

Hunting for Quarks

Jerry Gilfoyle for the CLAS Collaboration
University of Richmond



"The Periodic Table"

- JLab Mission
- What we know and don't know.
- The Neutron Magnetic Form Factor
- Experiments with CLAS
- More JLab Highlights
- Concluding Remarks

G_M^n Co-conspirators:

Jeff Lachniet

Will Brooks

Mike Vineyard

Brian Quinn

Kristen Greenholt (UG)

Rusty Burrell (UG)

What is the Mission of Jefferson Lab?

- Pursue basic research into the quark nature of the atomic nucleus..
- Map the geography of the transition from proton-neutron picture of nuclei to one based on quarks and gluons.
- Provide a testing ground for the theory of the color force Quantum Chromodynamics (QCD) and the nature of quark confinement.
- Probe the quark-gluon structure of hadronic matter and how it evolves within nuclei.



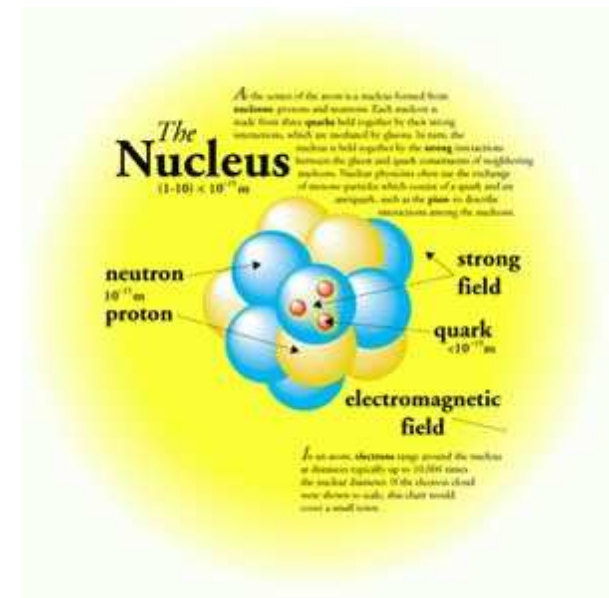
What Do We Know?

- The Universe is made of quarks and leptons and the force carriers.

BOSONS			force carriers spin = 0, 1, 2, ...		
Unified Electroweak spin = 1			Strong (color) spin = 1		
Name	Mass GeV/c ²	Electric charge	Name	Mass GeV/c ²	Electric charge
γ photon	0	0	g gluon	0	0
W⁻	80.4	-1			
W⁺	80.4	+1			
Z⁰	91.187	0			

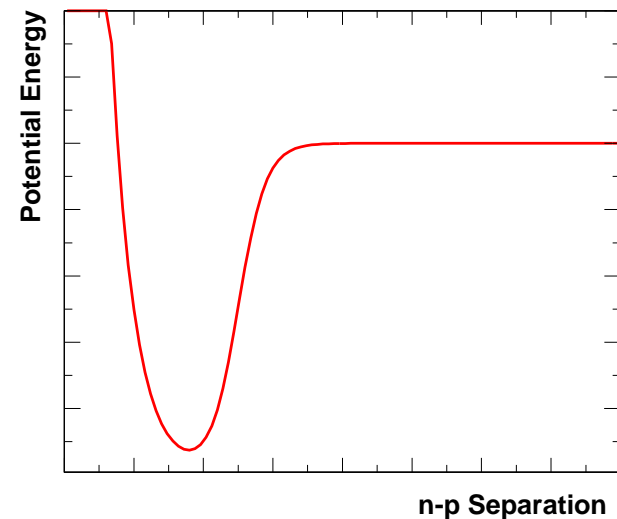
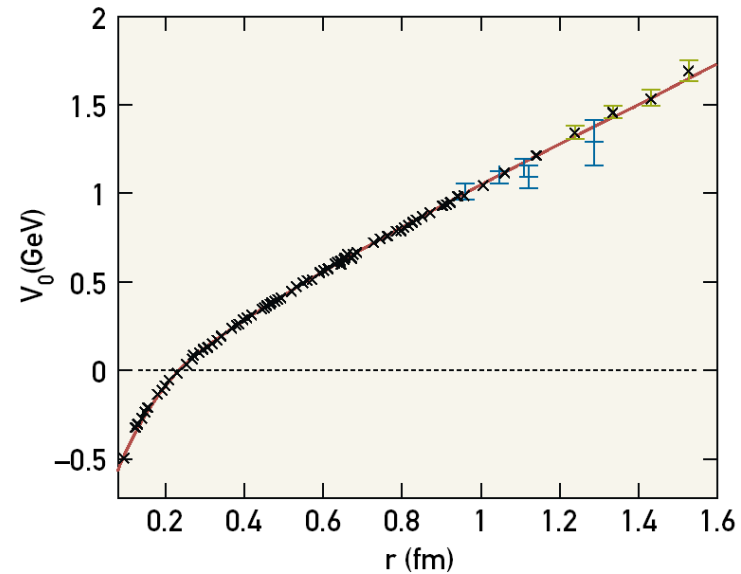
FERMIONS			matter constituents spin = 1/2, 3/2, 5/2, ...		
Leptons spin = 1/2			Quarks spin = 1/2		
Flavor	Mass GeV/c ²	Electric charge	Flavor	Approx. Mass GeV/c ²	Electric charge
ν_e electron neutrino	$<1 \times 10^{-8}$	0	u up	0.003	2/3
e electron	0.000511	-1	d down	0.006	-1/3
ν_μ muon neutrino	<0.0002	0	c charm	1.3	2/3
μ muon	0.106	-1	s strange	0.1	-1/3
ν_τ tau neutrino	<0.02	0	t top	175	2/3
τ tau	1.7771	-1	b bottom	4.3	-1/3

- The atomic nucleus is made of protons and neutrons bound by the strong force.
- The quarks are confined inside the protons and neutrons.
- Protons and neutrons are NOT confined.



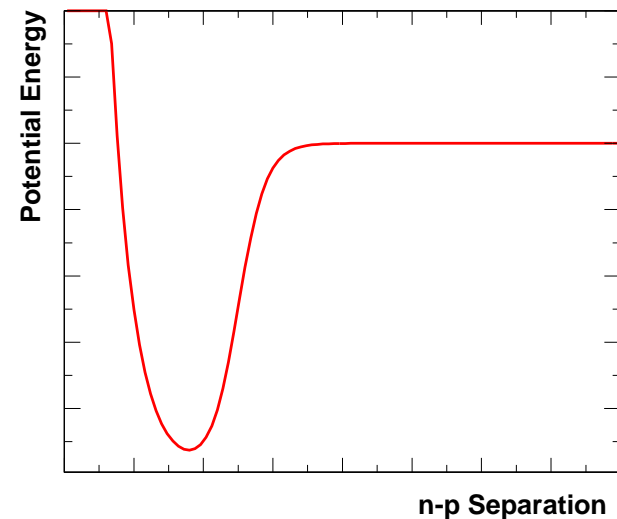
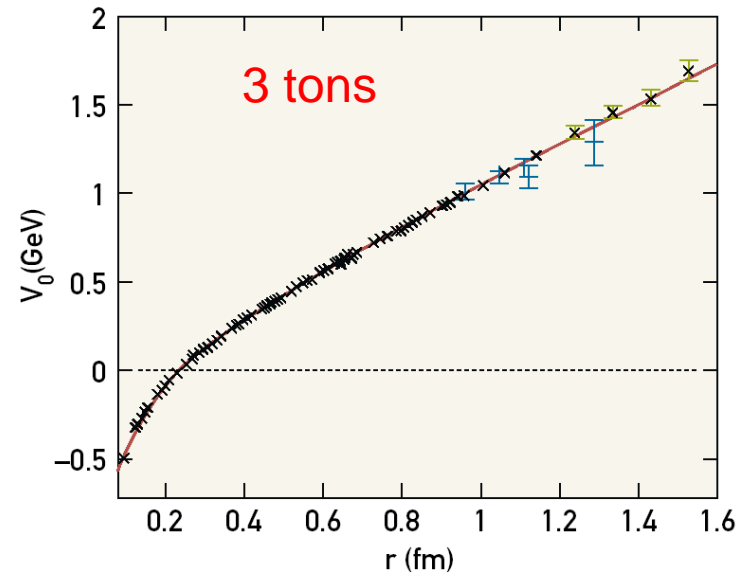
What is the Force?

- Quantum chromodynamics (QCD) looks like the right way to get the force at high energy.
- The hadronic model uses a phenomenological force fitted to data at low energy. This 'strong' force is the residual force between quarks.



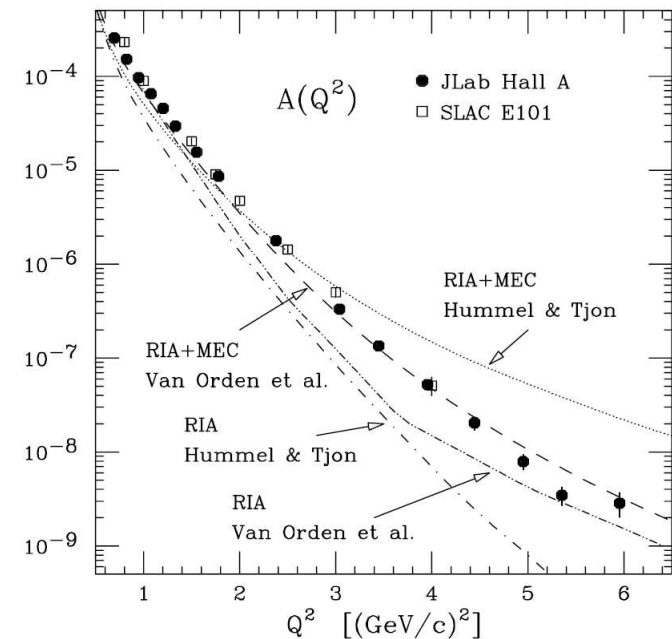
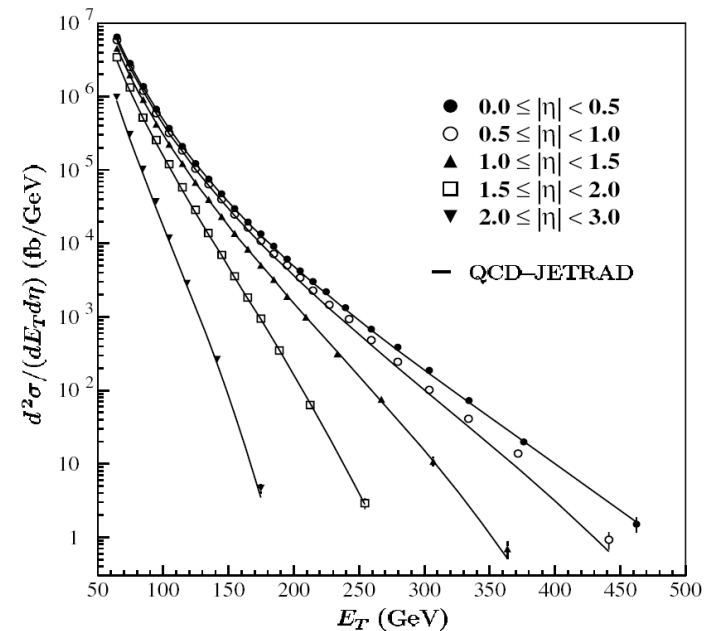
What is the Force?

- Quantum chromodynamics (QCD) looks like the right way to get the force at high energy.
- The hadronic model uses a phenomenological force fitted to data at low energy. This 'strong' force is the residual force between quarks.



How Well Do We Know It?

- We have a working theory of strong interactions: quantum chromodynamics or QCD (B.Abbott, *et al.*, Phys. Rev. Lett., **86**, 1707 (2001)).
- The coherent hadronic model (the standard model of nuclear physics) works too (L.C.Alexa, *et al.*, Phys. Rev. Lett., **82**, 1374 (1999)).



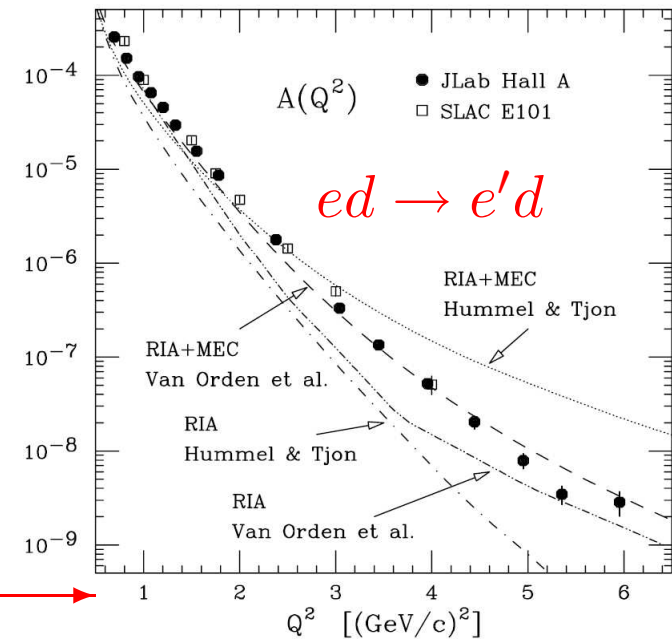
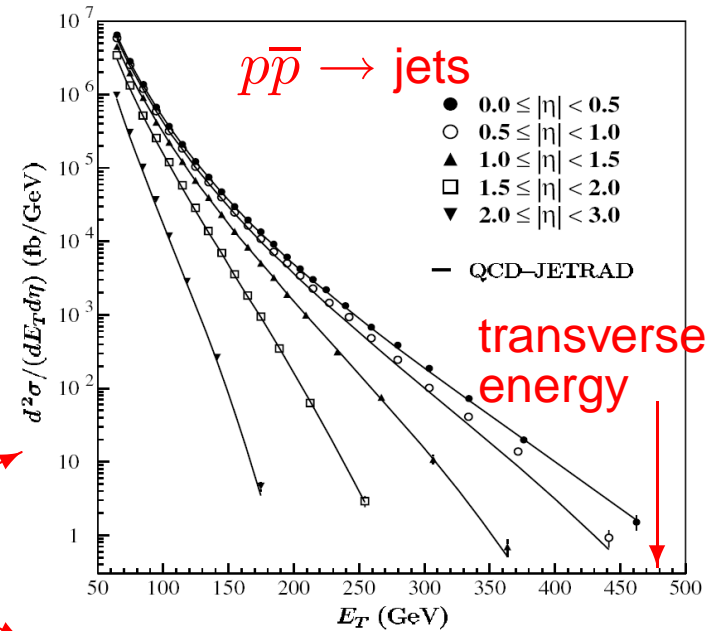
How Well Do We Know It?

- We have a working theory of strong interactions: quantum chromodynamics or QCD (B.Abbott, *et al.*, Phys. Rev. Lett., **86**, 1707 (2001)).

effective target area

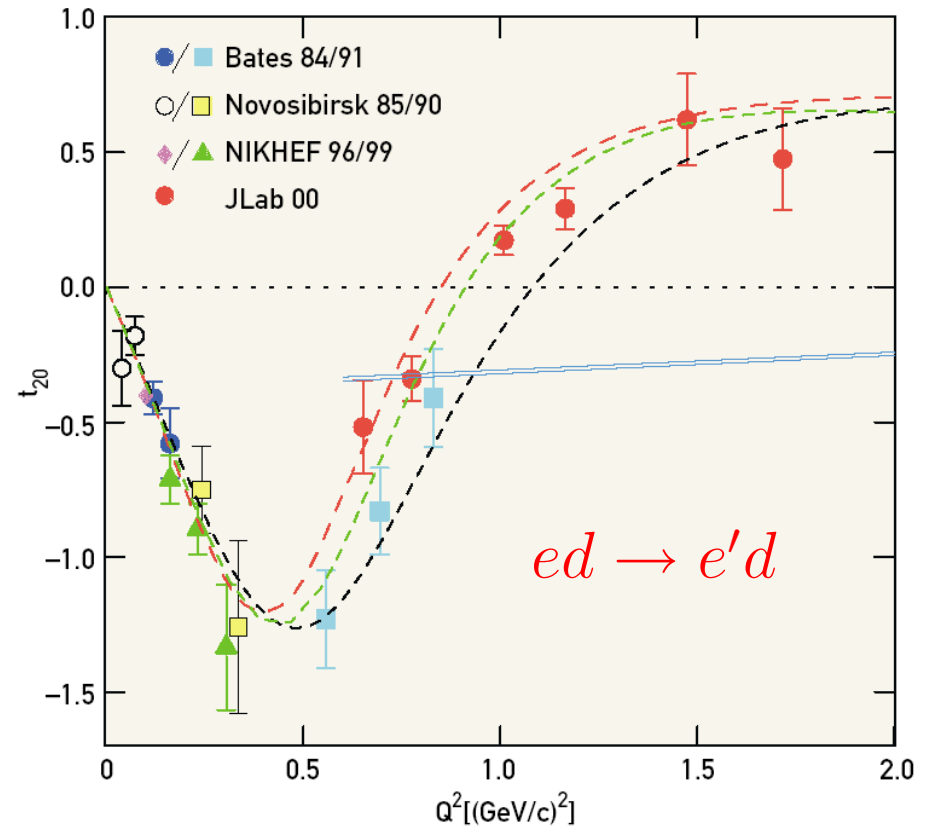
- The coherent hadronic model (the standard model of nuclear physics) works too (L.C.Alexa, *et al.*, Phys. Rev. Lett., **82**, 1374 (1999)).

4-momentum transfer squared



What Don't We Know?

1. We can't get QCD and the hadronic model to line up.
D. Abbott, *et al.*, Phys. Rev Lett. **84**, 5053 (2000).
2. NEED TO FIGURE OUT QCD AT THE ENERGIES OF NUCLEI!!



The Magnetic Form Factor of the Neutron (G_M^n)

- Fundamental quantity related to the distribution of charge and magnetization/currents in the proton and neutron (the nucleons).
- Part of a broad effort to understand how nucleons are 'constructed from the quarks and gluons of QCD'.*
- Needed to extract the distribution of quarks in the neutron.
- Elastic form factors (G_M^n , G_E^n , G_M^p , and G_E^p) provide key constraints on generalized parton distributions (GPDs) which promise to give us a three-dimensional picture of hadrons.
- Fundamental challenge for lattice QCD.

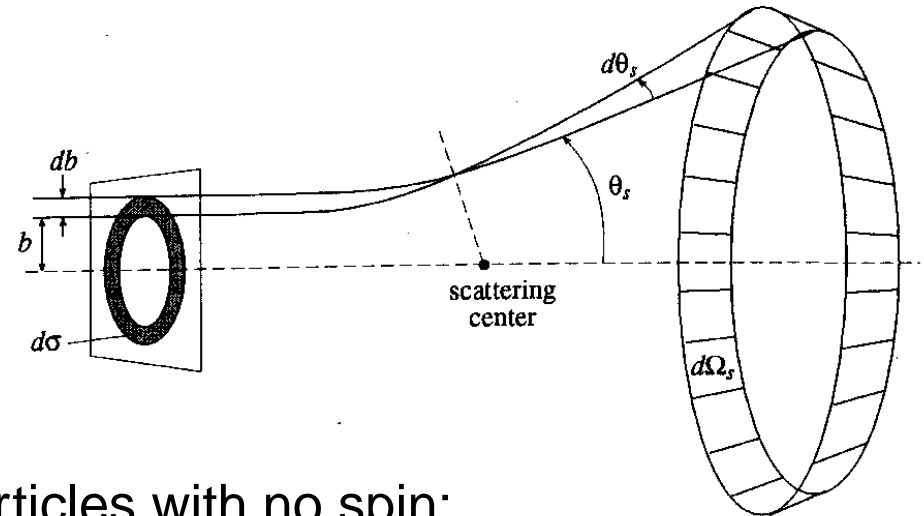
* 'Opportunities in Nuclear Science: A Long-Range Plan for the Next Decade', NSF/DOE

Nuclear Science Advisory Committee, April, 2002.

What is a Form Factor?

- Start with the cross section.

$$\frac{d\sigma}{d\Omega} = \frac{\text{scattered rate/solid angle}}{\text{incident rate} \times \text{targets/area}}$$



For elastic scattering by point particles with no spin:

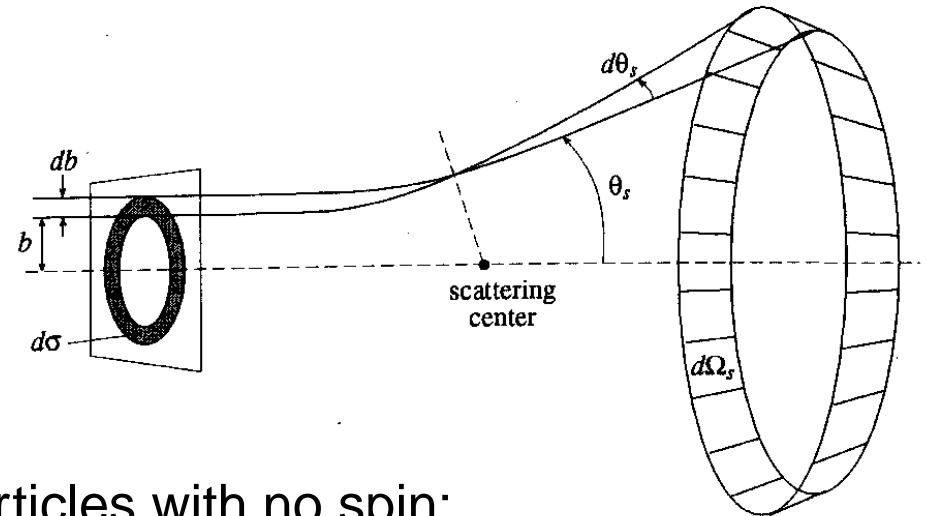
$$\frac{d\sigma}{d\Omega} = \frac{Z^2 \alpha^2 (\hbar c)^2}{16E^2 \sin^4(\theta/2)}$$

(Rutherford cross section, σ_{Ruth}).

What is a Form Factor?

- Start with the cross section.

$$\frac{d\sigma}{d\Omega} = \frac{\text{scattered rate/solid angle}}{\text{incident rate} \times \text{targets/area}}$$



For elastic scattering by point particles with no spin:

$$\frac{d\sigma}{d\Omega} = \frac{Z^2 \alpha^2 (\hbar c)^2}{16E^2 \sin^4(\theta/2)} \quad (\text{Rutherford cross section, } \sigma_{Ruth}).$$

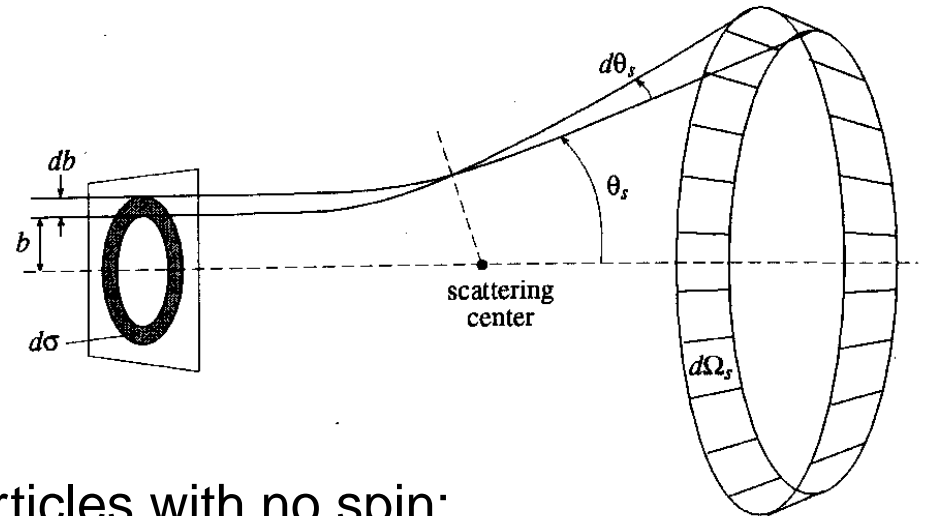
- Get the cross section for elastic scattering by point particles with spin.

$$\frac{d\sigma}{d\Omega} = \sigma_{Ruth} \left(1 - \beta^2 \sin^2 \frac{\theta}{2} \right) \quad (\text{Mott cross section})$$

What is a Form Factor?

- Start with the cross section.

$$\frac{d\sigma}{d\Omega} = \frac{\text{scattered rate/solid angle}}{\text{incident rate} \times \text{targets/area}}$$



For elastic scattering by point particles with no spin:

$$\frac{d\sigma}{d\Omega} = \frac{Z^2 \alpha^2 (\hbar c)^2}{16E^2 \sin^4(\theta/2)} \quad (\text{Rutherford cross section, } \sigma_{Ruth}).$$

- Get the cross section for elastic scattering by point particles with spin.

$$\frac{d\sigma}{d\Omega} = \sigma_{Ruth} \left(1 - \beta^2 \sin^2 \frac{\theta}{2} \right) \quad (\text{Mott cross section})$$

- What happens when the beam is electrons and the target is not a point?

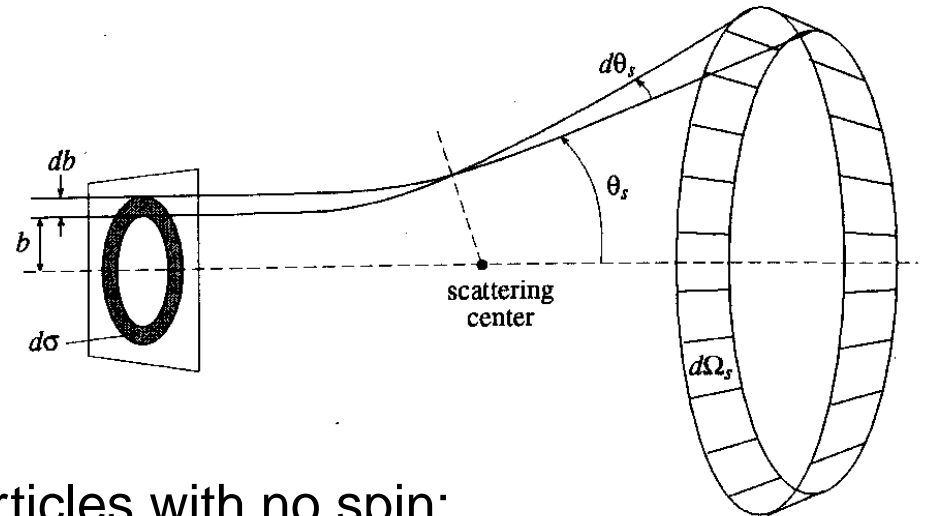
$$\frac{d\sigma}{d\Omega} = \sigma_{Ruth} \left(1 - \beta^2 \sin^2 \frac{\theta}{2} \right) |F(Q^2)|^2$$

where Q^2 is the 4-momentum transfer.

What is a Form Factor?

- Start with the cross section.

$$\frac{d\sigma}{d\Omega} = \frac{\text{scattered rate/solid angle}}{\text{incident rate} \times \text{targets/area}}$$



For elastic scattering by point particles with no spin:

$$\frac{d\sigma}{d\Omega} = \frac{Z^2 \alpha^2 (\hbar c)^2}{16 E^2 \sin^4(\theta/2)} \quad (\text{Rutherford cross section, } \sigma_{Ruth}).$$

- Get the cross section for elastic scattering by point particles with spin.

$$\frac{d\sigma}{d\Omega} = \sigma_{Ruth} \left(1 - \beta^2 \sin^2 \frac{\theta}{2} \right) \quad (\text{Mott cross section})$$

- What happens when the beam is electrons and the target is not a point?

$$\frac{d\sigma}{d\Omega} = \sigma_{Ruth} \left(1 - \beta^2 \sin^2 \frac{\theta}{2} \right) |F(Q^2)|^2$$

where Q^2 is the 4-momentum transfer.

THE FORM FACTOR!

Why Should You Care?

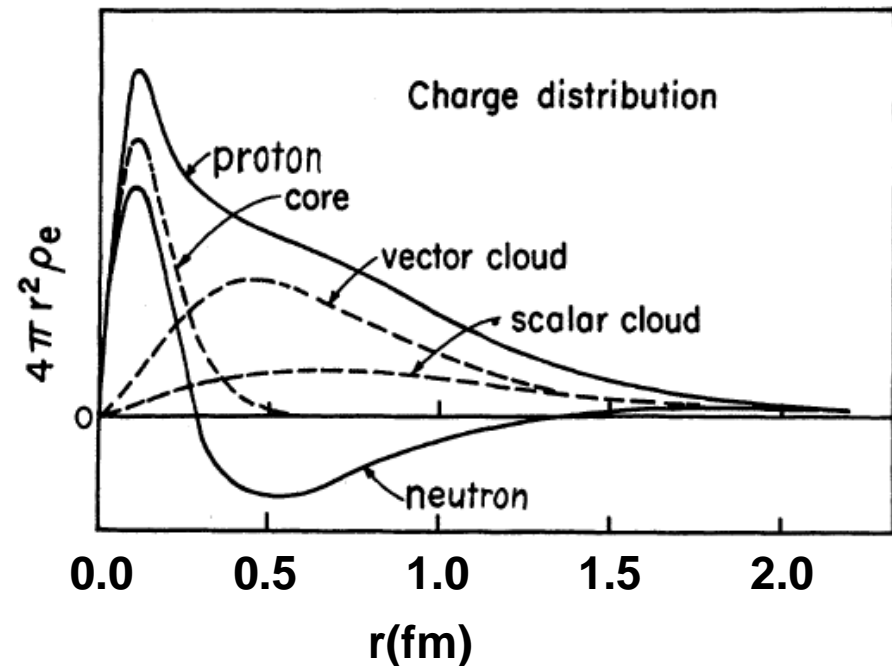
- The chain of reason.

$$\frac{d\sigma}{d\Omega} \rightarrow |F(Q^2)|^2 \Leftrightarrow F(Q^2) \leftarrow \rho(\vec{r}) \leftarrow \psi(\vec{r}) \leftarrow \begin{array}{l} \text{QCD,} \\ \text{Constituent quarks} \\ \text{Theory} \end{array}$$

Experiment Comparison

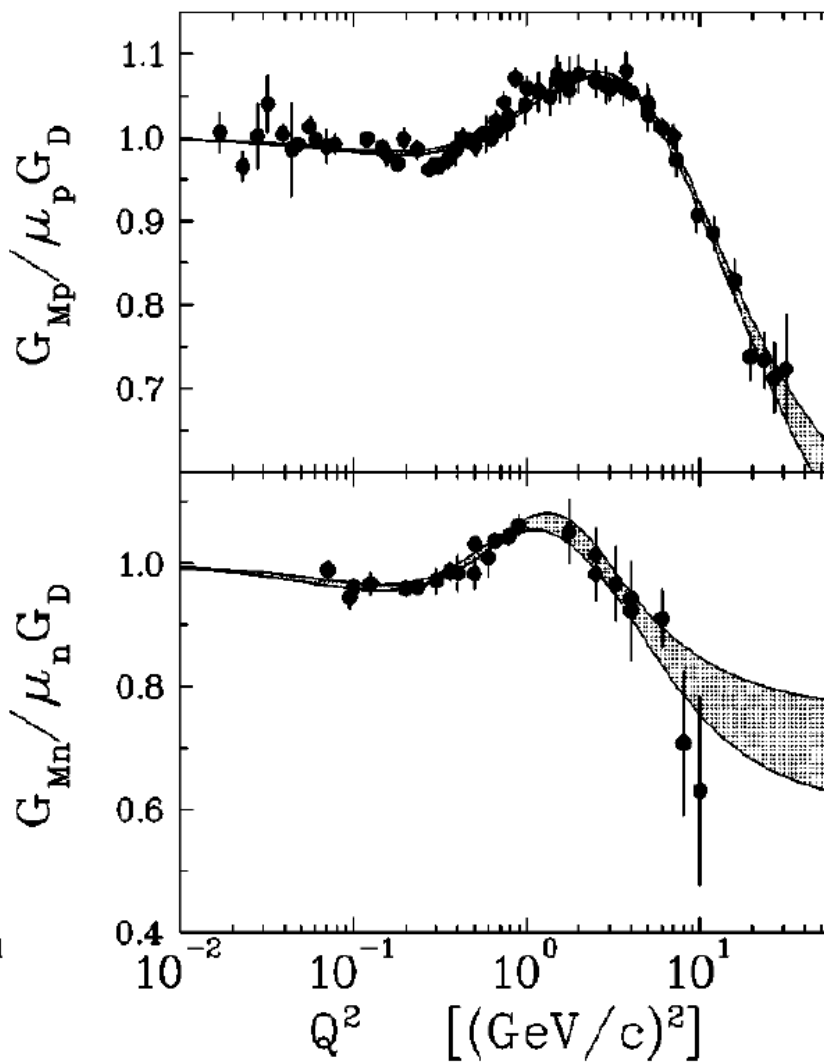
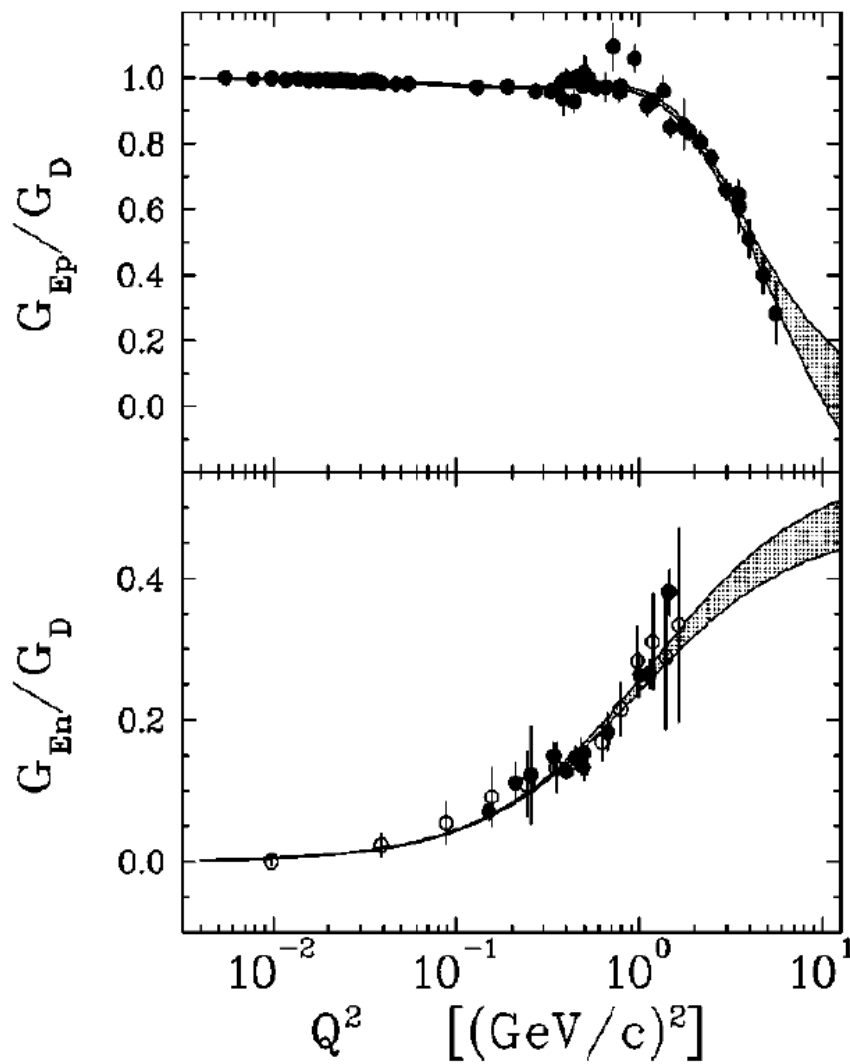
The form factors are the meeting ground between theory and experiment.

- The Fourier transform of the form factors are related to the charge and current distributions within the neutron. (R.M.Littauer, H.F.Schopper, R.R.Wilson, Phys. Rev. Lett., 7, 144 (1961).)



The World's Data.

From J.Kelly, Phys. Rev. C **70**, 068202 (2004).

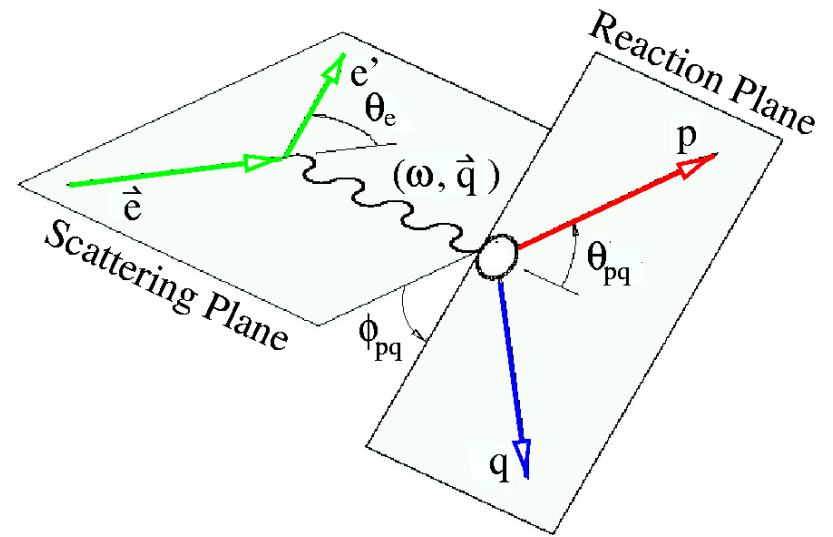


Why Should You Care Even More?

- Some kinematic definitions

$(\theta_{pq}, \phi_{pq}, b)$.

The coordinate b is the transverse distance from the z axis defined by \vec{q} .



- Recent analysis of existing form factor data by G. Miller (Phys. Rev. Lett. **99**, 112001 (2007)).

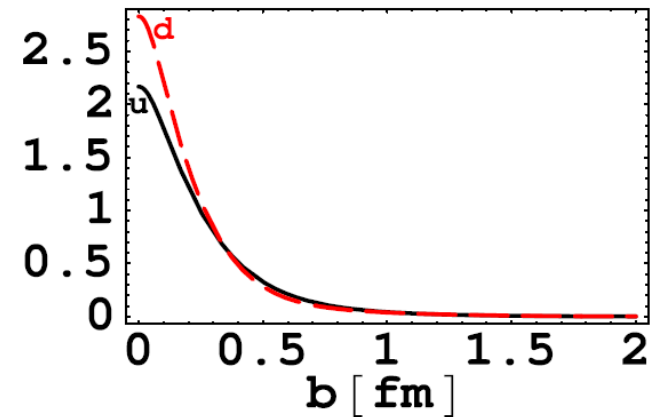
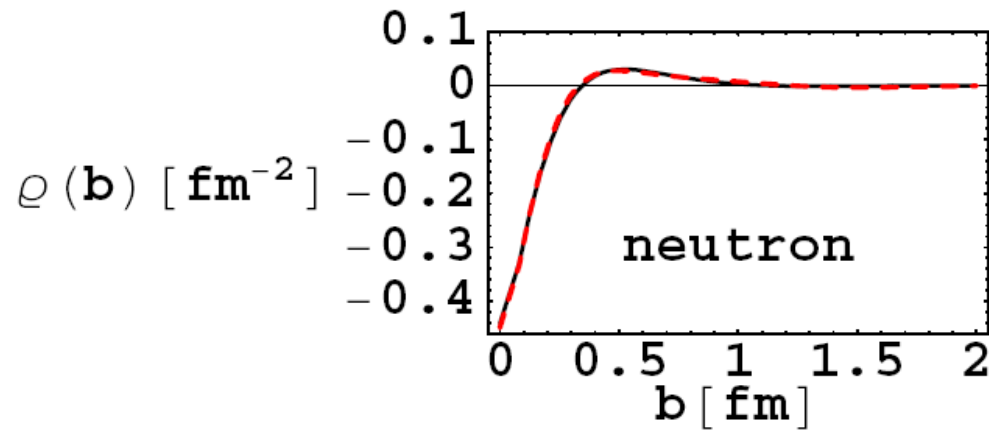
$$\rho(b) = \int_0^\infty dQ^2 \frac{1}{(2\pi)^2} \frac{G_E^n(Q^2) + \tau G_M^n(Q^2)}{1 + \tau} e^{i\vec{q}\cdot\vec{b}} \quad \tau = Q^2/4M_N^2$$

Why Should You Care Even More? - Part Deux

- Final form of $\rho(b)$.

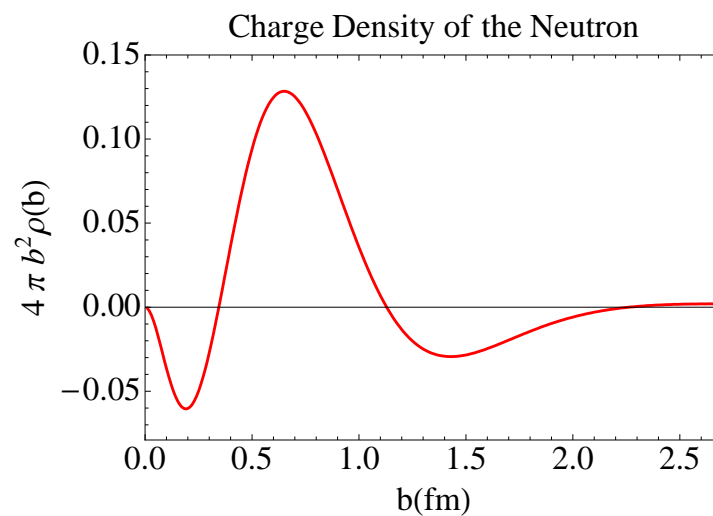
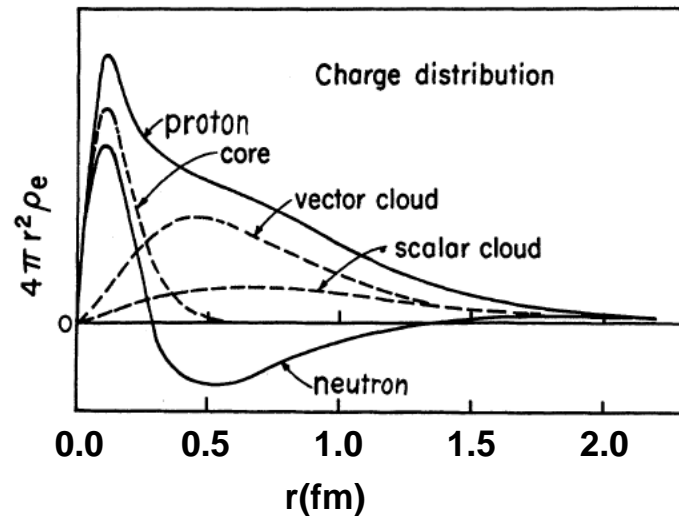
$$\rho(b) = \int_0^\infty dQ \frac{Q}{2\pi} J_0(Qb) \frac{G_E(Q^2) + \tau G_M Q^2}{1 + \tau} \quad \tau = Q^2 / 4M_N^2$$

- Miller results



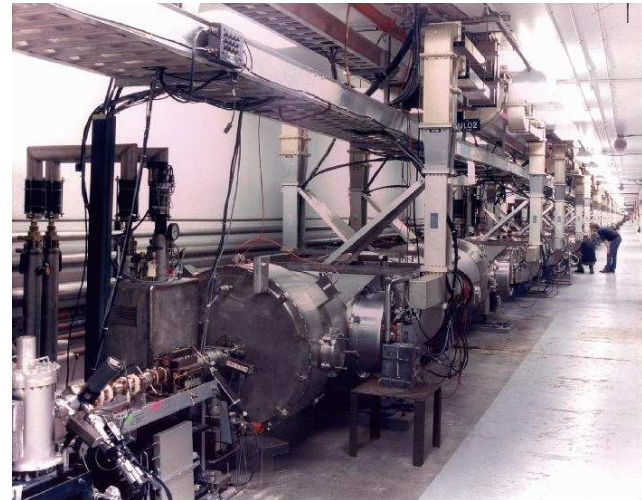
Why Should You Care Even More? - Part Three

- Comparison with previous results. Note that b and r are conceptually different.



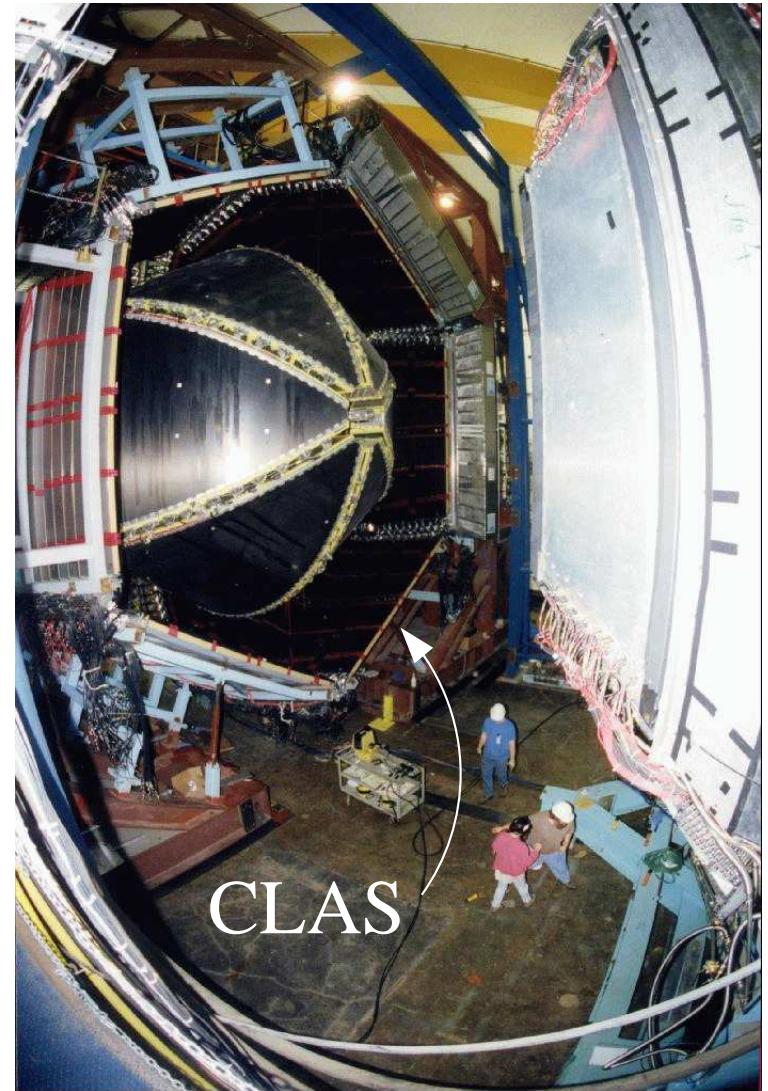
How Do We Measure G_M^n on a Neutron? (Step 1)

- Start at your local mile-long, high-precision, 6-GeV electron accelerator.
- The Continuous Electron Beam Accelerator Facility (CEBAF) produces beams of unrivaled quality.
- Electrons do up to five laps, are extracted, and sent to one of three experimental halls.
- All three halls can run simultaneously.



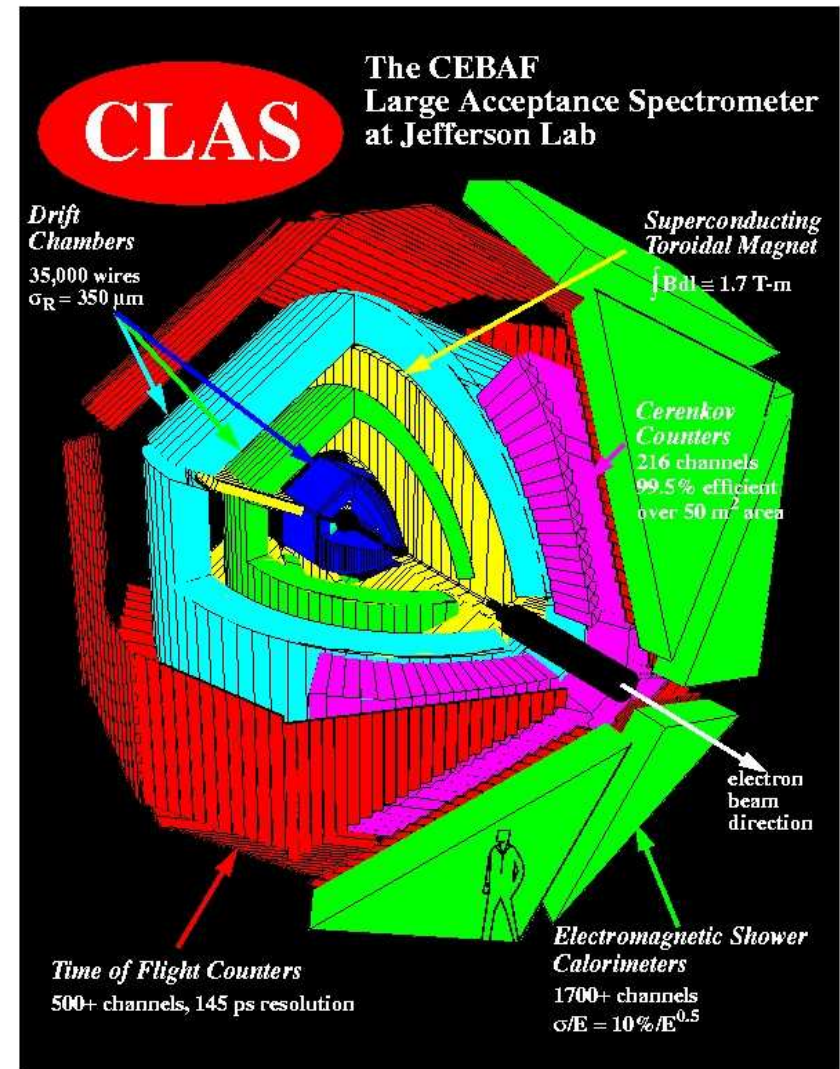
How Do We Measure G_M^n on a Neutron? (Step 2)

- Add one 45-ton, \$50-million radiation detector: the CEBAF large Acceptance Spectrometer (CLAS).
- CLAS covers a large fraction of the total solid angle.
- Has about 35,000 detecting elements in about 40 layers.

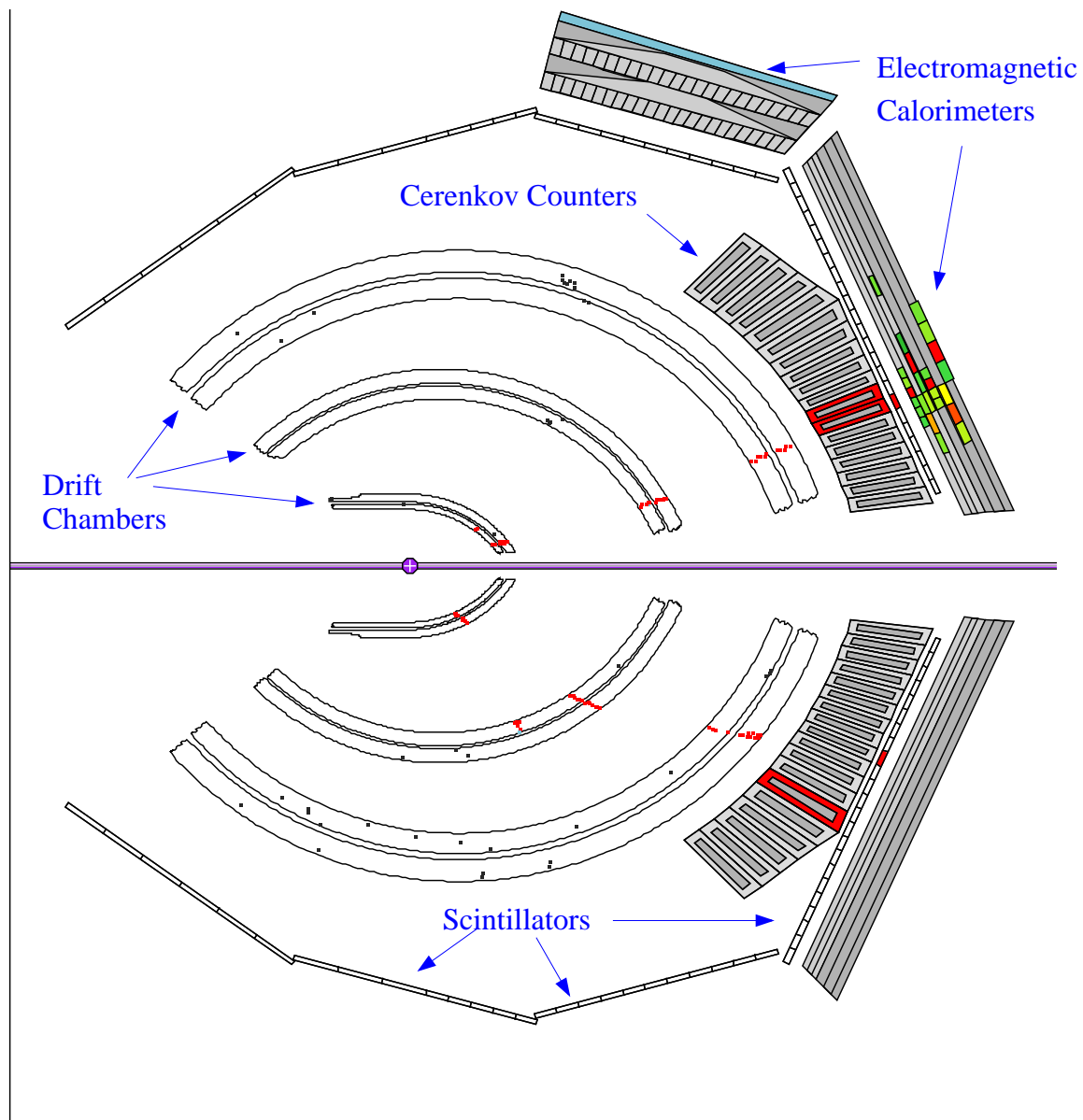


How Do We Measure G_M^n on a Neutron? (Step 2)

- Drift chambers map the trajectory of the collision. A toroidal magnetic field bends the trajectory to measure momentum.
- Other layers measure energy, time-of-flight, and particle identification.
- Each collision is reconstructed and the intensity pattern reveals the forces and structure of the colliding particles.



A CLAS Event



How Do We Measure G_M^n on a Neutron? (Step 3)

- Where's my target?

Use a dual target cell with liquid hydrogen and deuterium.

- How bad do the protons mess things up? They help!

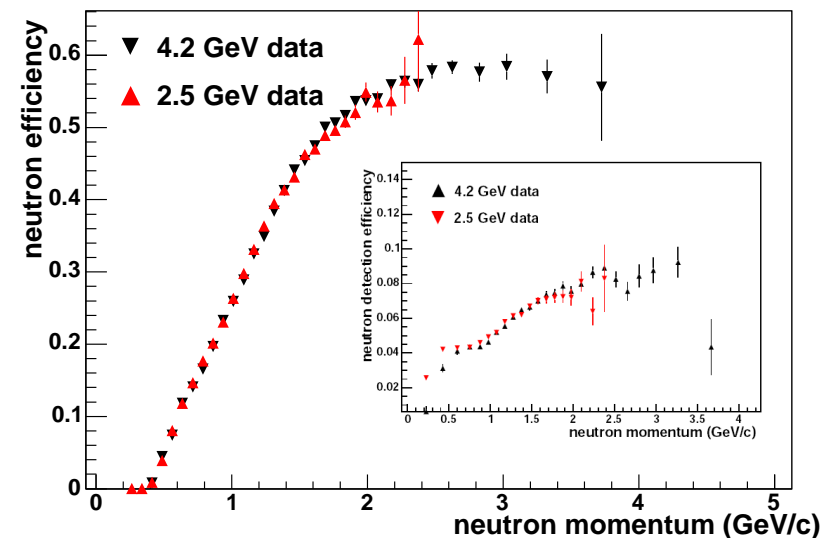
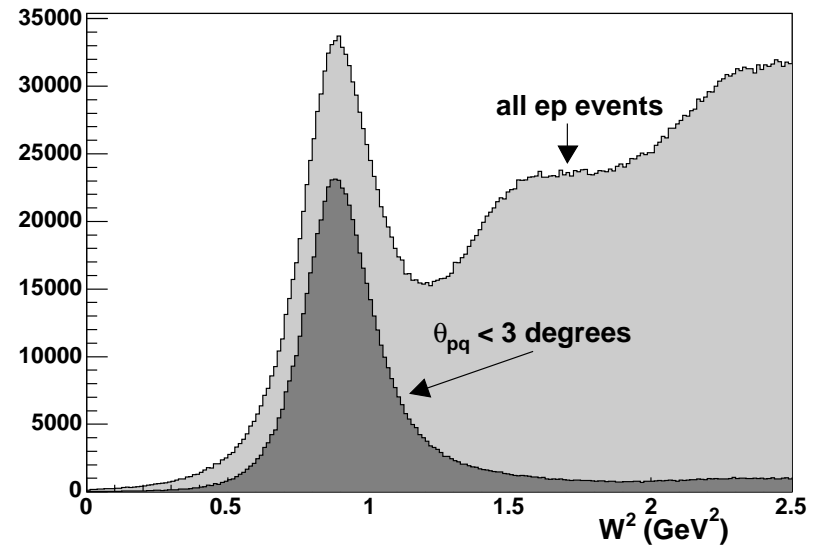


$$R = \frac{\frac{d\sigma}{d\Omega}(D(e, e'n))}{\frac{d\sigma}{d\Omega}(D(e, e'p))} = a(Q^2) \frac{\frac{G_E^n^2 + \tau G_M^n^2}{1 + \tau} + 2\tau G_M^n^2 \tan^2(\frac{\theta}{2})}{\frac{G_E^p^2 + \tau G_M^p^2}{1 + \tau} + 2\tau G_M^p^2 \tan^2(\frac{\theta}{2})}$$

- The ratio is less vulnerable to corrections like acceptance, efficiencies, *etc.*
- The dual target enables us to perform *in situ* detection calibrations.

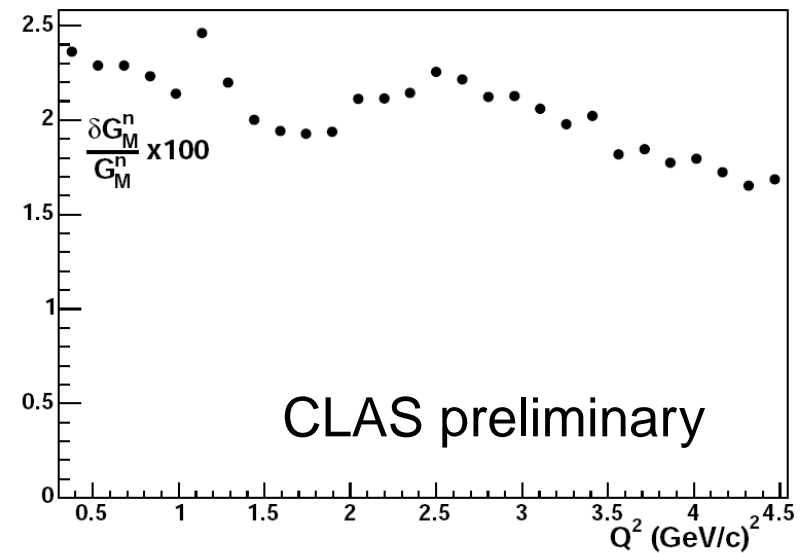
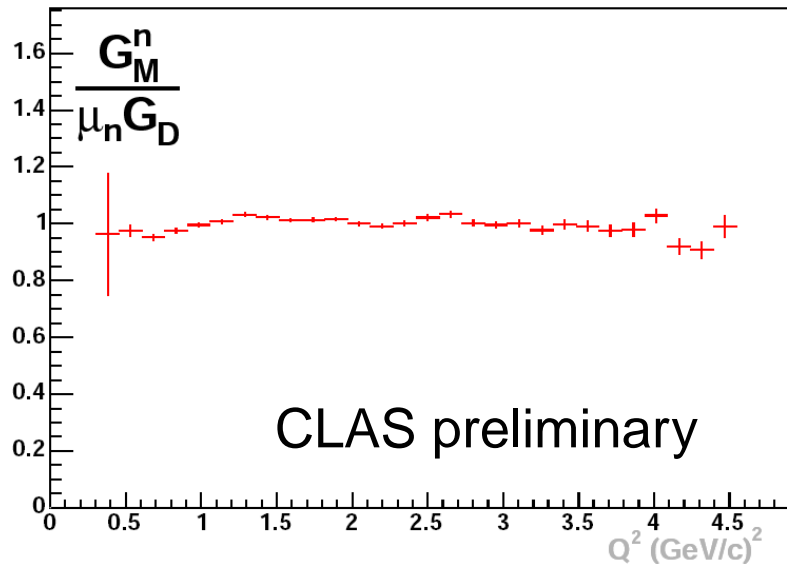
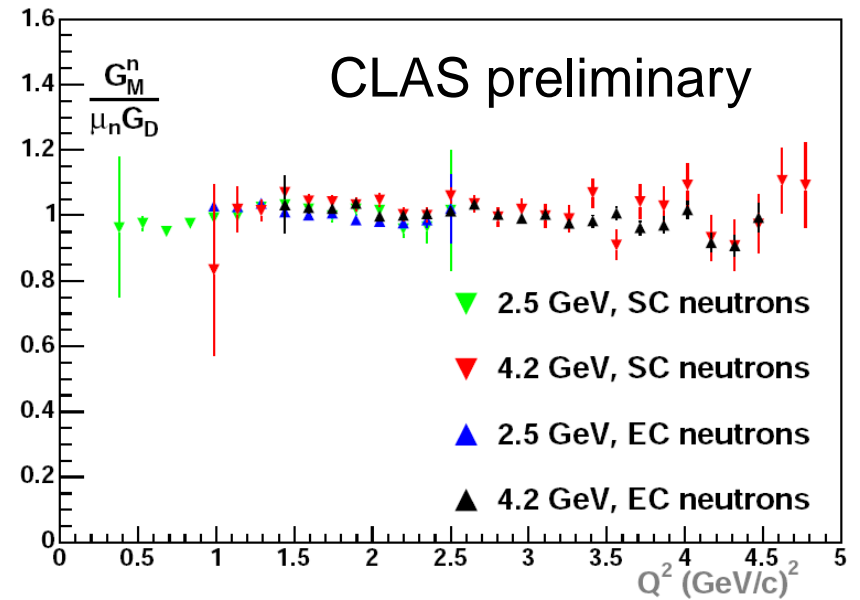
How Do We Measure G_M^n on a Neutron? (Step 4)

- Quasi-elastic event selection: Apply a maximum θ_{pq} cut to eliminate inelastic events plus a cut on W^2 (J.Lachniet thesis).
- Use the $ep \rightarrow e'\pi^+n$ reaction from the hydrogen target as a source of tagged neutrons in the TOF and calorimeter.

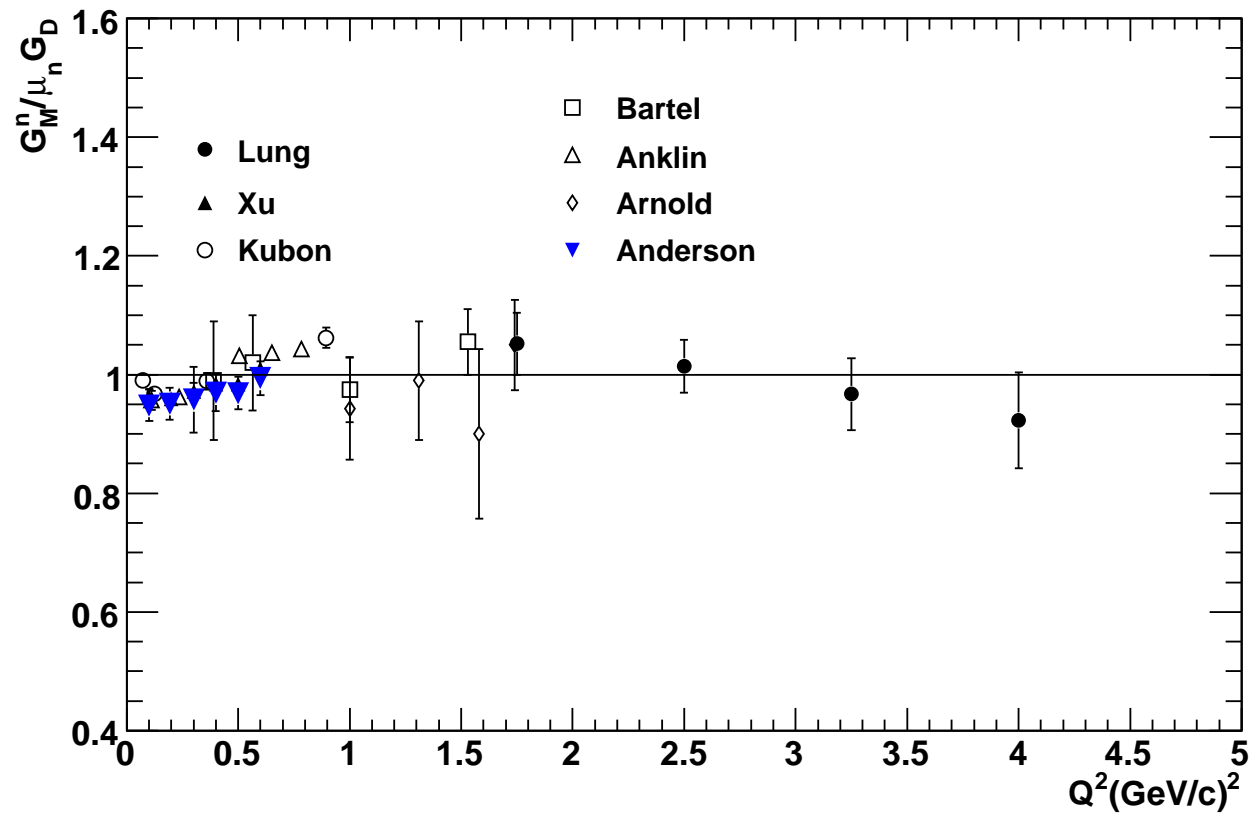


Results - Overlaps and Final Averages

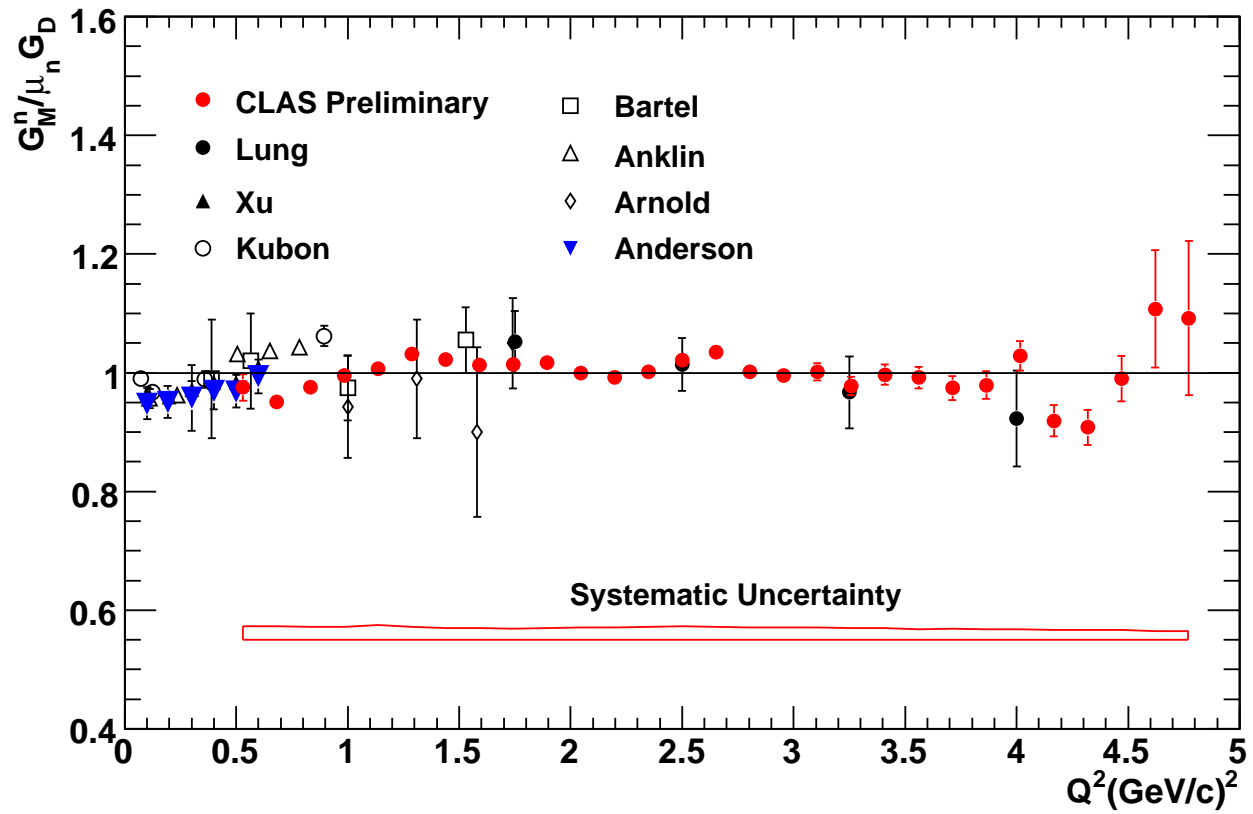
- Overlapping measurements of G_M^n scaled by the dipole are consistent.
- Weighted-average $G_M^n / \mu_n G_D$ and systematic uncertainty $\frac{\delta G_M^n}{G_M^n} \times 100 (< 2.5\%)$.



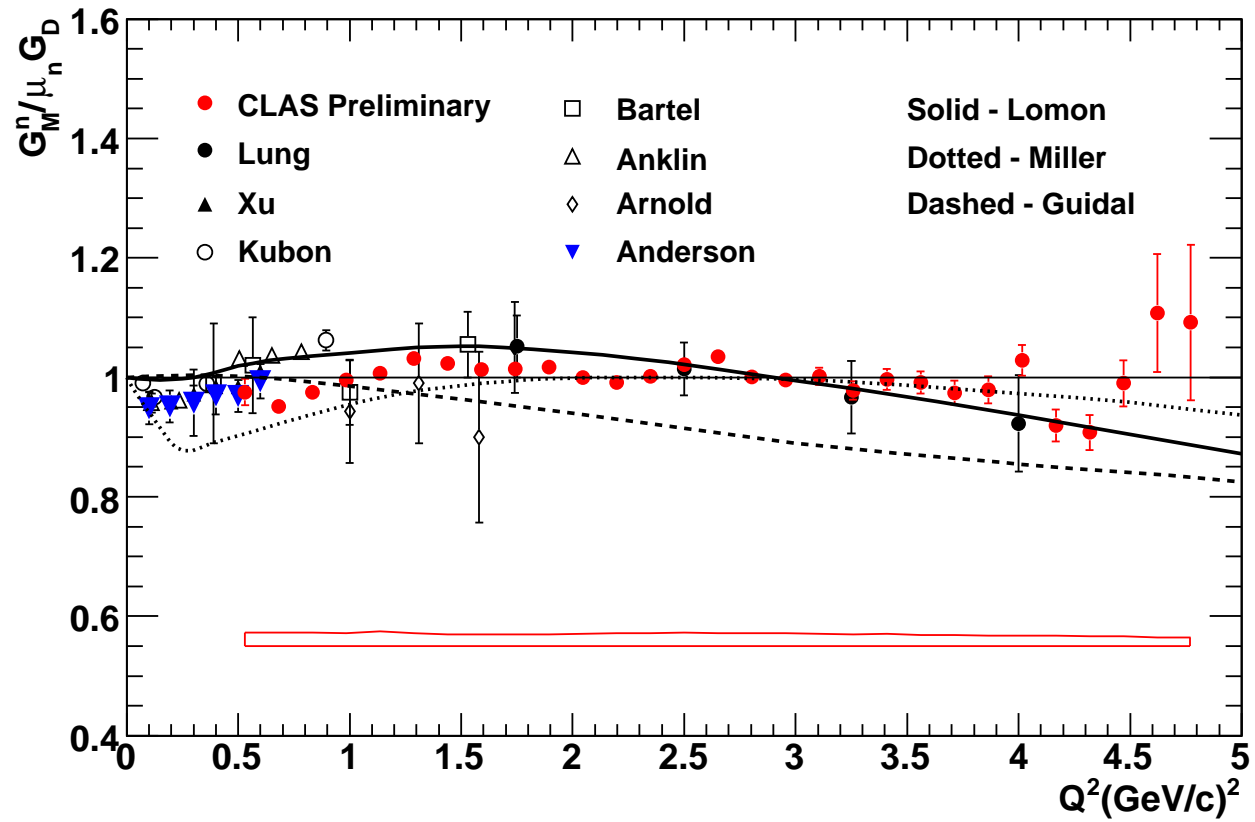
Results - Comparison with Existing Data and Theory



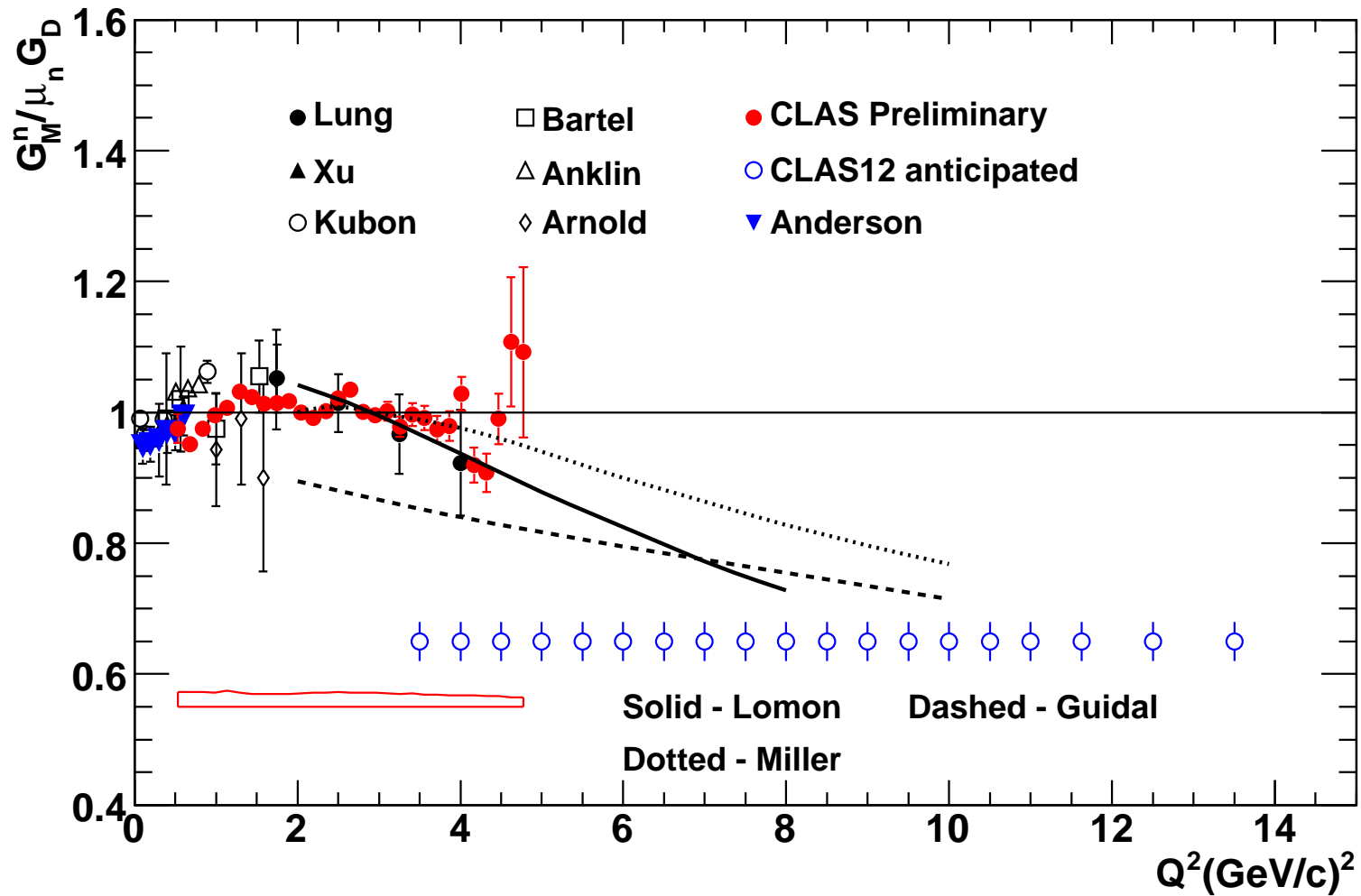
Results - Comparison with Existing Data and Theory



Results - Comparison with Existing Data and Theory

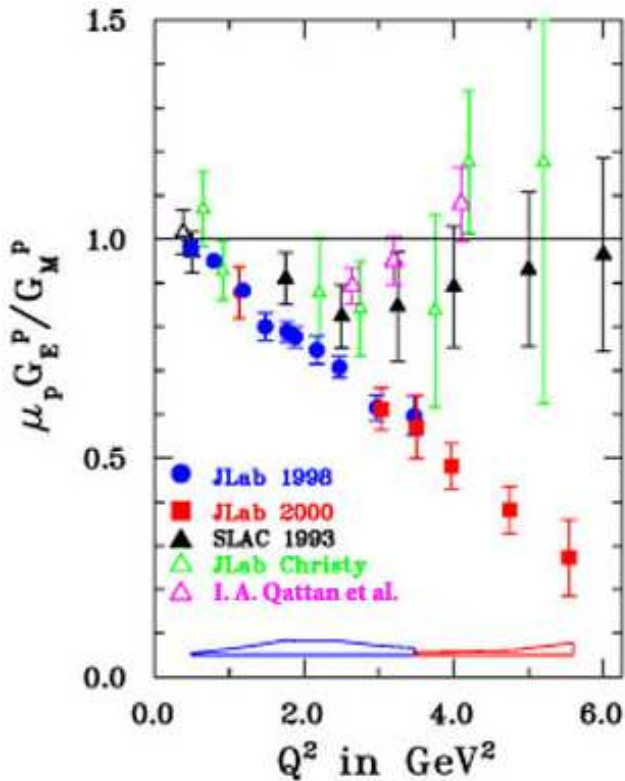


More To Come



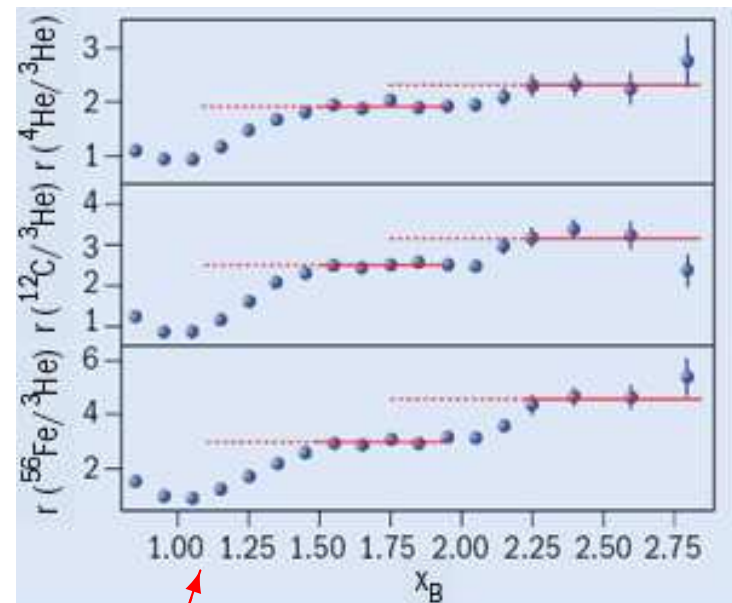
More Jefferson Lab Highlights

Ratio of charge and magnetization of the proton (G_E^p/G_M^p).



O. Gayou, *et al.*, Phys. Rev. Lett. **88**, 092301 (2002)

Short range correlations in nuclei.

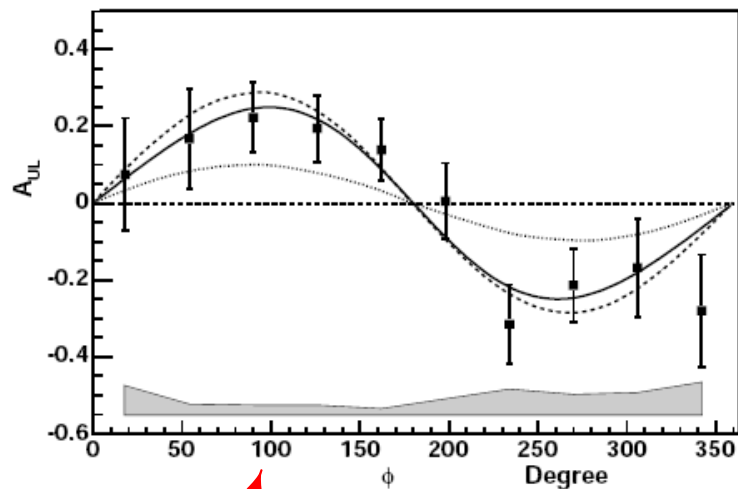


Fraction of nucleon momentum carried by struck quark.

K. S. Egiyan *et al.*, Phys. Rev. Lett. **96** (2006) 082501

Even More Jefferson Lab Highlights

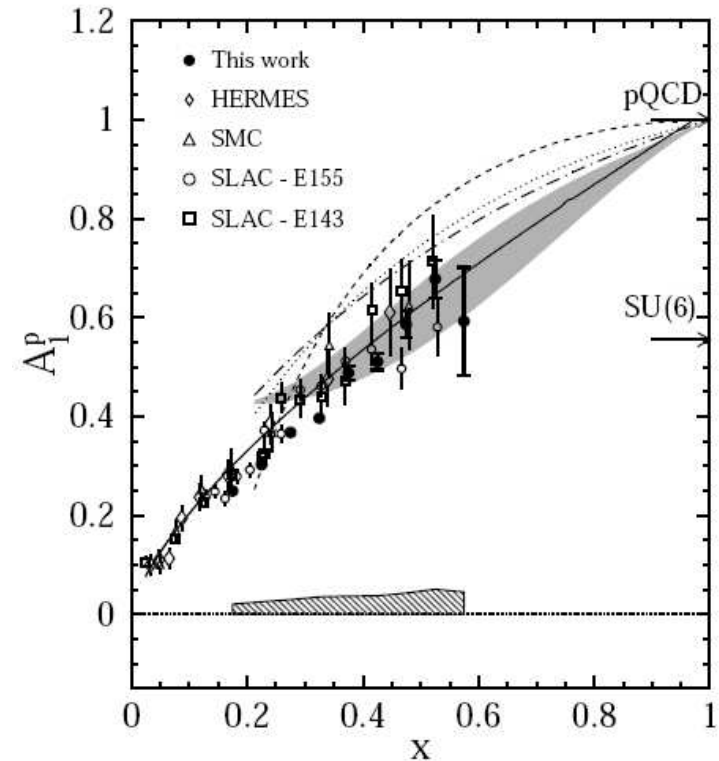
Towards a 3-dimensional picture of hadrons.



Azimuthal angle around the 3-momentum transfer \vec{q} .

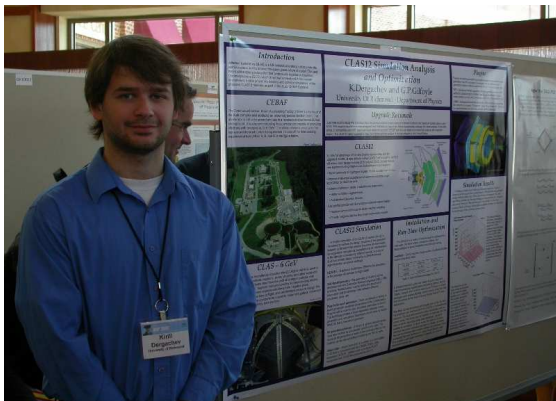
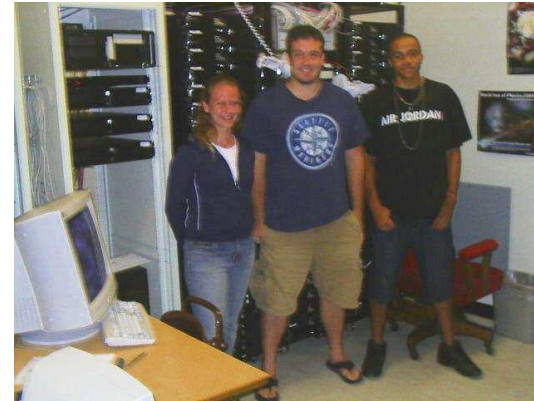
S. Stepanyan *et al.* (CLAS), Phys. Rev. Lett. **87** (2001) 182002

Where is the nucleon spin?.



K.V. Dharmawardane, *et al.* (CLAS), Phys. Lett. B **641**, 11 (2006)

Life on the Frontiers of Knowledge



Concluding Remarks

- JLab is a laboratory to test and expand our understanding of quark and nuclear matter, QCD, and the Standard Model.
- We continue the quest to unravel the nature of matter at greater and greater depths.
- Lots of new and exciting results are coming out.
- A bright future lies ahead with the JLab 12-GeV Upgrade.

