

Particle in a Box

The potential

$$V = 0 \quad 0 < x < a$$

$$= \infty \quad \textit{otherwise}$$

Eigenfunctions and eigenvalues

$$|\phi_n\rangle = \sqrt{\frac{2}{a}} \sin\left(\frac{n\pi x}{a}\right) \quad E_n = n^2 \frac{\hbar^2 \pi^2}{2ma^2}$$

Superposition

$$|\psi\rangle = \sum_{n=1}^{\infty} b_n |\phi_n\rangle \quad \langle \phi_m | \phi_n \rangle = \delta_{m,n}$$

Getting the coefficients

$$b_n = \langle \phi_n | \psi \rangle \quad P_n = |b_n|^2$$

Time Dependence

$$\Psi(x, t) = \sum_{n=1}^{\infty} b_n |\phi_n(x)\rangle e^{-i\omega_n t}$$

Free Particle

The potential

$$V = 0$$

Eigenfunctions and eigenvalues

$$|\phi(k)\rangle = \frac{1}{\sqrt{2\pi}} e^{\pm ikx} \quad E = \frac{\hbar^2 k^2}{2m}$$

Superposition

$$|\psi\rangle = \int_{-\infty}^{\infty} b(k) \phi(k) dk$$

$$\langle \phi(k') | \phi(k) \rangle = \delta(k - k')$$

Getting the coefficients

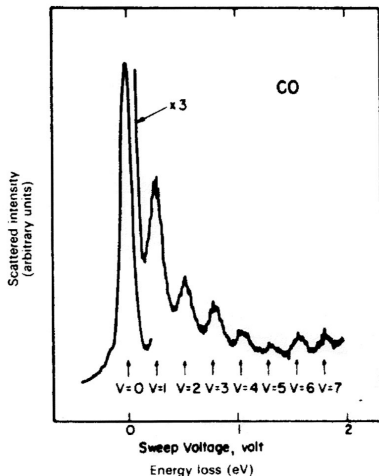
$$b(k) = \langle \phi(k) | \psi \rangle \quad P_n = |b(k)|^2$$

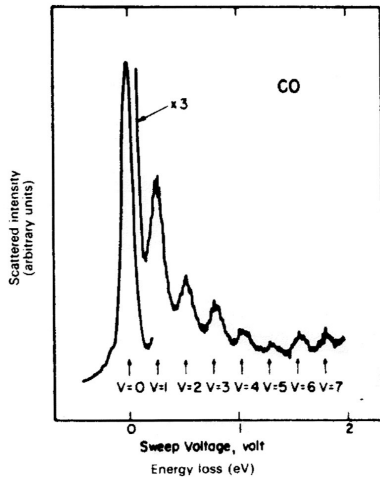
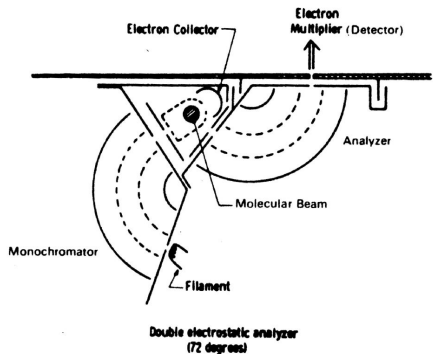
Time Dependence

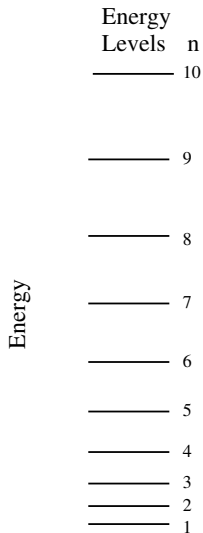
$$\Psi(x, t) = \int_{-\infty}^{\infty} b(k) \phi_k(x) e^{-i\omega(k)t} dk$$

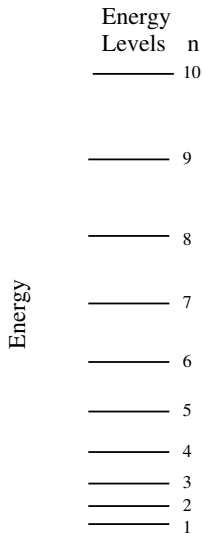
The excited states of the diatomic molecule carbon monoxide (CO) can be observed by crossing a beam of electrons with another beam of carbon monoxide. The energy spectrum of the scattered electrons is displayed here.

- 1 What is the simplest potential we used for a bound system?
- 2 What is the energy spectrum predicted for that potential? Does it fit here?
- 3 Find the eigenfunctions and eigenvalues of the harmonic oscillator. Does the energy spectrum reproduce the data?





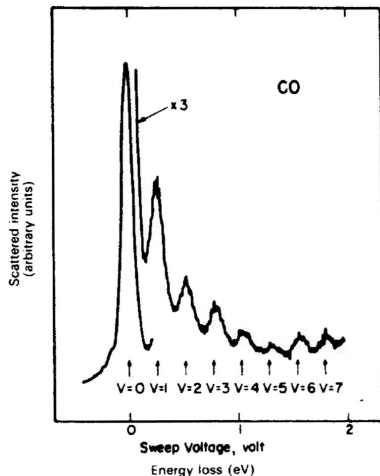


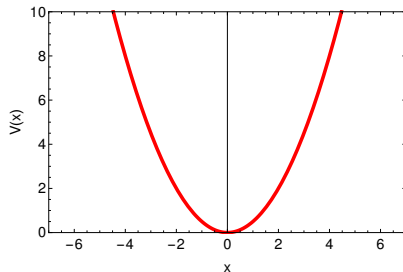


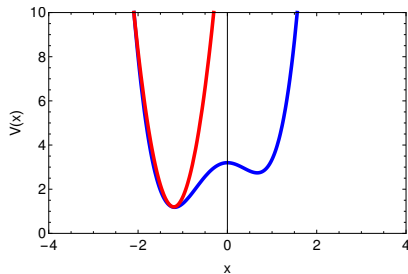
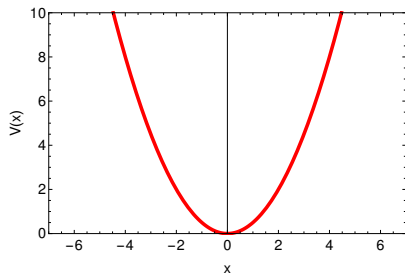
$$E_n = n^2 \frac{\hbar^2 \pi^2}{2ma^2}$$

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- 1 Each physical, measurable quantity, A , has a corresponding operator, \hat{A} , that satisfies the eigenvalue equation $\hat{A} \phi = a\phi$ and measuring that quantity yields the eigenvalues of \hat{A} .
- 2 Measurement of the observable A leaves the system in a state that is an eigenfunction of \hat{A} .
- 3 The state of a system is represented by a wave function Ψ which is continuous, differentiable and contains all the information about it.
 - The average value of any observable A is determined by
$$\langle A \rangle = \int_{all\ space} \Psi^* \hat{A} \Psi d\vec{r}.$$
 - The 'intensity' is proportional to $|\Psi|^2$.
- 4 The time development of the wave function is determined by

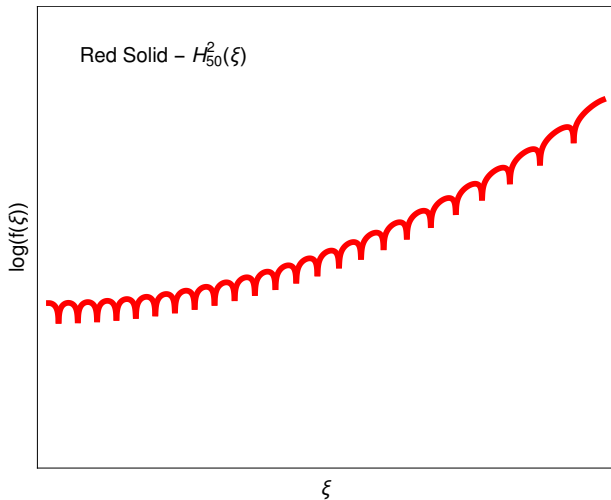
$$i\hbar \frac{\partial \Psi(\vec{r}, t)}{\partial t} = -\frac{\hbar^2}{2\mu} \nabla^2 \Psi(\vec{r}, t) + V(\vec{r}) \Psi(\vec{r}, t) \quad \mu \equiv \text{reduced mass.}$$

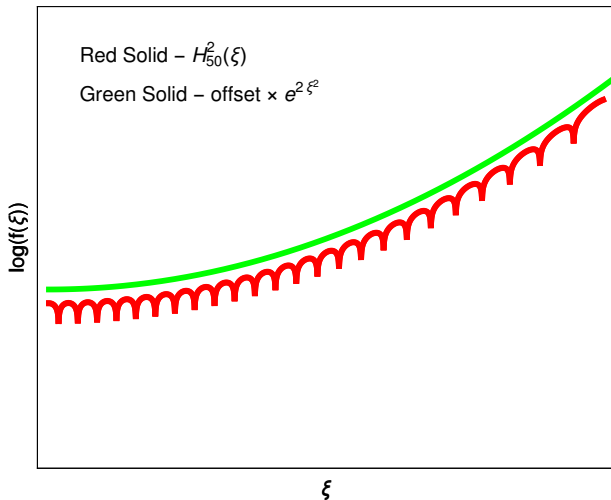
- 1 Potential energy: $\frac{\kappa X^2}{2}$
- 2 Hermite's equation:

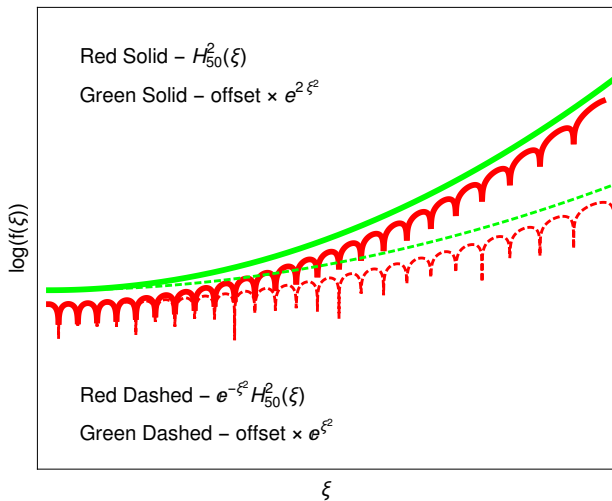
$$\frac{d^2 H}{d\xi^2} - 2\xi \frac{dH}{d\xi} + \left(\frac{\alpha}{\beta^2} - 1 \right) H = 0$$

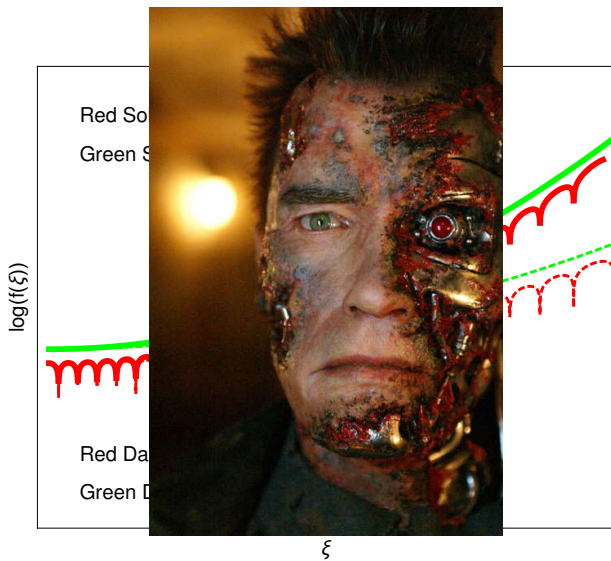
- 3 Second-order, linear, ordinary, homogeneous differential equation
 - 1 second-order: has a second derivative in it.
 - 2 linear: only derivatives to the first power.
 - 3 ordinary: one independent variable.
 - 4 homogeneous: equal to zero.
- 4 Method of Frobenius (19th century German mathematician)
 - 1 Used to generate an infinite series solution.
 - 2 Applies to equations of the form

$$u'' + \frac{p(z)}{z} u' + \frac{q(z)}{z^2} u = 0$$









$$H_0(\xi) = \frac{1}{\sqrt{\sqrt{\pi}}}$$

$$H_1(\xi) = \frac{1}{\sqrt{2\sqrt{\pi}}} 2\xi$$

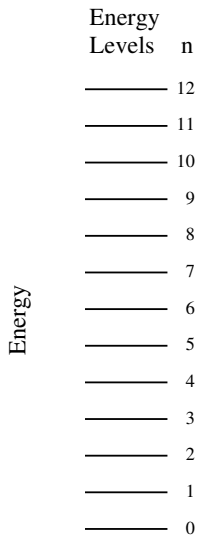
$$H_2(\xi) = \frac{1}{\sqrt{8\sqrt{\pi}}} (4\xi^2 - 2)$$

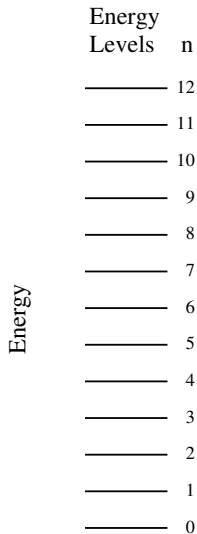
$$H_3(\xi) = \frac{1}{\sqrt{48\sqrt{\pi}}} (8\xi^3 - 12\xi)$$

$$H_4(\xi) = \frac{1}{\sqrt{384\sqrt{\pi}}} (16\xi^4 - 48\xi^2 + 12)$$

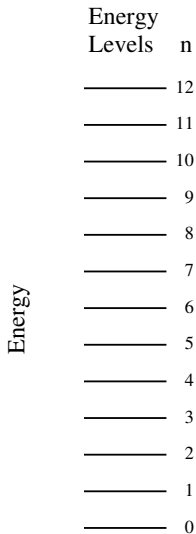
$$H_5(\xi) = \frac{1}{\sqrt{3840\sqrt{\pi}}} (32\xi^5 - 160\xi^3 + 120\xi)$$

$$H_6(\xi) = \frac{1}{\sqrt{46080\sqrt{\pi}}} (64\xi^6 - 480\xi^4 + 720\xi^2 - 120)$$

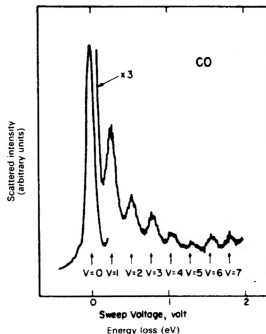


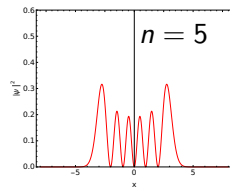
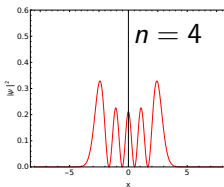
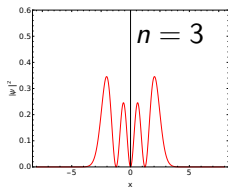
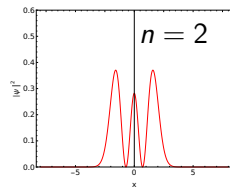
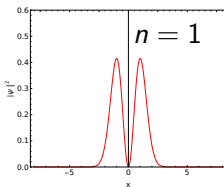
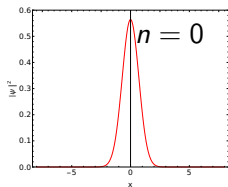


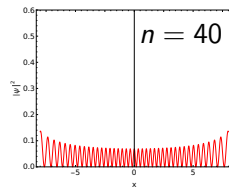
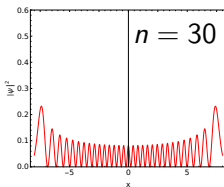
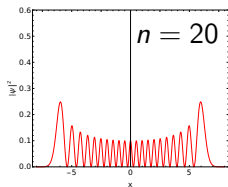
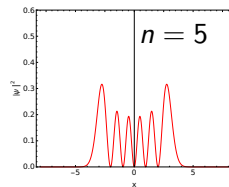
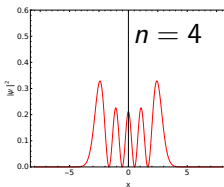
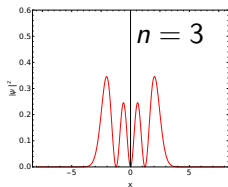
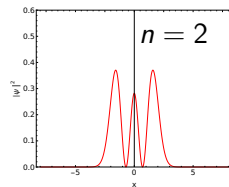
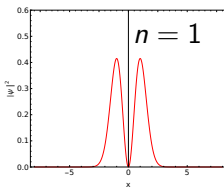
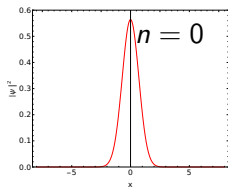
$$E_n = \left(n + \frac{1}{2} \right) \hbar \omega_0$$

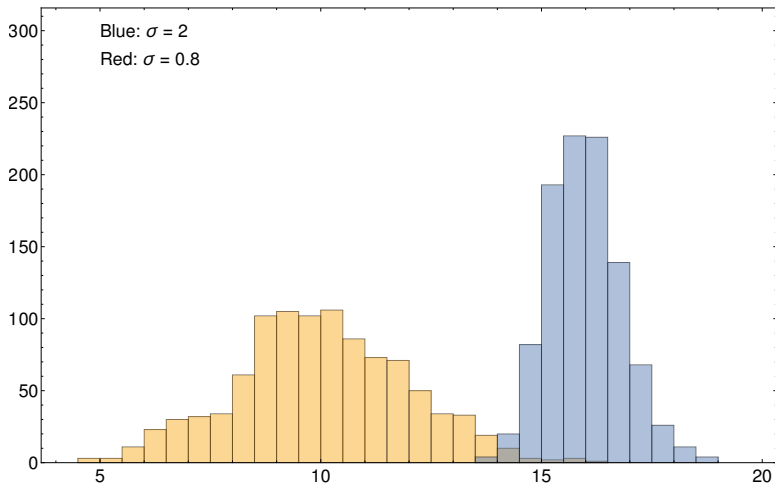


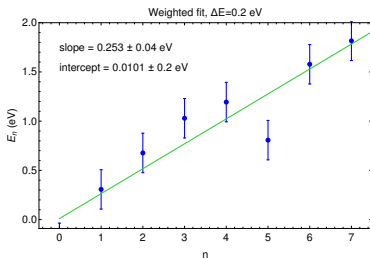
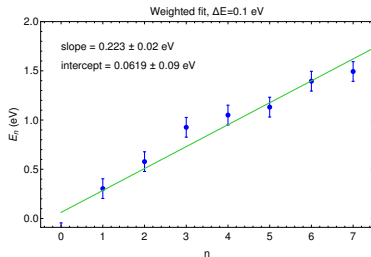
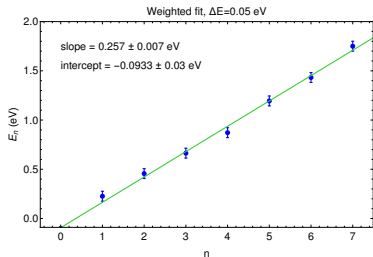
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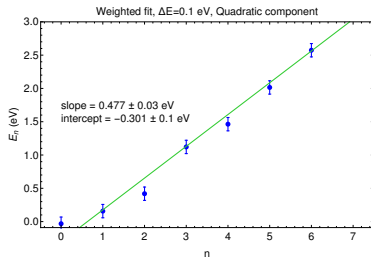
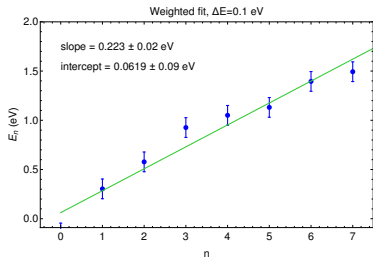












$$Q(\xi) = A e^{-\xi^2/2} H(\xi)$$



H.W.



$$\frac{d^2 H}{d\xi^2} - 2\xi \frac{dH}{d\xi} + \left(\frac{\alpha}{\beta^2} - 1\right) H = 0$$

Let

$$\text{Let } H(\xi) = \sum_{n=0}^{\infty} a_n \xi^n = a_0 + a_1 \xi + a_2 \xi^2 + a_3 \xi^3 + a_4 \xi^4 + a_5 \xi^5 + a_6 \xi^6 + \dots$$

$$\frac{dH}{d\xi} = a_1 + 2a_2 \xi + 3a_3 \xi^2 + 4a_4 \xi^3 + 5a_5 \xi^4 + 6a_6 \xi^5 + \dots$$

$$\frac{d^2H}{d\xi^2} = 2a_2 + 2 \cdot 3 a_3 \xi + 3 \cdot 4 a_4 \xi^2 + 4 \cdot 5 a_5 \xi^3 + 5 \cdot 6 a_6 \xi^4$$

$$\begin{aligned}
 & \left\{ 2a_2 + 2 \cdot 3 a_3 \xi + 3 \cdot 4 a_4 \xi^2 + 4 \cdot 5 a_5 \xi^3 + \dots \right\} + \xi^3 \left[4 \cdot 5 a_5 + \dots \right] \\
 & \left[-2a_1 \xi - 2 \cdot 2 a_2 \xi^2 - 2 \cdot 3 a_3 \xi^3 - 2 \cdot 4 a_4 \xi^4 + \dots \right] + \\
 & \left(\frac{\alpha}{\beta^2} - 1 \right) \left[a_0 + a_1 \xi + a_2 \xi^2 + a_3 \xi^3 + \dots \right] = 0 \\
 \hline
 & \xi^0 \left[1 \cdot 2 a_2 + \left(\frac{\alpha}{\beta^2} - 1 \cdot 2 \right) a_0 \right] + \\
 & \xi^1 \left[2 \cdot 3 a_3 + \left(\frac{\alpha}{\beta^2} - 1 - 2 \cdot 1 \right) a_1 \right] + \\
 & \xi^2 \left[3 \cdot 4 a_4 + \left(\frac{\alpha}{\beta^2} - 1 - 2 \cdot 2 \right) a_2 \right] +
 \end{aligned}$$

$$\sum^3 \left[4.5 a_5 + \left(\frac{\alpha}{\beta^2} - (-2.3) \right) a_3 \right] + \dots = 0$$