

*Jerry Gilfoyle for the CLAS Collaboration  
University of Richmond*

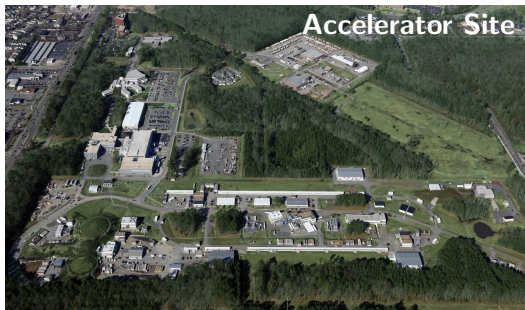


"The Periodic Table"

- Jefferson Lab's Mission
- What we know.
- What we don't know.
- What we measure.
- Experiments with CLAS12
- Concluding Remarks

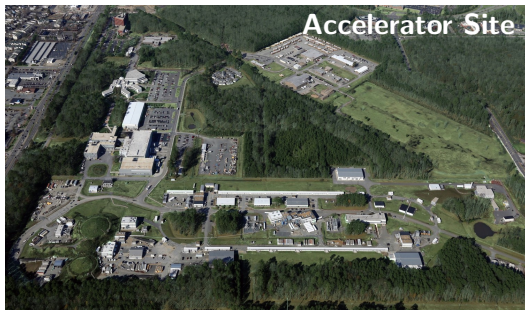
# What is the Mission of Jefferson Lab?

- Basic research into the quark nature of the atomic nucleus.
- Probe the quark-gluon structure of hadronic matter and how it evolves within nuclei.
- Test the theory of the color force Quantum Chromodynamics (QCD) and the nature of quark confinement.
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Solving QCD one of the seven Millenium Prize Problems from the Clay Mathematics Institute.

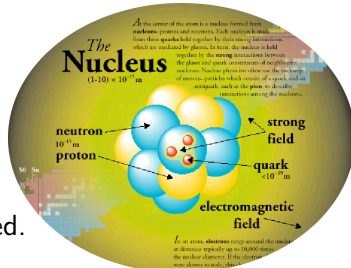
# What Do We Know?

- The Universe is made of quarks and leptons and the force carriers.

BOSONS			force carriers spin = 0, 1, 2, ...		
Unified Electroweak spin = 1			Strong (color) spin = 1		
Name	Mass GeV/c <sup>2</sup>	Electric charge	Name	Mass GeV/c <sup>2</sup>	Electric charge
$\gamma$ photon	0	0	<b>g</b> gluon	0	0
<b>W<sup>-</sup></b>	80.39	-1	Higgs Boson spin = 0		
<b>W<sup>+</sup></b>	80.39	+1	Name	Mass GeV/c <sup>2</sup>	Electric charge
<b>Z<sup>0</sup></b>	91.188	0	<b>H</b> Higgs	126	0
Z boson					

FERMIONS						matter constituents spin = 1/2, 3/2, 5/2, ...	
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- The atomic nucleus is made of protons and neutrons bound by the strong force.
- The quarks are confined inside the protons and neutrons.
- Protons and neutrons are NOT confined.



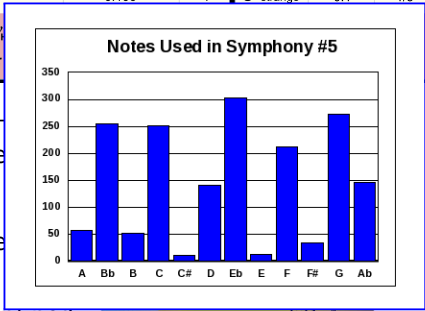
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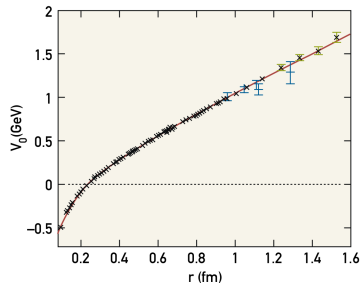
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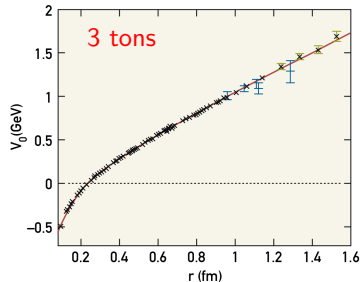
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- Quantum chromodynamics (QCD) looks like the right way to get the force at high energy.



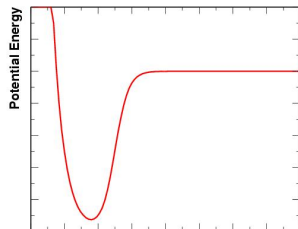
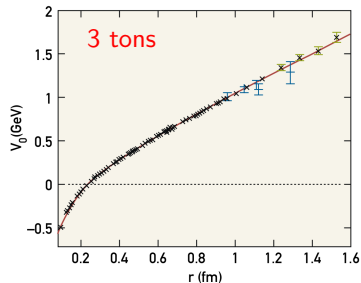
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- The hadronic model uses a phenomenological force fitted to data at low energy. This 'strong' force is the residual force between quarks.

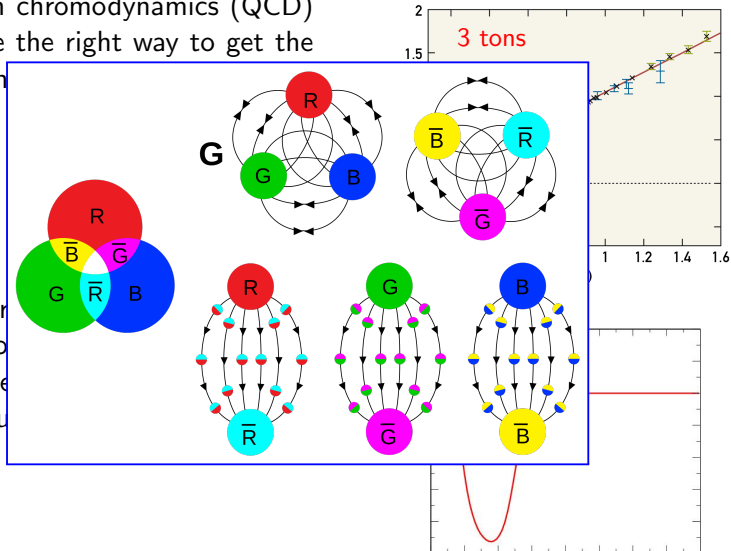




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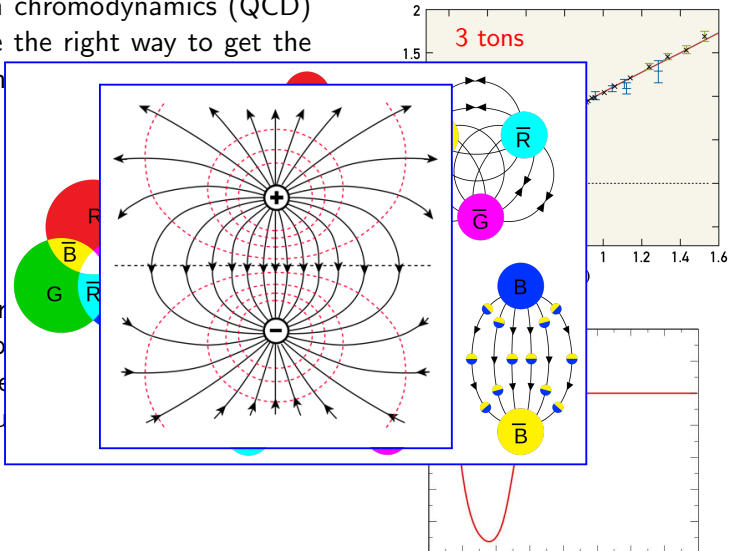
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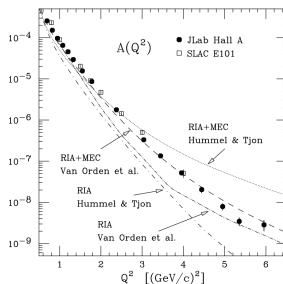
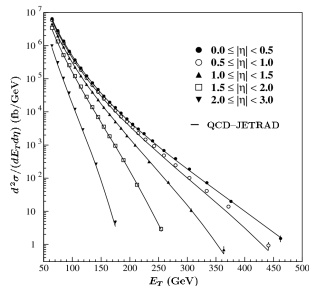
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# How Well Do We Know It?

- We have a working theory of strong interactions: quantum chromodynamics or QCD (B.Abbott, *et al.*, Phys. Rev. Lett., **86**, 1707 (2001)).
- The coherent hadronic model (the standard model of nuclear physics) works too (L.C.Alexa, *et al.*, Phys. Rev. Lett., **82**, 1374 (1999)).



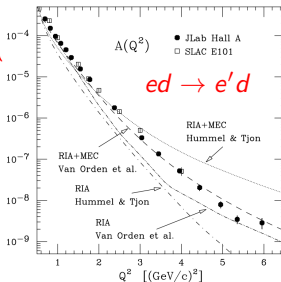
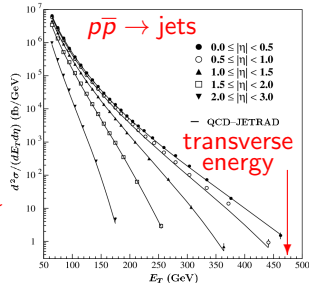
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effective target area

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4-momentum transfer squared



# What Don't We Know?

- Matter comes in pairs of quarks or triplets.
- We are mostly triplets (protons and neutrons).
- More than 99% of our mass is in nucleons.
- Proton  $\rightarrow$  2 ups + 1 down.
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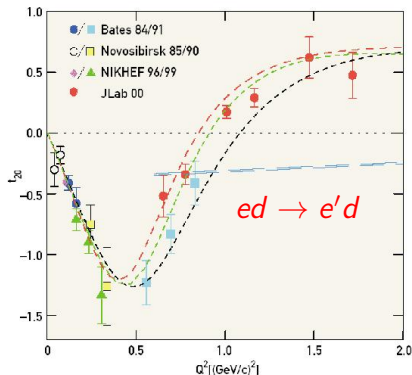
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- $m_n - m_p = 1.29333205(48) \text{ MeV}/c^2$  (exp)      Sz. Borsanyi et al. *Science* 347, 1452 (2015).
- $= 1.51(16)(23) \text{ MeV}/c^2$  (th)

# What Don't We Know?

- 1 We can't get QCD and the hadronic model to line up - D. Abbott, *et al.*, Phys. Rev Lett. **84**, 5053 (2000).
- 2 NEED TO FIGURE OUT QCD AT THE ENERGIES OF NUCLEI!!



# What Do We Measure?

The Magnetic Form Factor of the Neutron ( $G_M^n$ )

## The Magnetic Form Factor of the Neutron ( $G_M^n$ )

- Fundamental quantity related to the distribution of magnetization/currents in the neutron.
- Needed to extract the distribution of quarks in the neutron.
- Elastic form factors ( $G_M^n$ ,  $G_E^n$ ,  $G_M^p$ , and  $G_E^p$ ) provide key constraints on theory and the structure of hadrons.
- Part of a broad effort to understand how nucleons are 'constructed from the quarks and gluons of QCD'.\*

\* 'The Frontiers of Nuclear Science: A Long-Range Plan', NSF/DOE Nuclear Science Advisory Committee, April, 2007.

# How Do We Learn What's Inside the Nucleon?

- Nucleon elastic electromagnetic form factors (EEFFs) describe the distribution of charge and magnetization in the nucleon.
- They encode the deviations from point-particle behavior.
- Reveal the internal quark-gluon landscape of the nucleon and nuclei.
- We are in the region where the quarks get dressed.
- Rigorously test QCD in the non-perturbative regime.
- Jargon:  $G_E^p$ ,  $G_M^p$ ,  $G_E^n$ ,  $G_M^n$ .

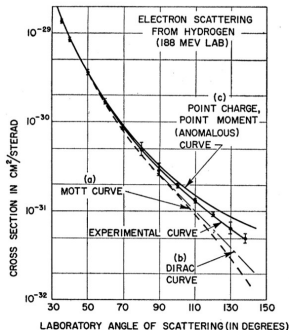
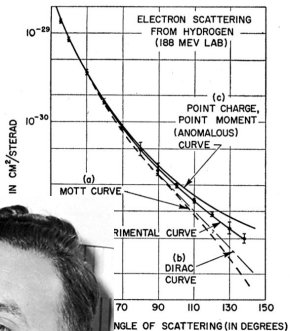


FIG. 5. Curve (a) shows the theoretical Mott curve for a spinless point proton. Curve (b) shows the theoretical curve for a point proton with the Dirac magnetic moment, curve (c) the theoretical curve for a point proton having the anomalous contribution in addition to the Dirac value of magnetic moment. The theoretical curves (b) and (c) are due to Rosenbluth.<sup>8</sup> The experimental curve falls between curves (b) and (c). This deviation from the theoretical curves represents the effect of a form factor for the proton and indicates structure within the proton, or alternatively, a breakdown of the Coulomb law. The best fit indicates a size of  $0.70 \times 10^{-12}$  cm.

McAllister and Hofstadter, PR 102, 851 (1956)

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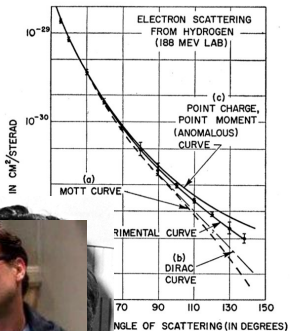
The theoretical Mott curve for a spinless nucleus shows the theoretical curve for a point nucleus having the anomalous contribution in the form of magnetic moment. The theoretical curve is due to Rosenbluth.<sup>8</sup> The experimental data (b) and (c). This deviation from the Mott curve is the effect of a form factor for the nucleus within the proton, or alternatively, the effect of a form factor for the nucleus within the neutron.

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The theoretical Mott curve for a spinless electron shows the theoretical curve for a point charge with no magnetic moment, curve (c) the theoretical curve for a point charge having the anomalous contribution in the form of magnetic moment. The theoretical curve due to Rosenbluth.<sup>8</sup> The experimental curves (b) and (c). This deviation from the Mott curve is the effect of a form factor for the electron within the proton, or alternatively, the Dirac law. The best fit indicates a size

Hofstadter, PR 102, 851 (1956)



Robert Hofstadter, Nobel Prize 1961

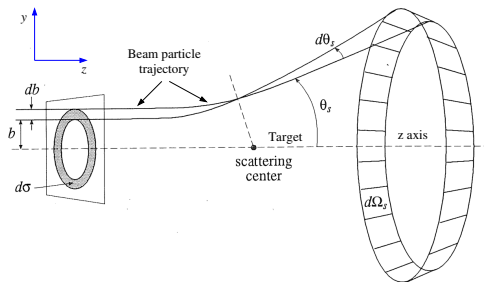
# What is a Form Factor?

- Start with the cross section.

$$\frac{d\sigma}{d\Omega} = \frac{\text{scattered rate/solid angle}}{\text{incident rate/surface area}}$$

- For elastic scattering use the Rutherford cross section.

$$\frac{d\sigma}{d\Omega} = \frac{Z_{tgt}^2 Z_{beam}^2 \alpha^2 (\hbar c)^2}{16E^2 \sin^4(\theta/2)}$$



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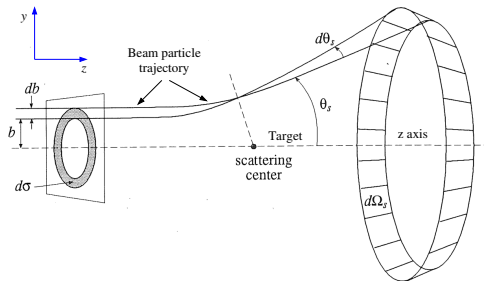
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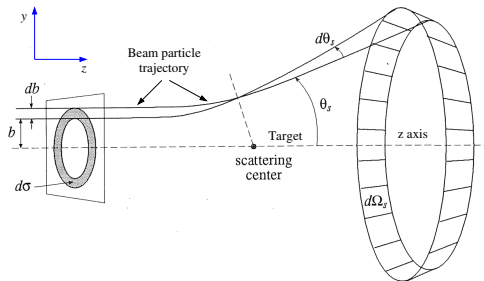
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- What happens when the beam is electrons and the target is not a point?

$$\frac{d\sigma}{d\Omega} = \frac{Z_{tgt}^2 \alpha^2 (\hbar c)^2}{16E^2 \sin^4(\theta/2)} \left(1 - \beta^2 \sin^2 \frac{\theta}{2}\right) |F(Q^2)|^2$$

where  $Q^2$  is the 4-momentum transfer.



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$$\frac{d\sigma}{d\Omega} = \frac{\text{scattered rate/solid angle}}{\text{incident rate/surface area}}$$

- For elastic scattering use the Rutherford cross section.

$$\frac{d\sigma}{d\Omega} = \frac{Z_{tgt}^2 Z_{beam}^2 \alpha^2 (\hbar c)^2}{16E^2 \sin^4(\theta/2)}$$

- Cross section for elastic scattering by point particles with spin.

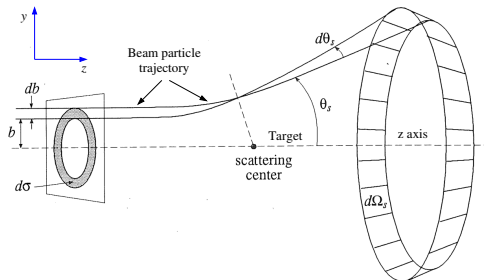
$$\frac{d\sigma}{d\Omega} = \frac{Z_{tgt}^2 Z_{beam}^2 \alpha^2 (\hbar c)^2}{16E^2 \sin^4(\theta/2)} \left(1 - \beta^2 \sin^2 \frac{\theta}{2}\right) \quad (\text{Mott cross section})$$

- What happens when the beam is electrons and the target is not a point?

$$\frac{d\sigma}{d\Omega} = \frac{Z_{tgt}^2 \alpha^2 (\hbar c)^2}{16E^2 \sin^4(\theta/2)} \left(1 - \beta^2 \sin^2 \frac{\theta}{2}\right) |F(Q^2)|^2$$

where  $Q^2$  is the 4-momentum transfer.

**THE FORM FACTOR!**



# Some Necessary Background

- EEFs cross section described with Dirac ( $F_1$ ) and Pauli ( $F_2$ ) form factors

$$\frac{d\sigma}{d\Omega} = \sigma_{Mott} \left[ (F_1^2 + \kappa^2 \tau F_2^2) + 2\tau (F_1 + \kappa F_2)^2 \tan^2 \left( \frac{\theta_e}{2} \right) \right]$$

where

$$\sigma_{Mott} = \frac{\alpha^2 E' \cos^2(\frac{\theta_e}{2})}{4E^3 \sin^4(\frac{\theta_e}{2})}$$

and  $\kappa$  is the anomalous magnetic moment,  $E$  ( $E'$ ) is the incoming (outgoing) electron energy,  $\theta$  is the scattered electron angle and  $\tau = Q^2/4M^2$ .

- For convenience use the Sachs form factors.

$$\frac{d\sigma}{d\Omega} = \frac{\sigma_{Mott}}{\epsilon(1+\tau)} (\epsilon G_E^2 + \tau G_M^2)$$

where

$$G_E = F_1 - \tau F_2 \quad \text{and} \quad G_M = F_1 + F_2 \quad \text{and} \quad \epsilon = \left[ 1 + 2(1 + \tau) \tan^2 \frac{\theta_e}{2} \right]^{-1}$$

# Why Should You Care?

- The chain of reason.

$$\frac{d\sigma}{d\Omega} \rightarrow |F(Q^2)|^2 \Leftrightarrow F(Q^2) \leftarrow \rho(\vec{r}) \leftarrow \psi(\vec{r}) \leftarrow \begin{matrix} \text{QCD,} \\ \text{Constituent quarks} \end{matrix}$$

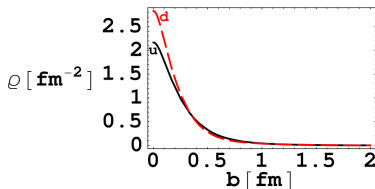
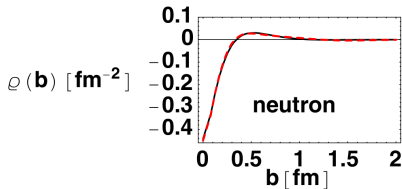
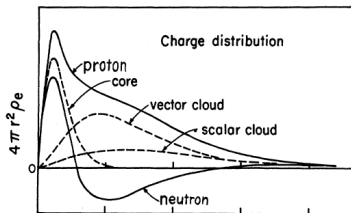
**Experiment**                      **Comparison**                      **Theory**

The form factors are the meeting ground between theory and experiment.

- The Fourier transform of the form factors are related to the charge and current distributions within the neutron.

# Why Should You Care Even More?

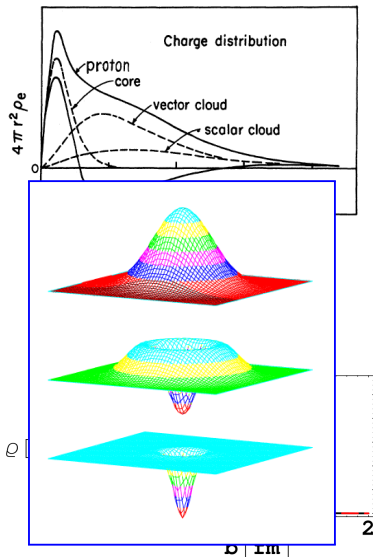
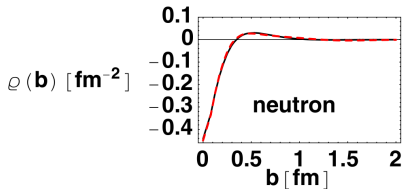
- The old picture of the neutron (and proton).
- What we know now - analysis of form factor data by G. Miller(Phys. Rev. Lett. **99**, 112001 (2007)).



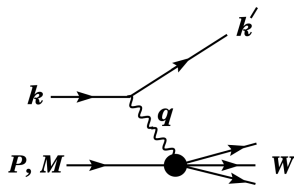


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# Some Kinematics Formulae



$\nu = E - E'$  where  $E$  and  $E'$  are the initial and final lepton energies in the nucleon rest frame.

$Q^2 = -q^2 = 2(EE' - \vec{k} \cdot \vec{k}') - m_l^2 - m_{l'}^2 \approx 4EE' \sin^2(\theta/2)$  where  $\theta$  is the lepton's scattering angle with respect to the lepton beam direction.

$x = \frac{Q^2}{2m\nu}$  where  $x$  is the fraction of the nucleon momentum carried by the struck quark.

$y = \frac{\nu}{E}$  is the fraction of the lepton's energy lost in the nucleon rest frame.

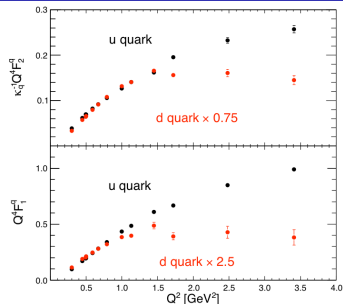
$W^2 = (P + q)^2 = M^2 + 2M\nu - Q^2$  is the square of the mass recoiling against the scattered lepton.

# What We'll Learn - Flavor Decomposition

- With all four EEFs we can unravel the contributions of the  $u$  and  $d$  quarks.
- Assume charge symmetry, no  $s$  quarks and use (Miller *et al.* Phys. Rep. **194**, 1 (1990))

$$F_{1(2)}^u = 2F_{1(2)}^P + F_{1(2)}^n \quad F_{1(2)}^d = 2F_{1(2)}^n + F_{1(2)}^P$$

- Evidence of di-quarks?  $d$ -quark scattering probes the diquark.



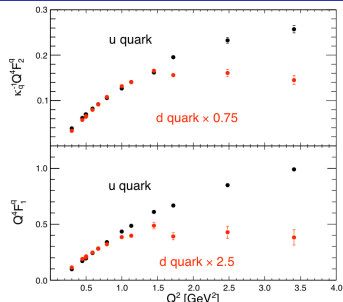
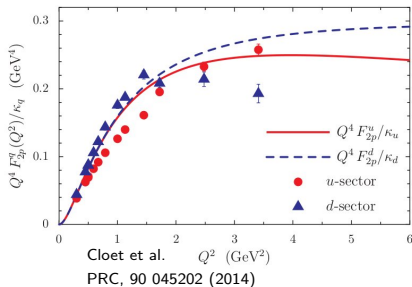
Cates et al. PRL **106**, 252003 (2011).

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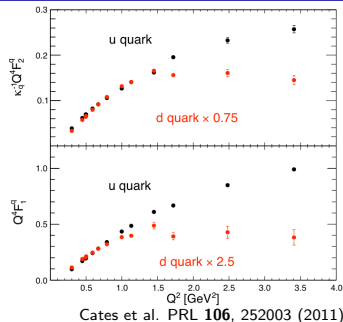
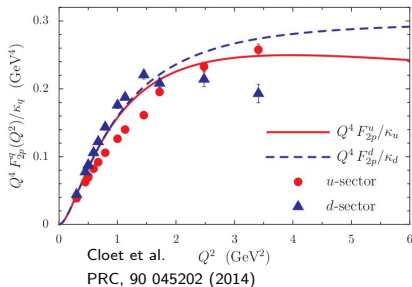
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The JLab program will double our reach in  $Q^2$  to  $\approx 8$  GeV<sup>2</sup>.

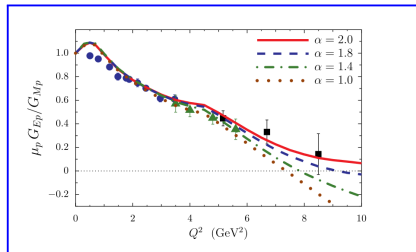
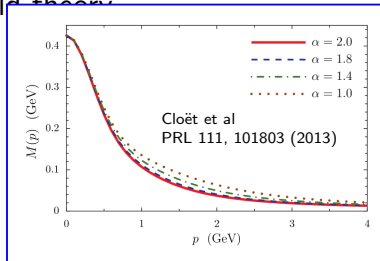
# What We'll Learn - Dyson-Schwinger Eqs

- Equations of motion of quantum field theory

- Infinite set of coupled integral equations.
- Inherently relativistic, non-perturbative, connected to QCD.
- Deep connection to confinement, dynamical chiral symmetry breaking.
- Infinitely many equations, gauge dependent  $\rightarrow$  Choose well!

- Recent results (Cloët et al).

- Model the nucleon dressed quark propagator as a quark-diquark.
- Damp the shape of the mass function  $M(p)$ .



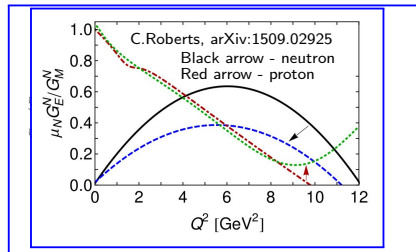
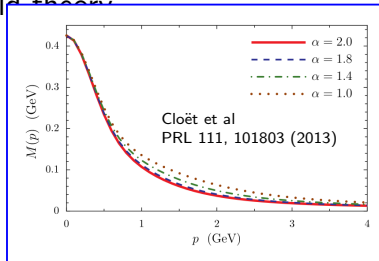
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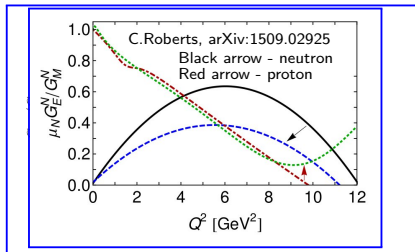
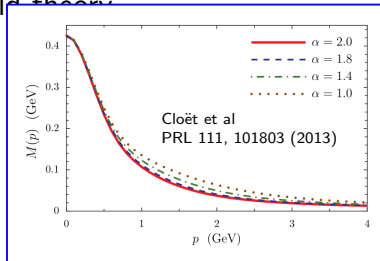
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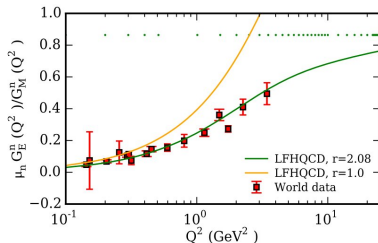
Position of zero in  $\mu_p G_E^p / G_M^p$  and  $\mu_n G_E^n / G_M^n$  sensitive to shape of  $M(p)$ !





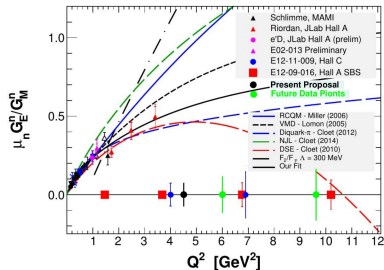
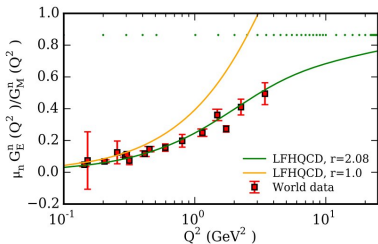
# What We'll Learn - Light Front Holographic QCD

- 1 Based on connections between light-front dynamics, it's holographic mapping to anti-de Sitter space, and conformal quantum mechanics.
- 2 Recent paper by Sufian *et al.* (Phys. Rev. D95, 01411 (2017)) included calculations of the electromagnetic form factors that include higher order Fock components  $|qqqq\bar{q}\rangle$ .
- 3 Obtain good agreement with all the form factor data with only three parameters, e.g.  $\mu_n G_E^n / G_M^n$ .



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# What We'll Learn - The Campaign

## The JLab Lineup

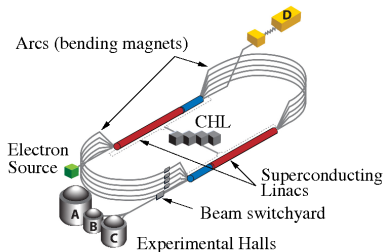
Quantity	Exp	Method	Target	$Q^2$ (GeV <sup>2</sup> )	Hall	Status
$G_M^p$	E12-07-108	Elastic scattering	$LH_2$	2.0 – 15.7	A	PRL*
$G_E^p/G_M^p$	E12-07-109	Polarization transfer	$LH_2$	6.4 – 10.5	A	fall, '23
$G_M^n$	E12-07-104	$E - p/e - n$ ratio	$LD_2, LH_2$	4.5 – 10.0	B	Complete <sup>†</sup>
$G_M^n$	E12-09-019	$E - p/e - n$ ratio	$LD_2, LH_2$	1.9 – 9.9	A	Complete <sup>†</sup>
$G_E^n/G_M^n$	E12-09-016	Double polarization asymmetry	polarized $^3\text{He}$	2.1 – 8.4	A	Complete <sup>†</sup>
$G_E^n/G_M^n$	E12-17-004	Polarization transfer	$LD_2$	4.3	A	Summer, '23
$G_E^n/G_M^n$	E12-11-009	Polarization transfer	$LD_2$	up to 6.9	A	To be scheduled

\* Phys. Rev. Lett., 128, 102002 (2022).

<sup>†</sup> Data collection complete.

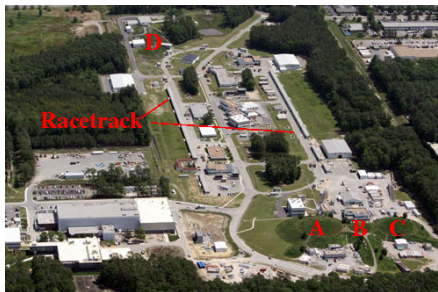
# How Do We Measure $G_M^n$ on a Neutron? (Step 1)

- Start at your local mile-long, high-precision, 12-GeV electron accelerator.
- The Continuous Electron Beam Accelerator Facility (CEBAF) produces beams of unrivaled quality.
- Electrons do up to five laps, are extracted, and sent to one of three experimental halls.
- All four halls can run simultaneously.



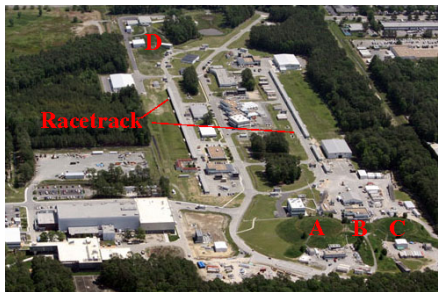
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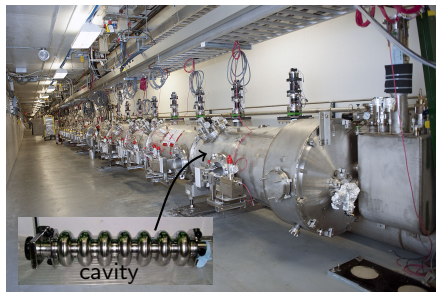
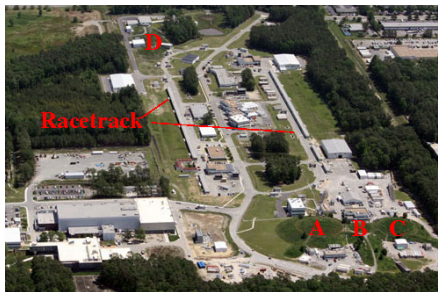
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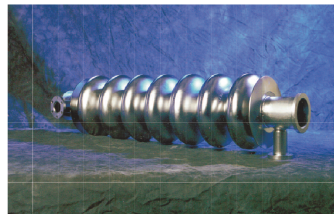
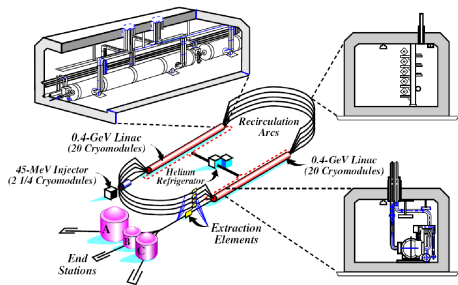
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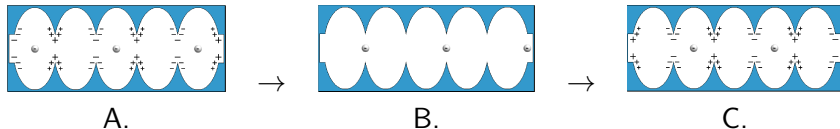
# How Does CEBAF Do That?

Accelerate your electrons to high energy.



*Cavity*

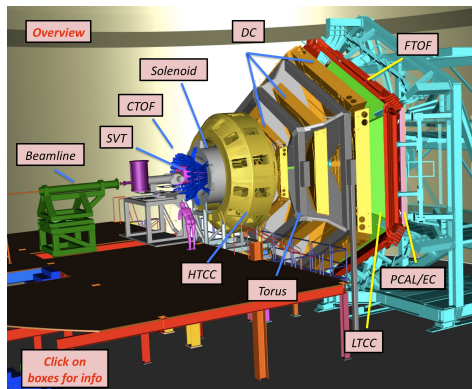
What happens inside the cavity? Feed it with oscillating, radio-frequency power at 1.5 GHz! In each half beam buckets are about 2 picoseconds long and arrive every 2 nanoseconds.





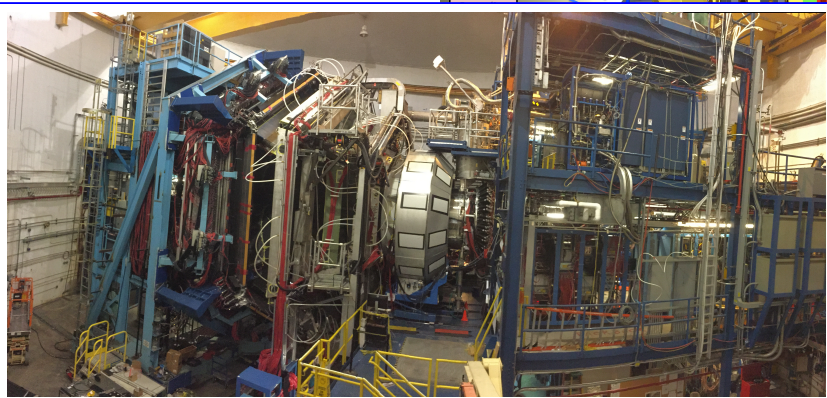
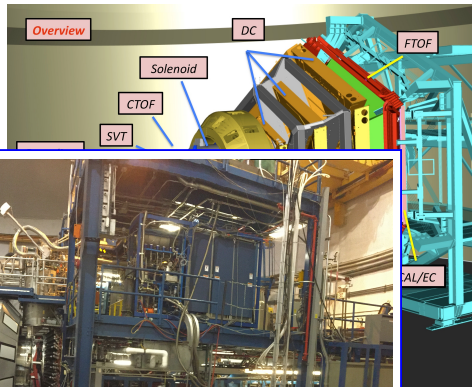
# How Do We Measure $G_M^n$ on a Neutron? (Step 2)

- Add one 45-ton, \$80-million radiation detector: the CEBAF Large Acceptance Spectrometer (CLAS12).
- CLAS covers a large fraction of the total solid angle at forward angles.
- Has about 62,000 detecting elements in about 40 layers.



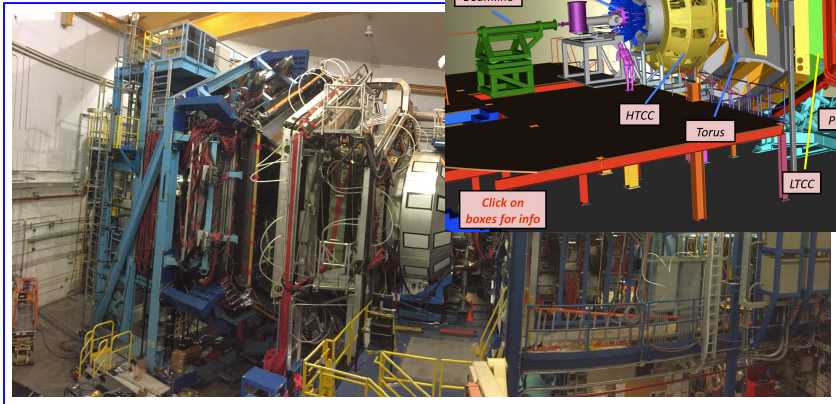
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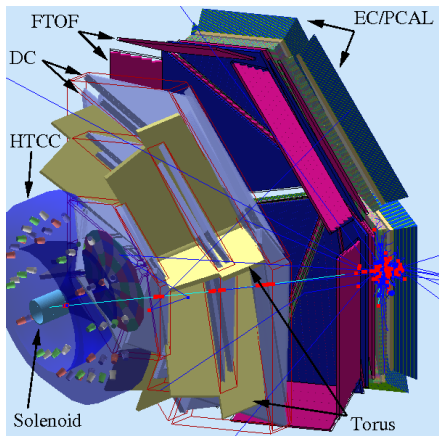
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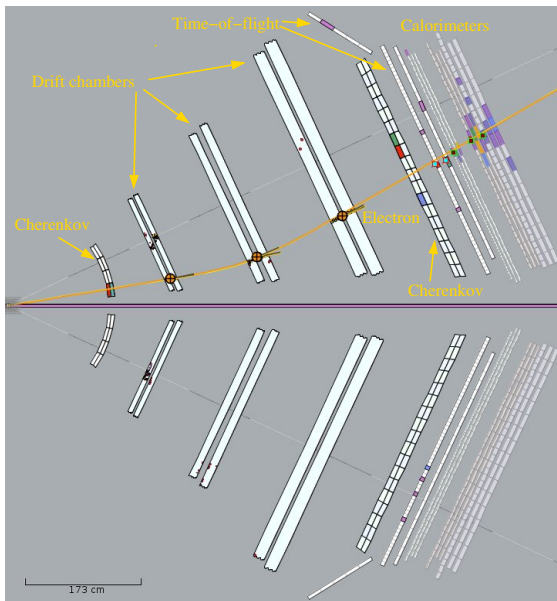


## How Do We Measure $G_M^n$ on a Neutron? (Step 2a)

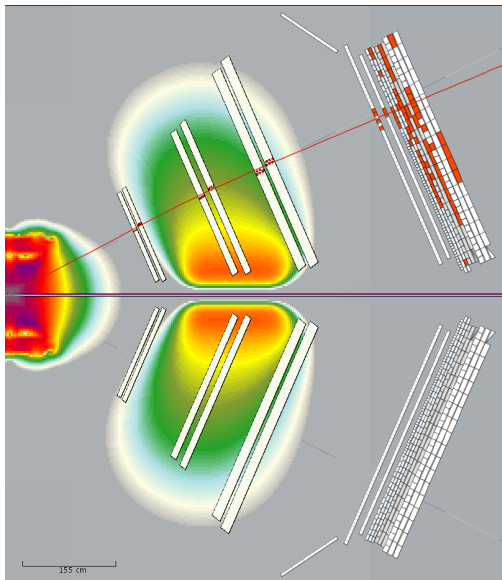
- Drift chambers map the trajectories. A toroidal magnetic field bends the particles to measure momentum.
- Other layers measure energy, time-of-flight, and particle identification.
- Each collision is reconstructed and the intensity pattern reveals the forces and structure of the colliding particles.
- Scatter electrons off protons and deuterons (proton+neutron).



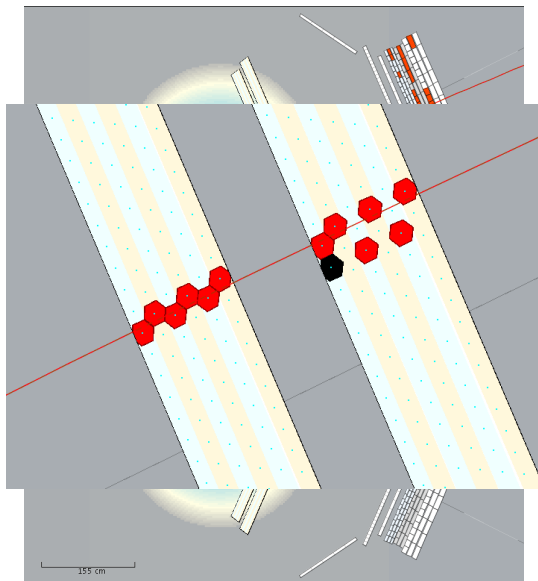
# A CLAS12 Event



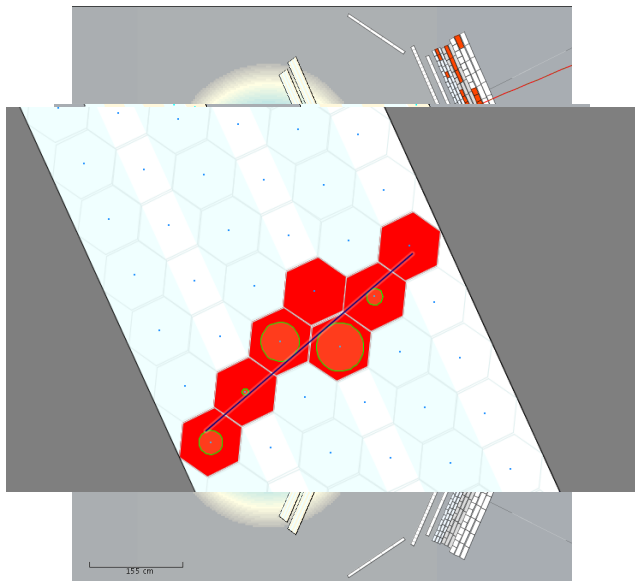
# A Simulated CLAS12 Event - Drift Chamber close-up



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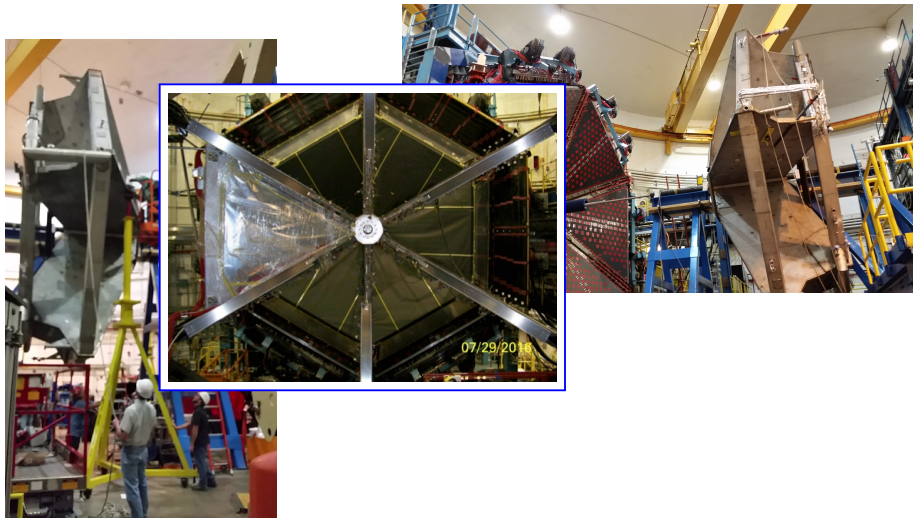




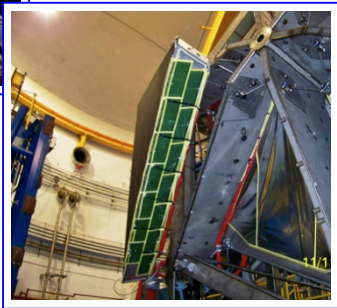
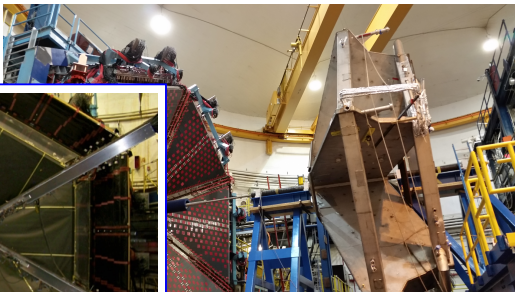
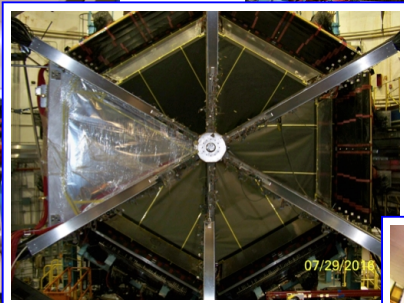
# A Real CLAS12 Event - Building the Drift Chambers



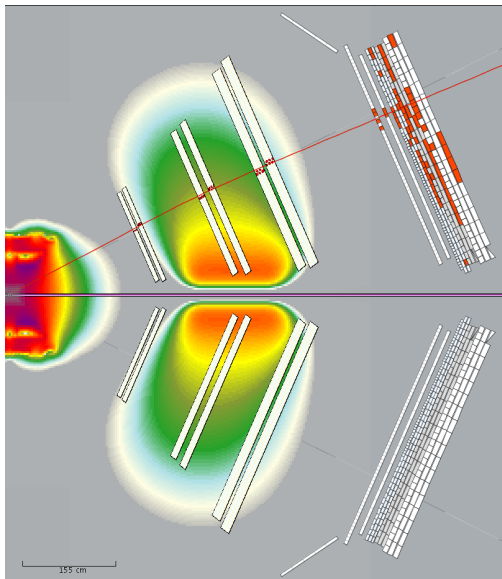
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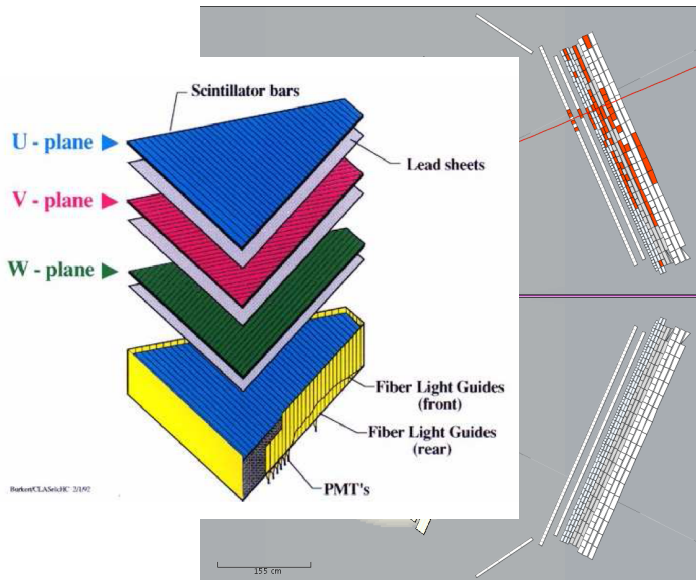
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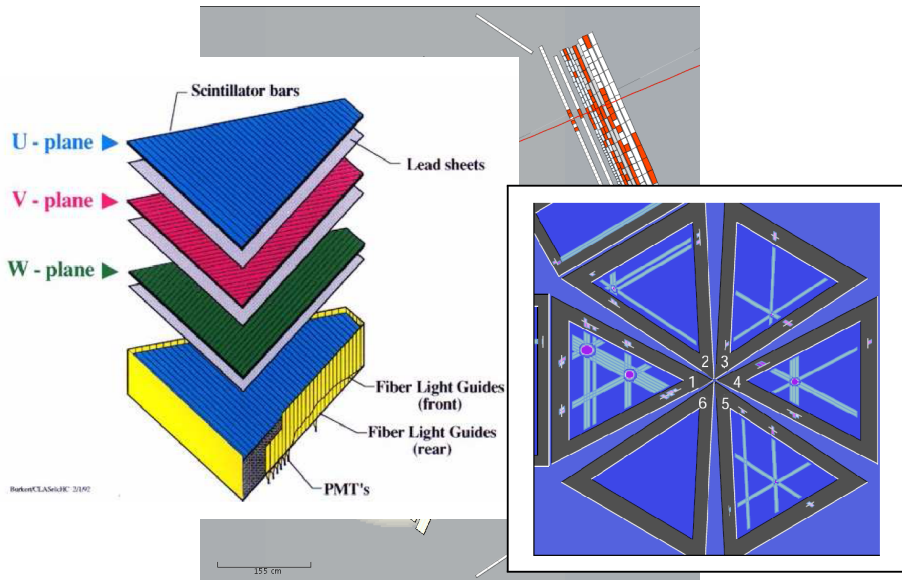
# A Simulated CLAS12 Event - Calorimeter close-up



# A Simulated CLAS12 Event - Calorimeter close-up

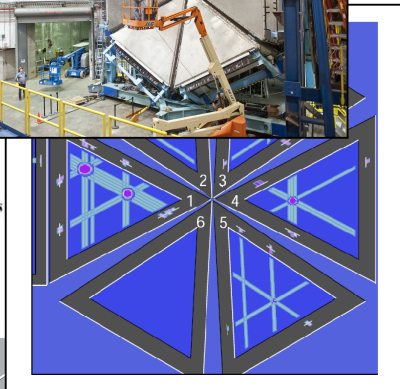
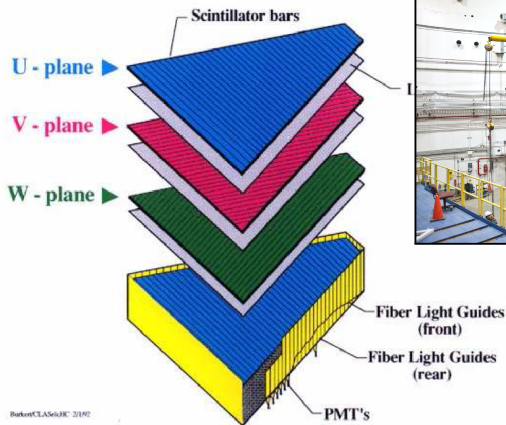


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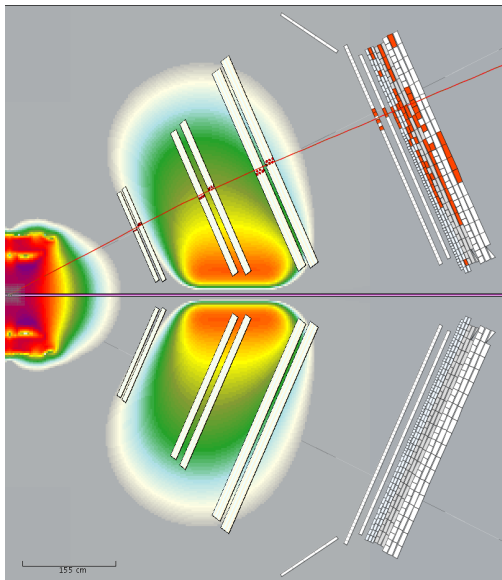


Berkley CLAS12/HC 2/10/12

# A Simulated CLAS12 Event - Calorimeter close-up

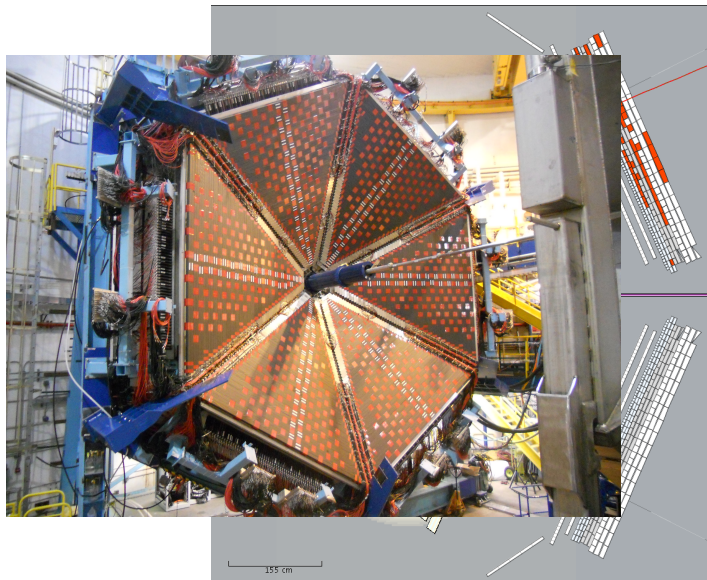


# A Simulated CLAS12 Event - Time-of-Flight close-up

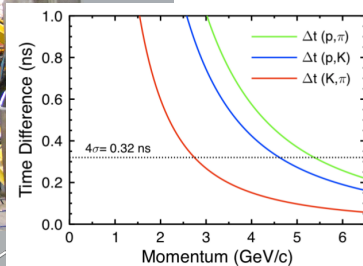
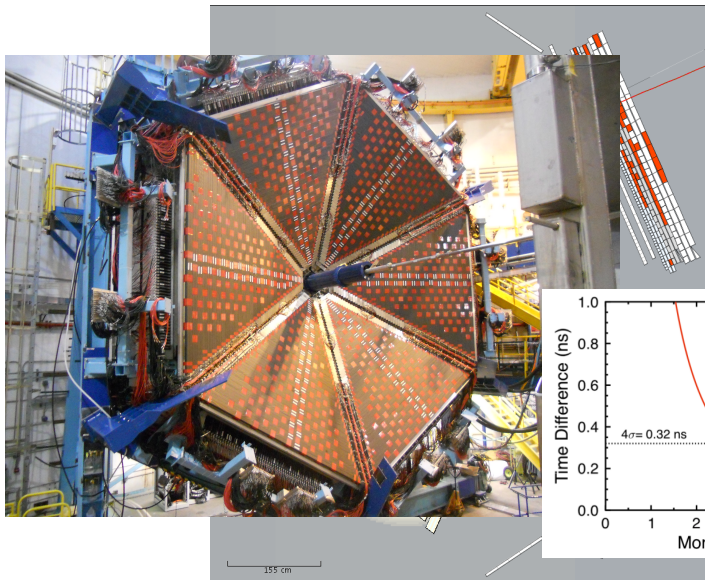




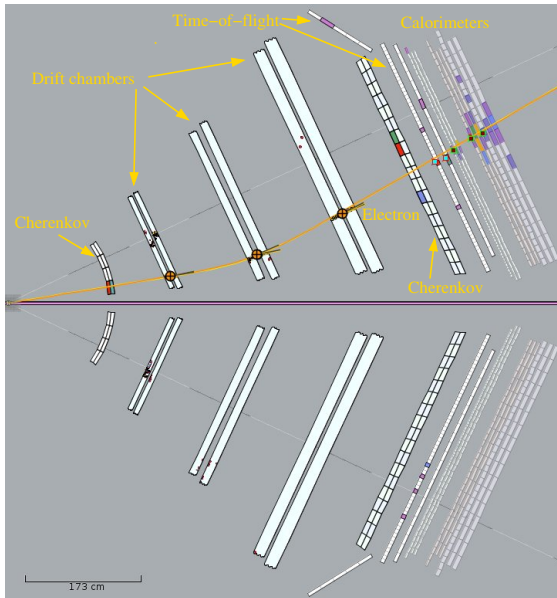
# A Simulated CLAS12 Event - Time-of-Flight close-up



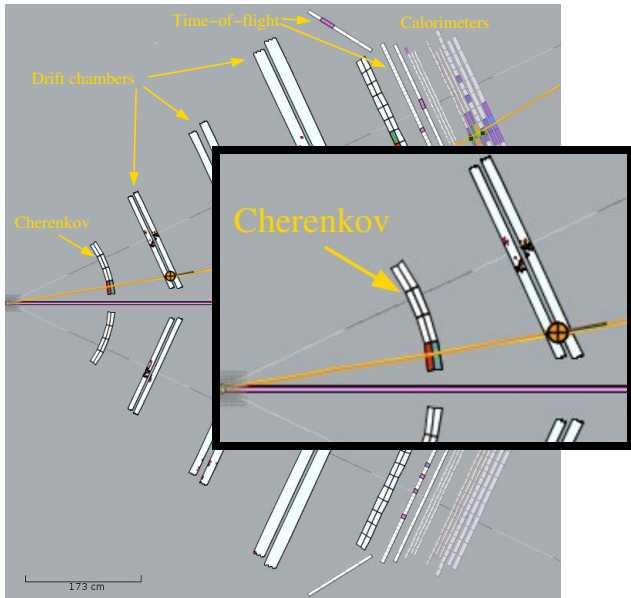
# A Simulated CLAS12 Event - Time-of-Flight close-up



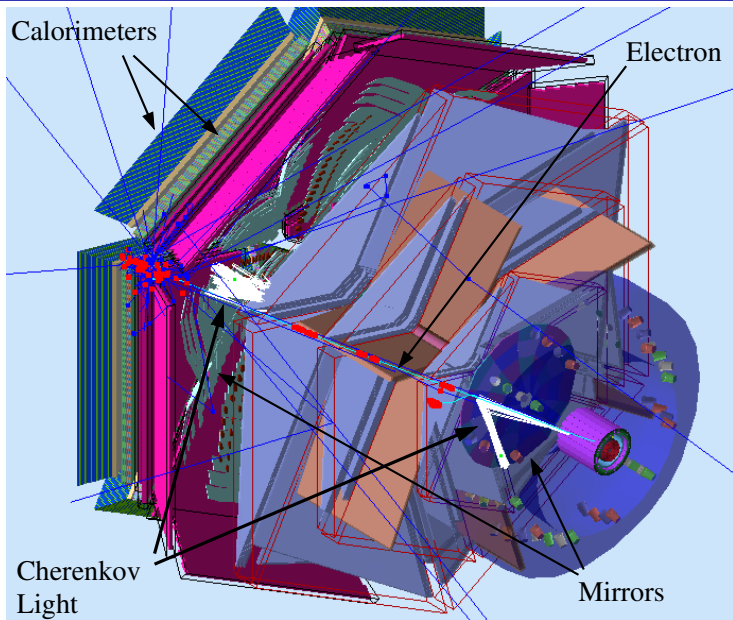
# A CLAS12 Event



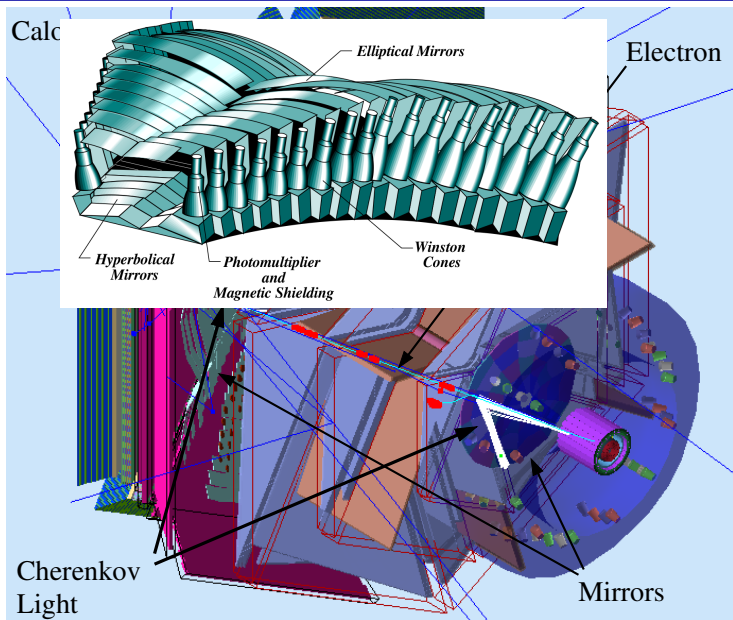
# A CLAS12 Event - Cherenkov close-up



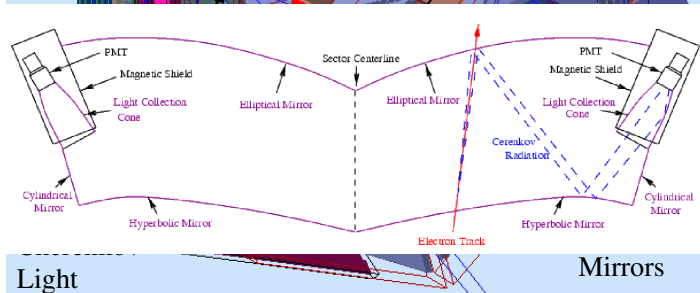
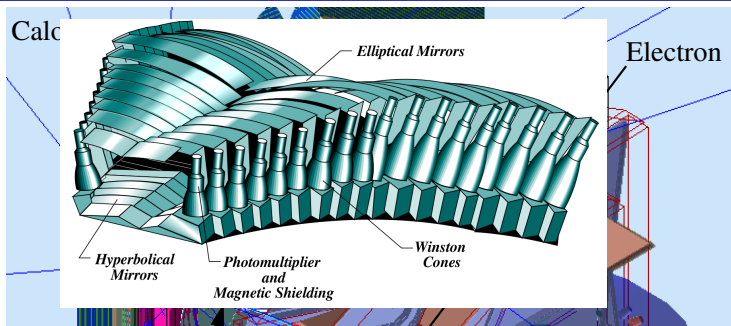
# A CLAS12 Event - Cherenkov close-up



# A CLAS12 Event - Cherenkov close-up

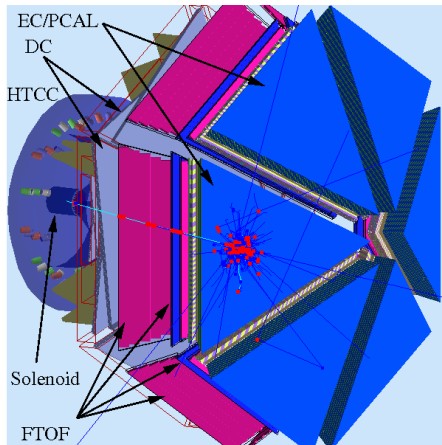


# A CLAS12 Event - Cherenkov close-up



## Forward Detector

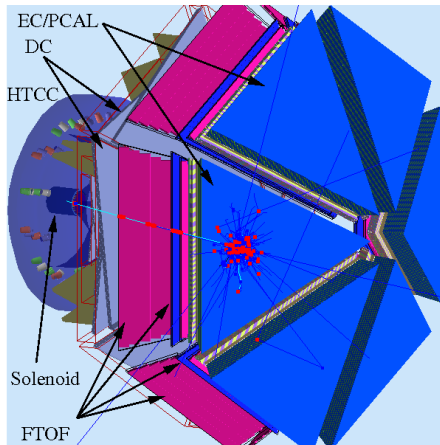
## Central Detector



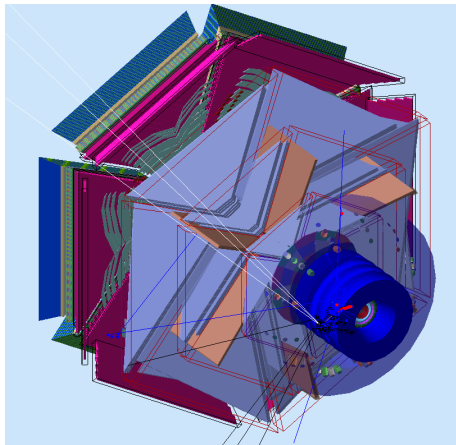


# Simulated CLAS12 Events

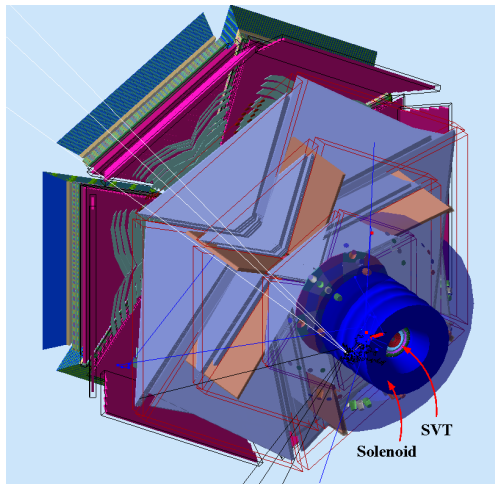
## Forward Detector



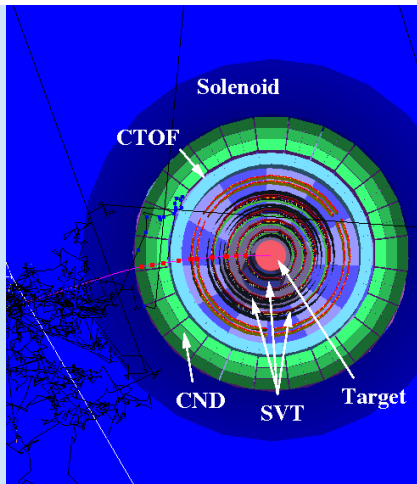
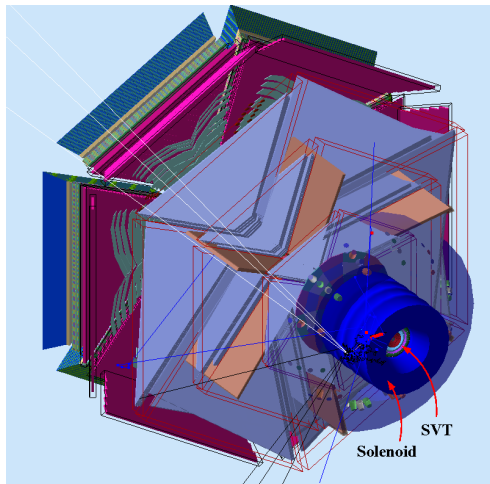
## Central Detector



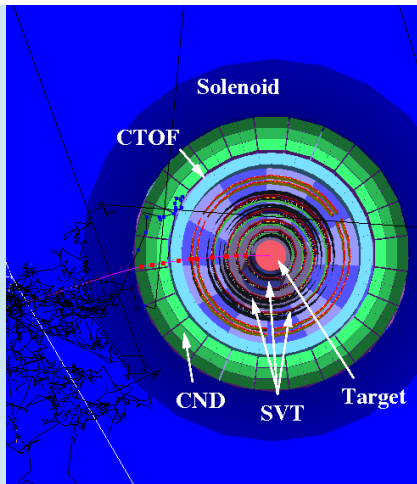
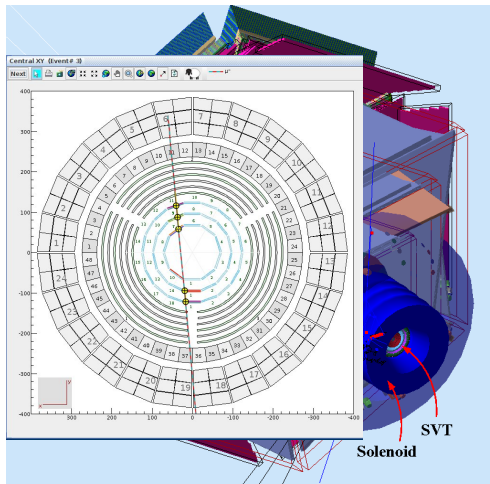
# Simulated CLAS12 Events - Silicon Vertex Tracker (SVT)



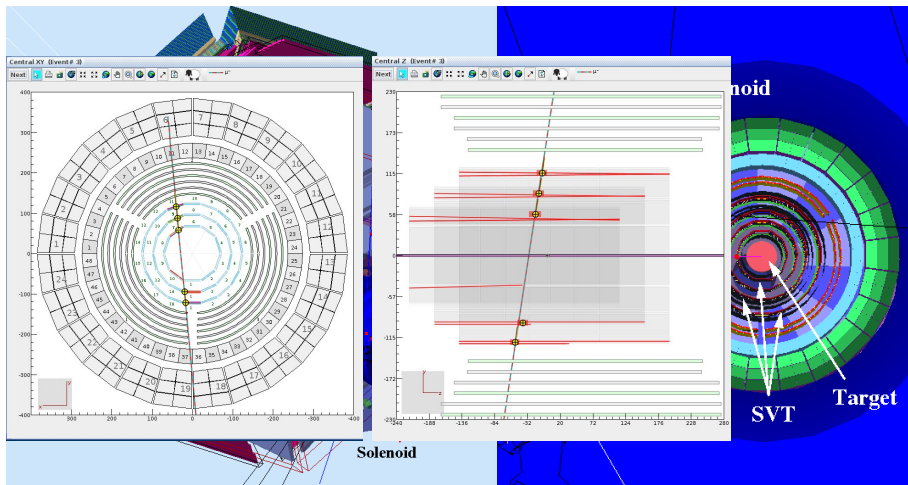
# Simulated CLAS12 Events - Silicon Vertex Tracker (SVT)



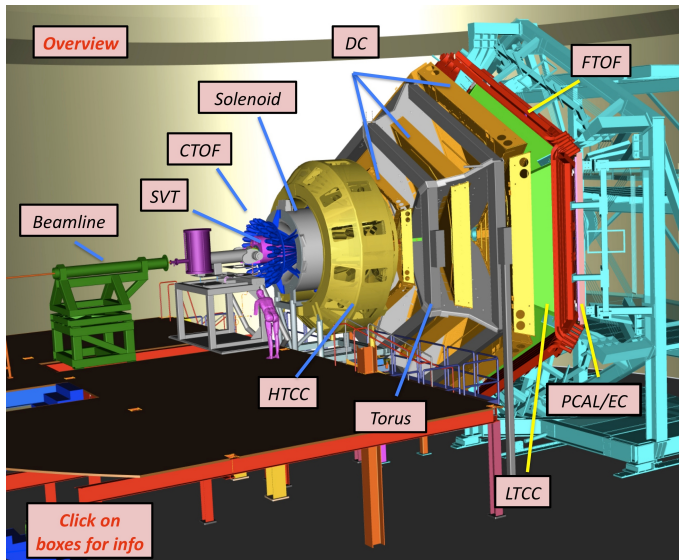
# Simulated CLAS12 Events - Silicon Vertex Tracker (SVT)



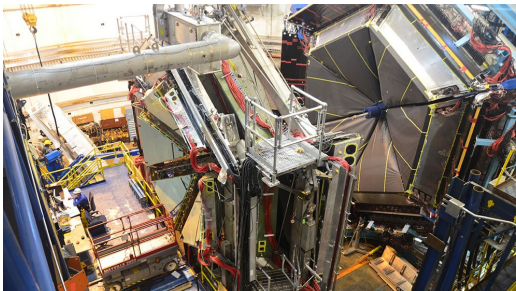
# Simulated CLAS12 Events - Silicon Vertex Tracker (SVT)



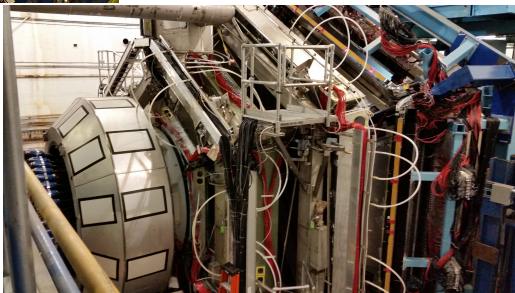
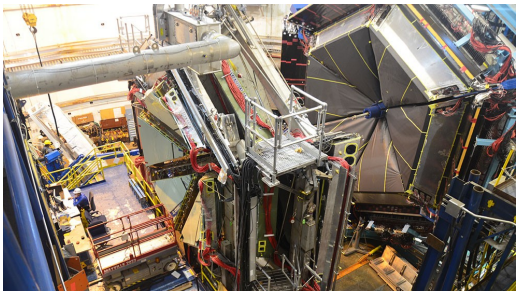
# Putting It All Together - 1



## Putting It All Together - 2



## Putting It All Together - 2





# Putting It All Together - 3

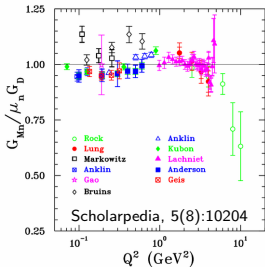
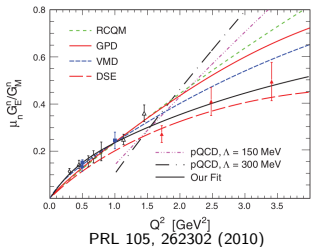


# Putting It All Together - 3

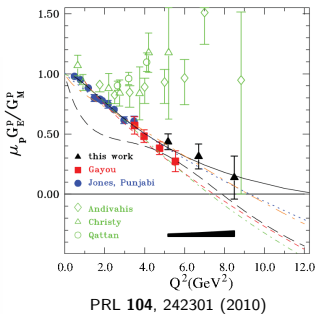
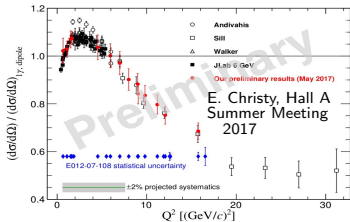


# Where We Are Now.

- $G_M^p$  well known over large  $Q^2$  range.
- The ratio  $G_E^p/G_M^p$  from polarization transfer measurements diverged from previous Rosenbluth separations.
  - Two-photon exchange (TPE).
  - Effect of radiative corrections.
- Neutron magnetic FF  $G_M^n$  still follows dipole.
- High- $Q^2$   $G_E^n$  opens up flavor decomposition



JLab E012-07-108,  $e-p$  elastic cross section



# How Do We Measure $G_M^n$ on a Neutron? (Step 3)

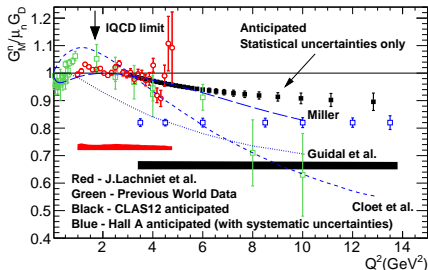
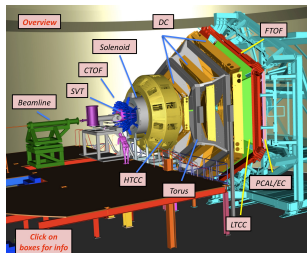
- E12-07-104 in Hall B (Gilfoyle, Hafidi, Brooks).
- Ratio Method on Deuterium:

$$R = \frac{\frac{d\sigma}{d\Omega} [{}^2\text{H}(e, e' n)_{QE}]}{\frac{d\sigma}{d\Omega} [{}^2\text{H}(e, e' p)_{QE}]}$$

$$= a \times \frac{\sigma_{\text{Mott}} \left( \frac{(G_E^n)^2 + \tau (G_M^n)^2}{1 + \tau} + 2\tau \tan^2 \frac{\theta_e}{2} (G_M^n)^2 \right)}{\frac{d\sigma}{d\Omega} [{}^1\text{H}(e, e' p)]}$$

where  $a$  is nuclear correction.

- Precise neutron detection efficiency needed to keep systematics low.
  - tagged neutrons from  ${}^2\text{H}(e, e' pn)$ .
  - $LH_2$  target.
- Kinematics:  $Q^2 = 3.5 - 13.0 \text{ (GeV}/c)^2$ .
- Beamtime: 40 days.
- Systematic uncertainties  $< 2.5\%$  across full  $Q^2$  range.
- Half of Run Group B done January, 2020.



# How Do We Measure $G_M^n$ on a Neutron? (Step 3)

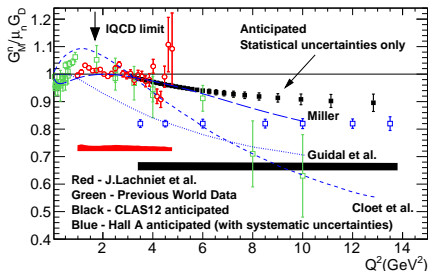
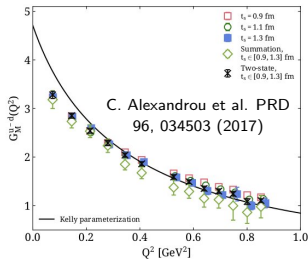
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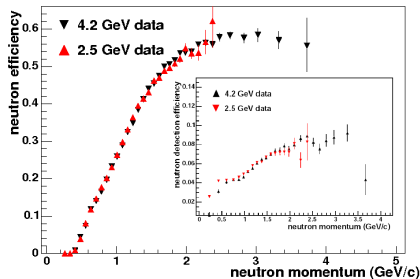
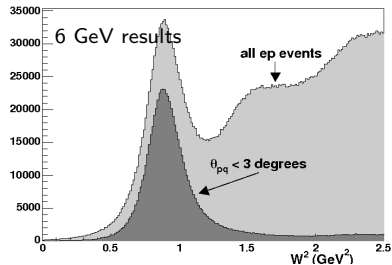
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- Systematic uncertainties  $< 2.5\%$  across full  $Q^2$  range.
- Half of Run Group B done January, 2020.



# How Do We Measure $G_M^n$ on a Neutron? (Step 4)

- Quasi-elastic event selection: Apply a maximum  $\theta_{pq}$  cut to eliminate inelastic events plus a cut on  $W^2$  (J.Lachniet thesis).
- Use the  $ep \rightarrow e'\pi^+n$  reaction from the hydrogen target as a source of tagged neutrons in the TOF and calorimeter.



# How Do We Measure $G_M^n$ on a Neutron? (Step 5)

Analyzing the data - CLAS12 computing requirements.

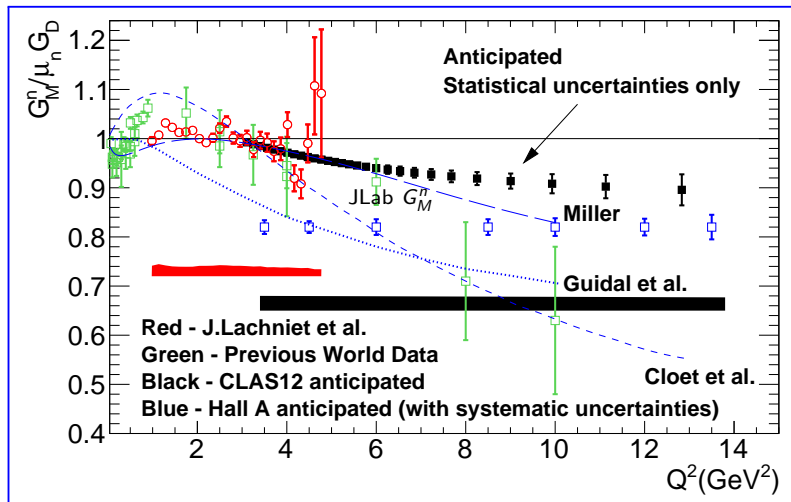
	Cores	Disk(TBytes)	Tape (TByte/year)
DAQ			1,270
Calibration	173		
Reconstruction	1,387	508	5,080
Simulation	8,139	318	1,558
Reconstruction Studies	1,214	508	
Physics Analysis	607	889	
Sum	11,520	2,223	7,938

We'll collect 5-10 TByte/day!

Intel Many-Integrated  
CoProcessor computer



# Anticipated Results





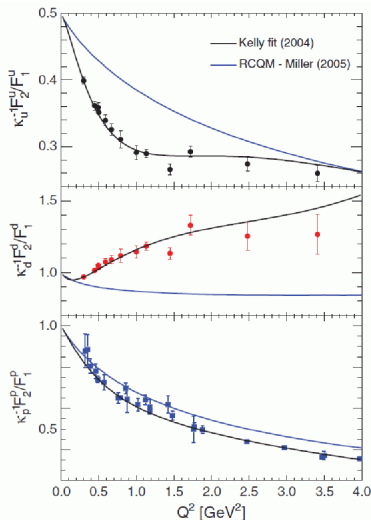
# Nuclear Structure - Flavor Decomposition

- By measuring all four EEFs we have an opportunity to unravel the contributions of the  $u$  and  $d$  quarks.
- Assume charge symmetry, no  $s$  quarks and use (Miller *et al.* Phys. Rep. **194**, 1 (1990))

$$F_{1(2)}^u = 2F_{1(2)}^p + F_{1(2)}^n$$

$$F_{1(2)}^d = 2F_{1(2)}^n + F_{1(2)}^p$$

- $u$  and  $d$  are different.
- AND different from the proton and neutron form factors.
- Evidence of di-quarks,  $s$  quark influence, ...?



Gordon Cates, Sean Riordan *et al.*, PRL **106**, 252003 (2011).

## Concluding Remarks

- JLab is a laboratory to test and expand our understanding of quark and nuclear matter, QCD, and the Standard Model.
- We continue the quest to unravel the nature of matter at greater and greater depths.
- Lots of new and exciting results are coming out.
- A bright future lies ahead in the 12 GeV Era.

U. S. Department of Energy's

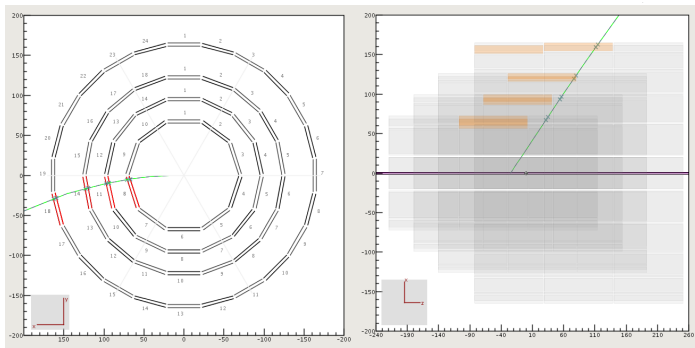
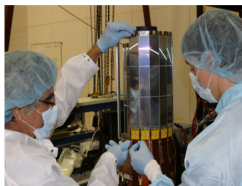
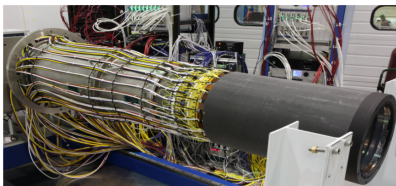


THOMAS JEFFERSON NATIONAL ACCELERATOR FACILITY

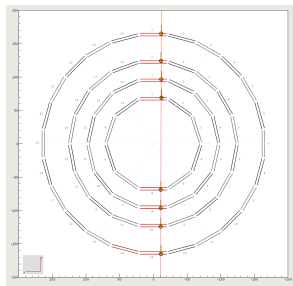
# Additional Slides

# What's going on now?

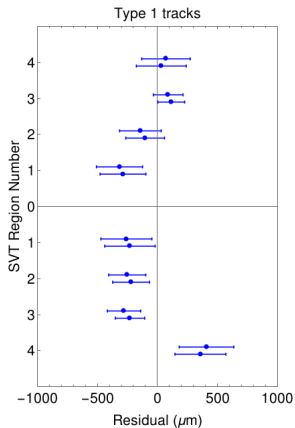
Alignment and commissioning of the silicon vertex tracker (SVT).



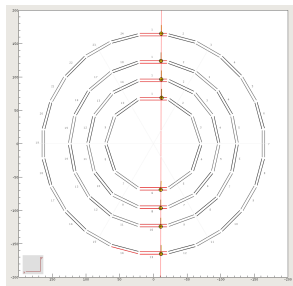
# Check alignment with Type1 cosmic ray tracks



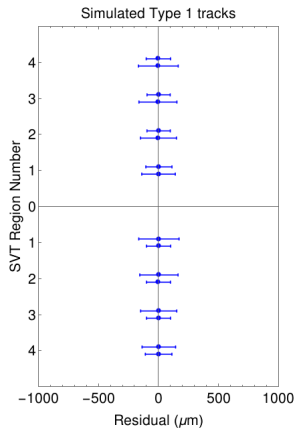
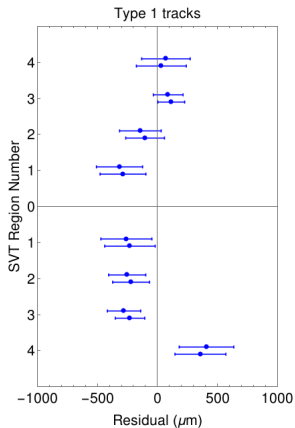
Type 1 tracks.

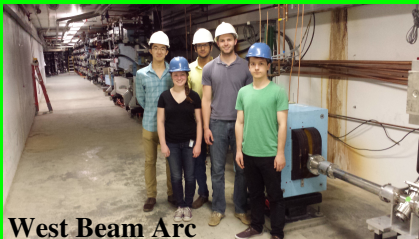


# Check alignment with Type1 cosmic ray tracks



Type 1 tracks.

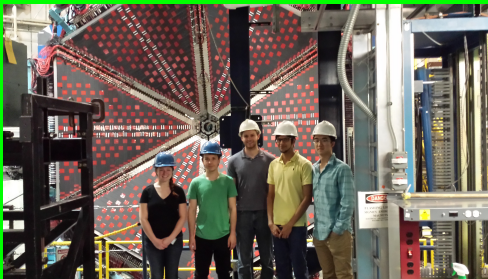




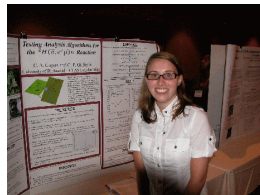
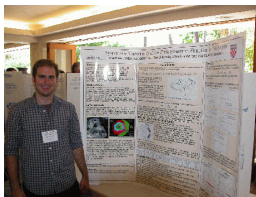
**West Beam Arc**

**Students from  
Richmond (and one  
from Surrey) visit  
JLab**

**CLAS12 detector**

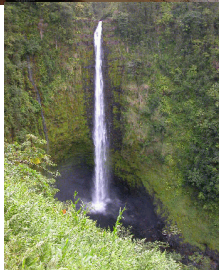
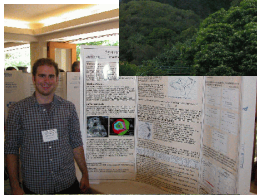
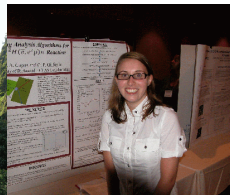


# Life on the Frontiers of Knowledge





# Life on the Frontiers of Knowledge



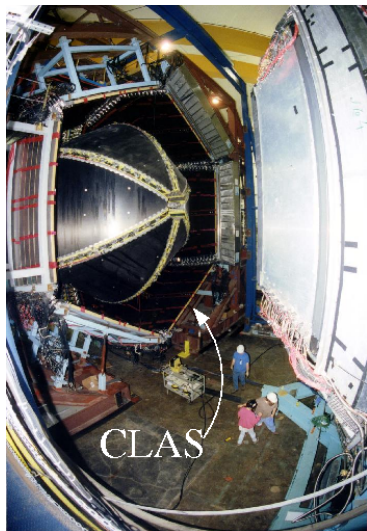
# Some Facts of Life On The Frontier

- Work at Jefferson Lab in Newport News.
  - 700 physicists, engineers, technicians, and staff.
  - Vibrant intellectual environment - talks, visitors, educational programs...
  - Lots going on.
- Richmond group part of CLAS Collaboration.
  - operates CLAS12.
  - ~190 physicists, 40 institutions, 13 countries.
  - Part of Software Group - emphasis on software development.
  - Past Surrey masters students (and Richmond undergrads) have presented posters at meetings, appeared on JLab publications,....
- Run-Group B consists of seven experiments (including  $G_M^n$ ) and ran in spring 2019.



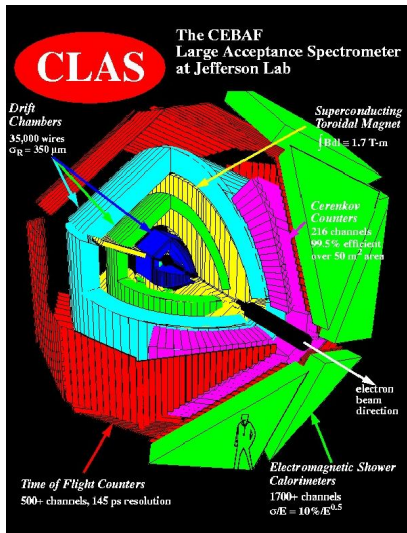
## How Do We Measure $G_M^n$ on a Neutron? (Step 2)

- Add one 45-ton, \$50-million radiation detector: the CEBAF Large Acceptance Spectrometer (CLAS).
- CLAS covers a large fraction of the total solid angle.
- Has about 35,000 detecting elements in about 40 layers.

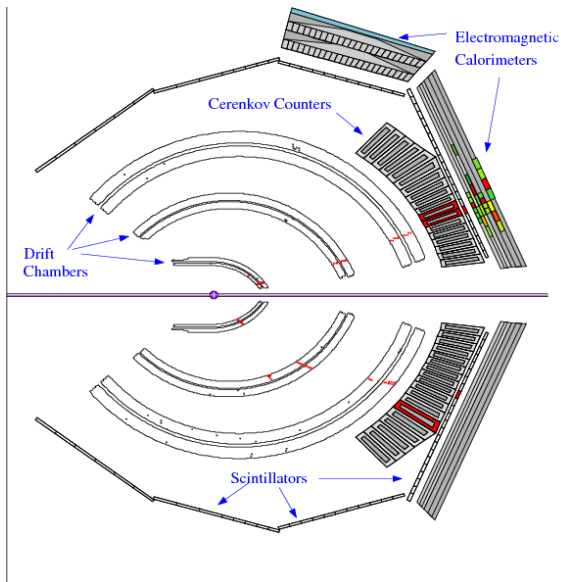


# How Do We Measure $G_M^n$ on a Neutron? (Step 2a)

- Drift chambers map the trajectories. A toroidal magnetic field bends the particles to measure momentum.
- Other layers measure energy, time-of-flight, and particle identification.
- Each collision is reconstructed and the intensity pattern reveals the forces and structure of the colliding particles.



# A CLAS Event



## How Do We Measure $G_M^n$ on a Neutron? (Step 3)

- Where's my target?

Use a dual target cell with liquid hydrogen and deuterium.

- How bad do the protons mess things up? They help!

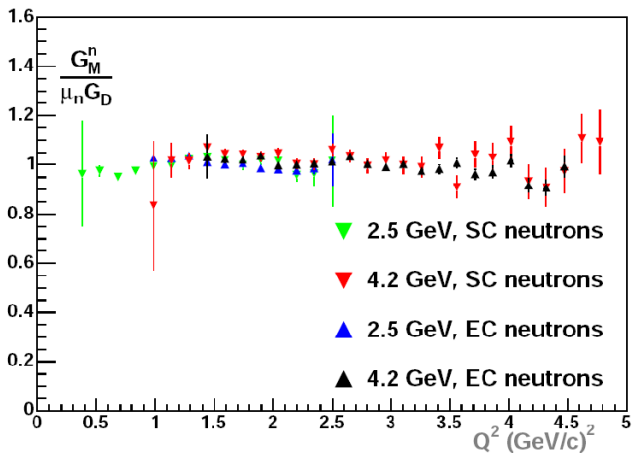


$$R = \frac{\frac{d\sigma}{d\Omega}(D(e, e'n))}{\frac{d\sigma}{d\Omega}(D(e, e'p))} = a(Q^2) \frac{\frac{G_E^{n^2} + \tau G_M^{n^2}}{1 + \tau} + 2\tau G_M^{n^2} \tan^2(\frac{\theta}{2})}{\frac{G_E^{p^2} + \tau G_M^{p^2}}{1 + \tau} + 2\tau G_M^{p^2} \tan^2(\frac{\theta}{2})}$$

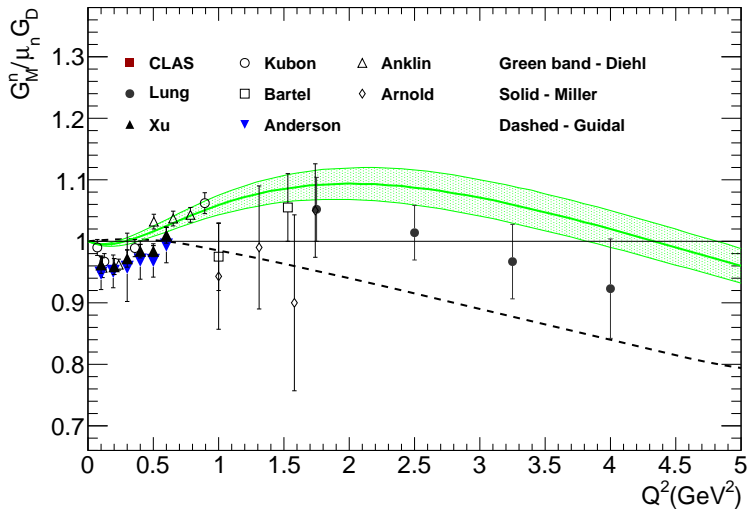
- The ratio is less vulnerable to corrections like acceptance, efficiencies, *etc.*
- Use the dual target to perform *in situ* detector calibrations.

# Results - Overlaps and Final Average

Overlapping measurements of  $G_M^n$  scaled by the dipole are consistent.

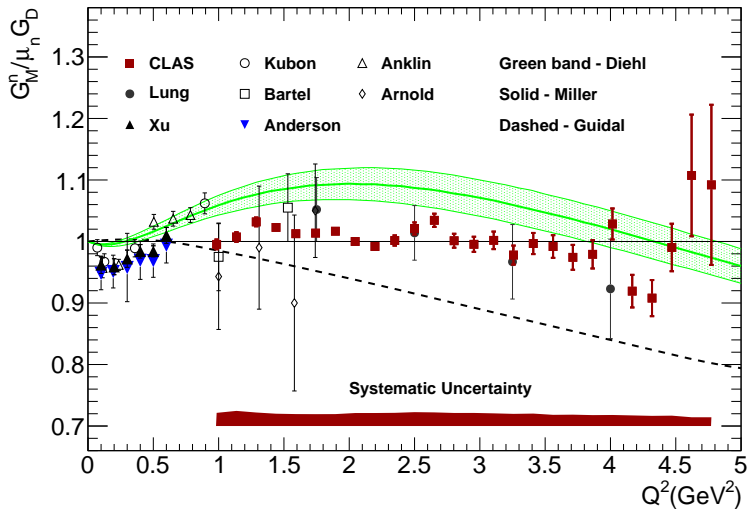


# Results - Comparison with Existing Data





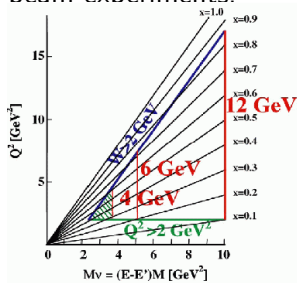
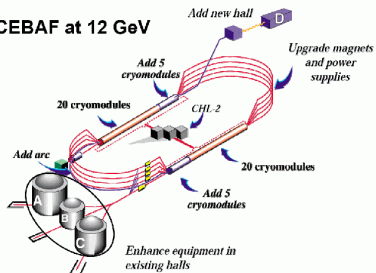
# Results - Comparison with Existing Data



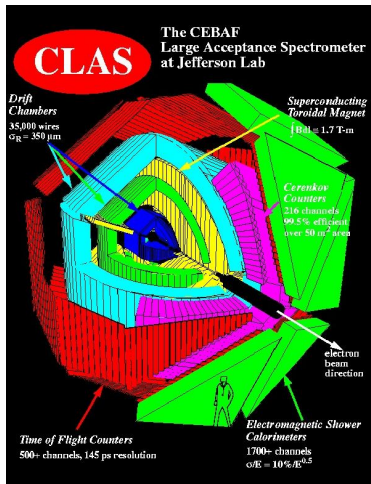
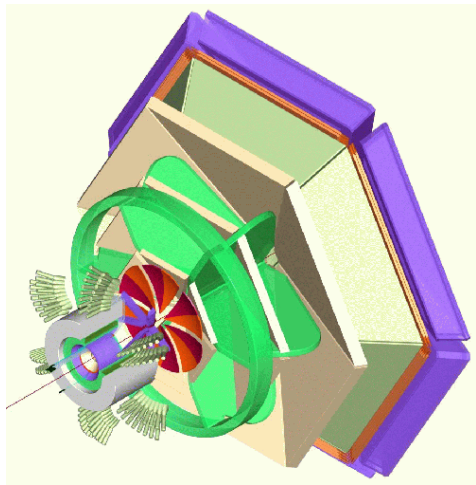
# JLab 12 GeV Upgrade - Better Accelerator

- The electron beam energy at JLab (CEBAF) has been doubled from 6 GeV to 12 GeV.
- Halls A, B and C will be upgraded to accommodate the new physics opportunities.
- A new hall (Hall D) will house a large-acceptance detector built around a solenoidal magnet for photon beam experiments.

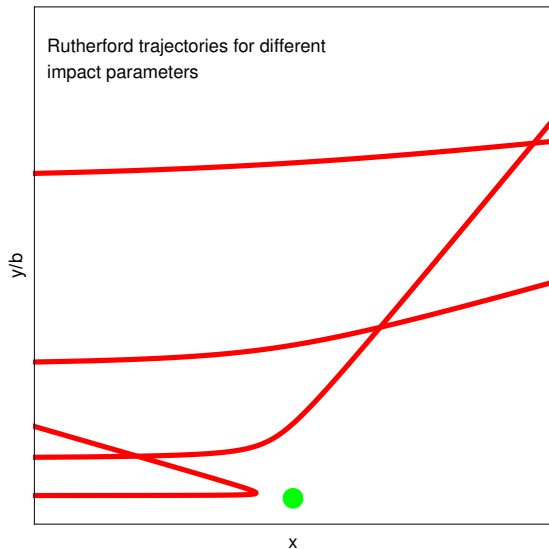
## CEBAF at 12 GeV



# JLab 12 GeV Upgrade - New Detectors

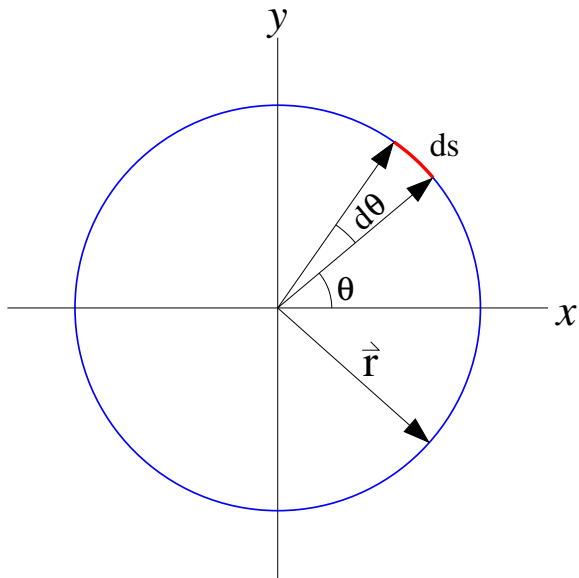


# Rutherford Trajectories



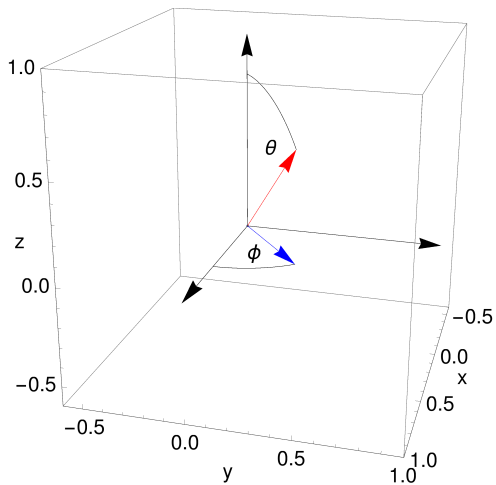
# What is an Angle?

# What is an Angle?

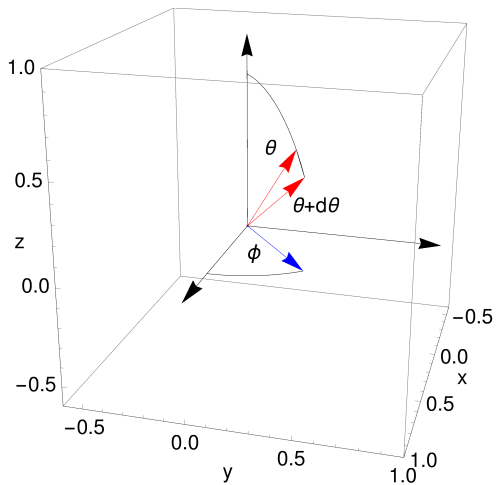


$$d\theta = \frac{ds}{|\vec{r}|}$$

# Solid Angle

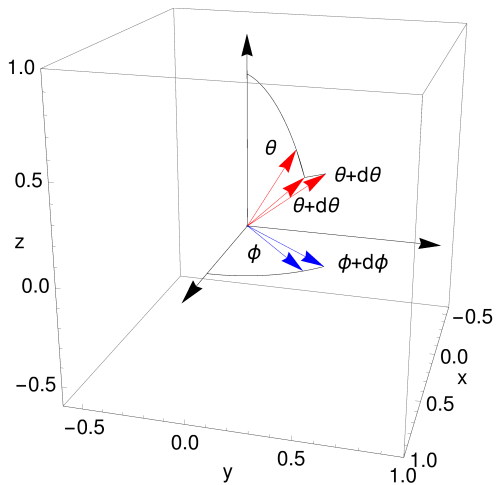


# Solid Angle

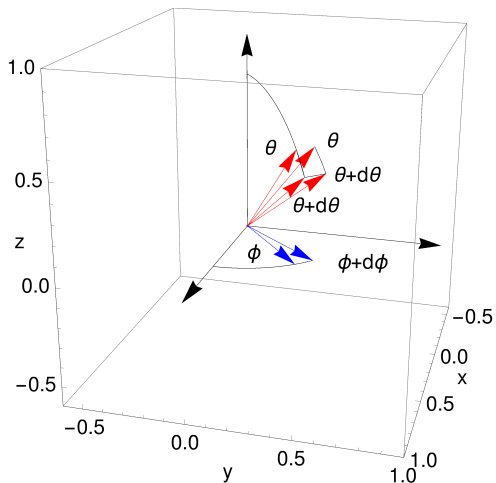




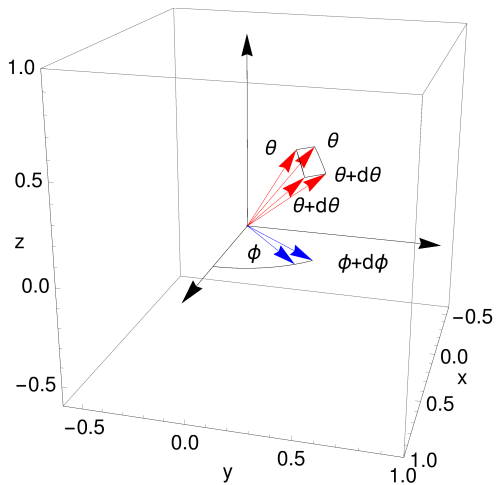
# Solid Angle



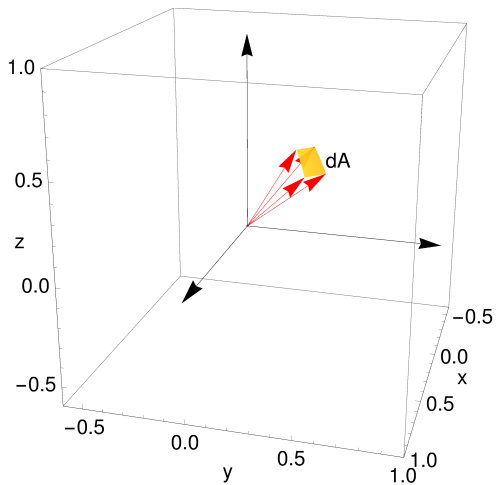
# Solid Angle



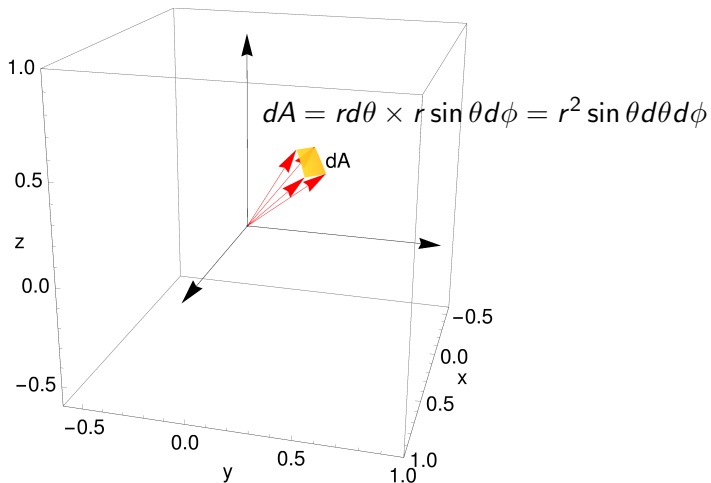
# Solid Angle



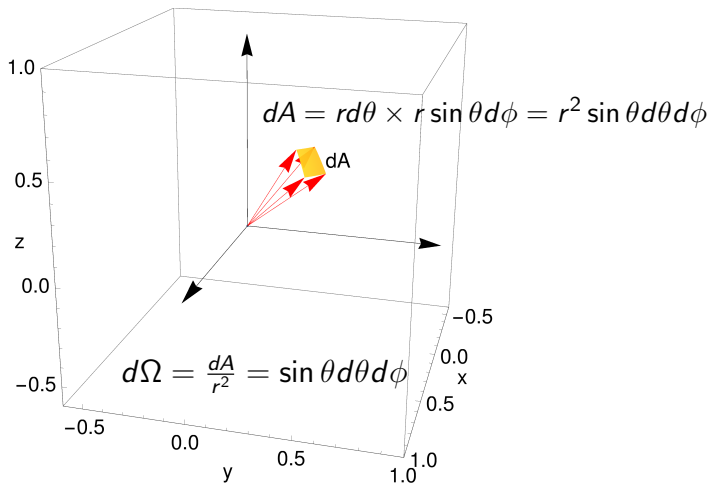
# Solid Angle



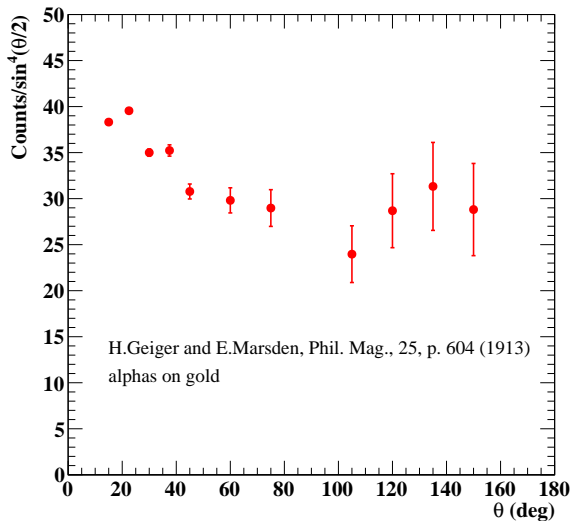
# Solid Angle



# Solid Angle

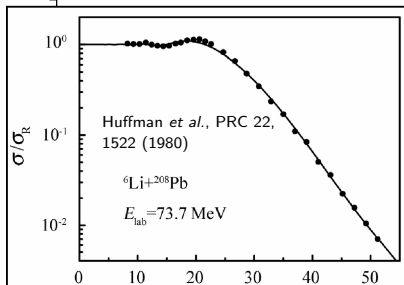
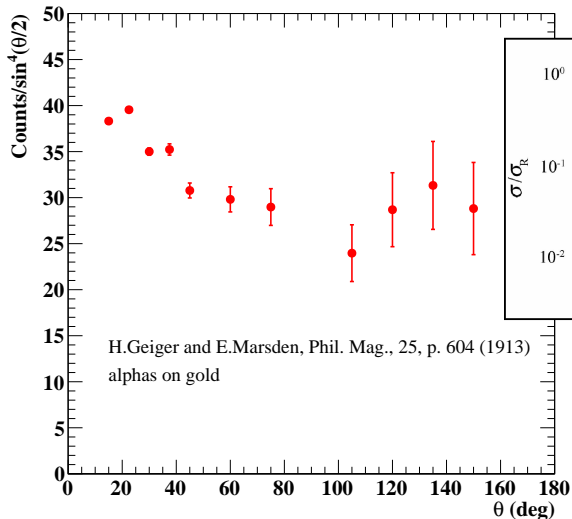


# Rutherford Scattering Results From Rutherford



2016-12-18 10:38:14

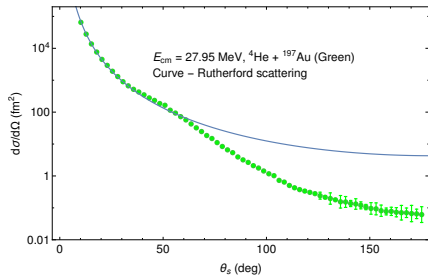
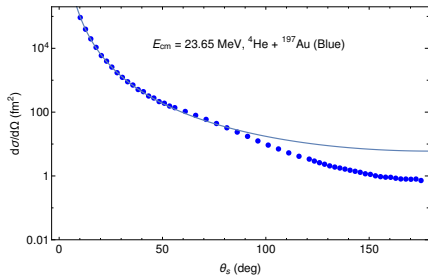
# Rutherford Scattering Results From Rutherford



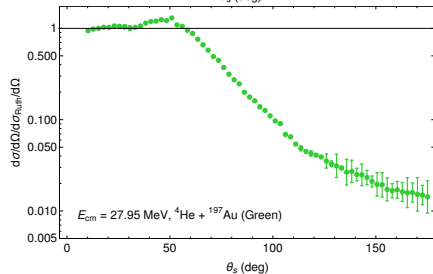
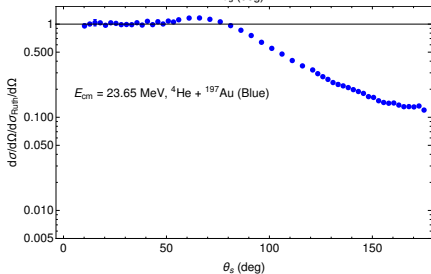
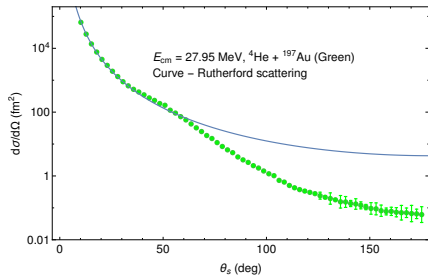
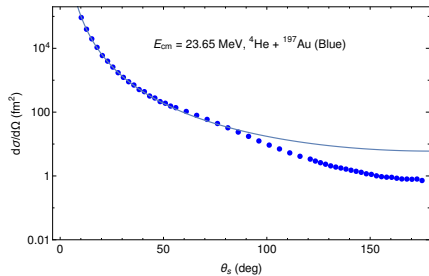
2016-12-18 10:38:14



# Recent Rutherford Scattering Results



# Recent Rutherford Scattering Results



# Standard Model

