 \\ 0 & e^{ik_i 2a} \end{pmatrix} \quad \mathbf{p}_i^{-1} = \begin{pmatrix} e^{ik_i 2a} & 0 \\ 0 & e^{-ik_i 2a} \end{pmatrix}$$

$$E = \frac{\hbar^2 k^2}{2m} \quad k = \sqrt{\frac{2m(E - V)}{\hbar^2}} \quad T = \frac{\text{transmitted flux}}{\text{incident flux}} \quad R = \frac{\text{reflected flux}}{\text{incident flux}} \quad \text{flux} = |\psi|^2 v$$

$$V(r) = \frac{Z_1 Z_2 e^2}{r} \quad E = \frac{1}{2} \mu v^2 + V(r) \quad \vec{R}_{cm} = \frac{\sum_i m_i \vec{r}_i}{\sum_i m_i} \quad \mu = \frac{m_1 m_2}{m_1 + m_2}$$

$$\psi(x) = \sum_{n=1}^{\infty} a_n x^n \quad \langle K \rangle = \frac{3}{2} kT \quad n(v) = 4\pi N \left( \frac{m}{2\pi k_B T} \right)^{3/2} v^2 e^{-mv^2/2k_B T} \quad \vec{L} = \vec{r} \times \vec{p} = \mathcal{I} \vec{\omega}$$

$$\mathcal{I} = \sum_i m_i r_1^2 = \int r^2 dm \quad KE_{rot} = \frac{L^2}{2\mathcal{I}} \quad E_\ell = \frac{\ell(\ell+1)\hbar^2}{2\mathcal{I}} \quad V_{coul} = \frac{Z_1 Z_2 e^2}{r} \quad ME = \frac{p_r^2}{2\mu} + \frac{L^2}{2\mu r^2} + V(r)$$

$$L_z |nlm\rangle = m\hbar |nlm\rangle \quad L^2 |nlm\rangle = \ell(\ell+1)\hbar^2 |nlm\rangle$$

$$\sin A \sin B = \frac{1}{2} [\cos(A-B) - \cos(A+B)] \quad \cos A \cos B = \frac{1}{2} [\cos(A-B) + \cos(A+B)]$$

$$\sin A \cos B = \frac{1}{2} [\sin(A-B) + \sin(A+B)] \quad \sin A + \sin B = 2 \sin \left( \frac{A+B}{2} \right) \cos \left( \frac{A-B}{2} \right)$$

### Constants

Speed of light ( $c$ )	$2.9979 \times 10^8 \text{ m/s}$	fermi ( $fm$ )	$10^{-15} \text{ m}$
Boltzmann constant ( $k_B$ )	$1.381 \times 10^{-23} \text{ J/K}$	angstrom ( $\text{\AA}$ )	$10^{-10} \text{ m}$
	$8.62 \times 10^{-5} \text{ eV/k}$	electron-volt ( $eV$ )	$1.6 \times 10^{-19} \text{ J}$
Planck constant ( $h$ )	$6.621 \times 10^{-34} \text{ J-s}$	MeV	$10^6 \text{ eV}$
	$4.1357 \times 10^{-15} \text{ eV-s}$	GeV	$10^9 \text{ eV}$
Planck constant ( $\hbar$ )	$1.0546 \times 10^{-34} \text{ J-s}$	Electron charge ( $e$ )	$1.6 \times 10^{-19} \text{ C}$
	$6.5821 \times 10^{-16} \text{ eV-s}$	$e^2$	$\hbar c/137$
Planck constant ( $\hbar c$ )	$197 \text{ MeV-fm}$	Electron mass ( $m_e$ )	$9.11 \times 10^{-31} \text{ kg}$
	$1970 \text{ eV-}\text{\AA}$		$0.511 \text{ MeV}/c^2$
Proton mass ( $m_p$ )	$1.67 \times 10^{-27} \text{ kg}$	atomic mass unit ( $u$ )	$1.66 \times 10^{-27} \text{ kg}$
	$938 \text{ MeV}/c^2$		$931.5 \text{ MeV}/c^2$
Neutron mass ( $m_n$ )	$1.68 \times 10^{-27} \text{ kg}$		
	$939 \text{ MeV}/c^2$		

### Integrals and Derivatives

$$\frac{df}{du} = \frac{df}{dx} \frac{du}{dx} \quad \frac{d}{dx}(x^n) = nx^{n-1} \quad \frac{d}{dx}(\sin x) = \cos x \quad \frac{d}{dx}(\cos x) = -\sin x \quad \frac{d}{dx}(e^{ax}) = ae^{ax}$$

$$\frac{d}{dx}(\ln ax) = \frac{1}{x} \quad \int x^n dx = \frac{x^{n+1}}{n+1} \quad \int e^{ax} dx = \frac{e^{ax}}{a} \quad \int \frac{1}{x} = \ln x \quad \int \frac{1}{\sqrt{x^2+a^2}} dx = \ln \left[ x + \sqrt{x^2+a^2} \right]$$

$$\int \frac{x}{\sqrt{x^2+a^2}} dx = \sqrt{x^2+a^2} \quad \int \frac{x^2}{\sqrt{x^2+a^2}} dx = \frac{1}{2} x \sqrt{x^2+a^2} - \frac{1}{2} a^2 \ln \left[ x + \sqrt{x^2+a^2} \right]$$

$$\int \frac{x^3}{\sqrt{x^2+a^2}} dx = \frac{1}{3} (-2a^2 + x^2) \sqrt{x^2+a^2} \quad \int x^2 \sin(ax) dx = \frac{2x \sin(ax)}{a^2} - \frac{(a^2 x^2 - 2) \cos(ax)}{a^3}$$

