

Measuring G_M^n in CLAS12

- Use the ratio $R = \frac{e-n}{e-p}$ in quasi-elastic kinematics to extract G_M^n .

$$R = \frac{\frac{d\sigma}{d\Omega} [{}^2\text{H}(e,e'n)_{QE}]}{\frac{d\sigma}{d\Omega} [{}^2\text{H}(e,e'p)_{QE}]} = \frac{\sigma_{mott}^n \left(G_E^{n2} + \frac{\tau_p}{\epsilon_p} G_M^{n2} \right) \left(\frac{1}{1+\tau_n} \right)}{\sigma_{mott}^p \left(G_E^{p2} + \frac{\tau_p}{\epsilon_p} G_M^{p2} \right) \left(\frac{1}{1+\tau_p} \right)}$$

- Data Set:

- 1 Run Group B: Spring 2019, Fall 2019, Jan 2020
- 2 Run Group A: Fall 2018, Spring 2019, used to extract CLAS12 neutron detection efficiency (NDE) for the $e = n$ events in R .
- 3 Beam energies: 10.6 GeV and 10.2 GeV.
- 4 Torus polarity: inbending and out bending.

- Outline:

- 1 G_M^n : reaction, corrections
- 2 NDE: reaction, event selection, method, preliminary results
- 3 Acceptance Matching
- 4 Quasielastic Event Selection - method, very preliminary results

Extracting G_M^n

Use the ratio $R = \frac{e-n}{e-p}$ in quasi-elastic kinematics to extract G_M^n .

$$R = \frac{\sigma_{mott}^n \left(G_E^{n2} + \frac{\tau_n}{\epsilon_n} G_M^{n2} \right) \left(\frac{1}{1+\tau_n} \right)}{\sigma_{mott}^p \left(G_E^{p2} + \frac{\tau_p}{\epsilon_p} G_M^{p2} \right) \left(\frac{1}{1+\tau_p} \right)}$$

where and

$$\tau = \frac{Q^2}{4M^2} \quad \epsilon = \left(1 + 2(1 + \tau) \tan^2 \frac{\theta_e}{2} \right)^{-1} \quad \sigma_{Mott} = \frac{\alpha^2 E' \cos^2(\frac{\theta_e}{2})}{4E^3 \sin^4(\frac{\theta_e}{2})}$$

Solving for G_M^n

$$G_M^n = \sqrt{\left[R \left(\frac{\sigma_{mott}^p}{\sigma_{mott}^n} \right) \left(\frac{1 + \tau_n}{1 + \tau_p} \right) \left(G_E^{p2} + \frac{\tau_p}{\epsilon_p} G_M^{p2} \right) - G_E^{n2} \right] \frac{\epsilon_n}{\tau_n}}$$

Corrections to R

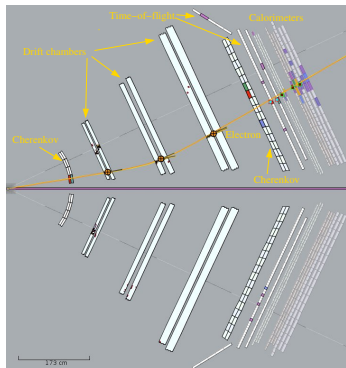
$$R(Q^2) = R_c(Q^2) = f_{NDE} f_{nuclear} f_{radiative} f_{fermi} \dots R_{obs}$$

NDE*	acceptance matching
radiative effects	nuclear effects
fermi motion	θ_{pq} range
momentum corrections	angle corrections

* NDE - neutron detection efficiency

Neutron Detection Efficiency (NDE) - 1

- Use the ${}^1\text{H}(e, e'\pi^+n)$ reaction as a source of tagged neutrons.
- Event Selection
 - 1 Use standard CLAS12 reconstruction code and select ${}^1\text{H}(e, e'\pi^+)X_n$ events where X_n can be any number of neutrals, *i.e.* include all neutrals as neutron candidates.
 - 2 Use standard Run Group A cuts*:
 - 1 calorimeter fiducial cuts
 - 2 sampling fraction in calorimeter
 - 3 correlations in calorimeter
 - 4 deposited energy in calorimeter
 - 5 HTCC photoelectrons cut
 - 6 vertex cut
 - 7 $e - \pi^+$ vertex difference cut
 - 8 χ^2PID cut.

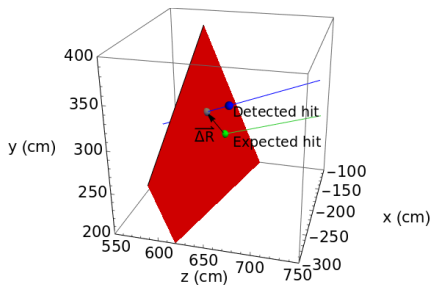
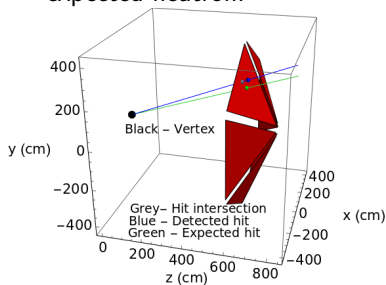


* CLAS12 RG-A - Analysis Note Overview and Procedures - Phase I, Towards SIDIS CLAS12

Neutron Detection Efficiency (NDE) - 2

- Tagging Neutrons:

- 1 Assume the reaction is ${}^1\text{H}(e, e'\pi^+)n$ and use the e' and π^+ information to predict the trajectory of the assumed neutron.
- 2 'Swim' the neutron through the CLAS12 detector to see if strikes the fiducial region of the detector.
- 3 If the 'swum' neutron DOES NOT strike CLAS12 drop the event.
- 4 If the 'swum' neutron DOES strike CLAS12, continue. This is an expected neutron.

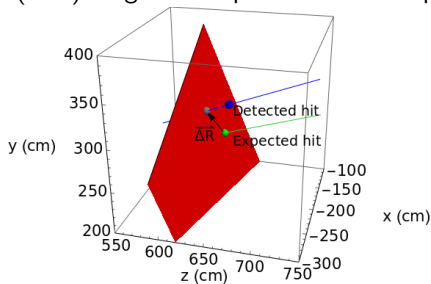
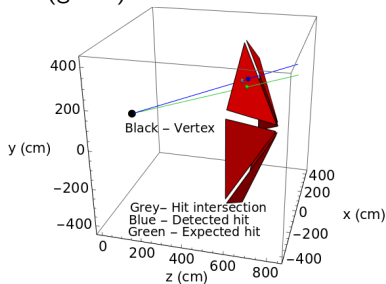


Red panels: ECAL front face, left - full size, right - close-up view

Neutron Detection Efficiency (NDE) - 3

- Tagging Neutrons (continued):

- ⑤ Now search the neutrals in the event to and see if one of those neutrals lies 'near' the predicted neutron track. See plot below. If a neutron is found this is a detected event.
- ⑥ Geometry of neutron trajectories in CLAS12 for expected neutron (green) and a detected neutron (blue). Right-hand panel is a close-up.

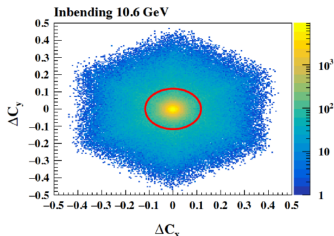


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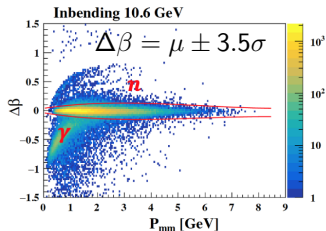
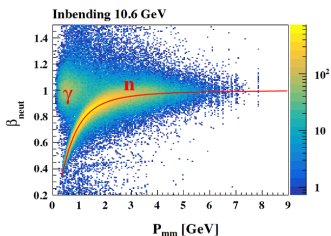
Neutron Detection Efficiency (NDE) - 4

- Tagging Neutrons (continued):

- ⑦ Cut on 'nearby' tracks, $\Delta C_x \Delta C_y$ cut.



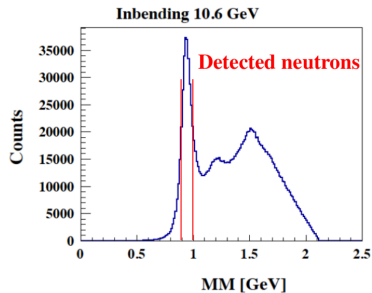
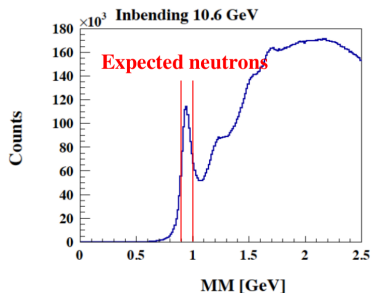
- ⑧ Cut on $\Delta\beta = \frac{P_n}{E_{miss}} - \beta_n$ where $\beta_n = \frac{\ell_n}{c\Delta t_n}$. These are detected neutrons.



Neutron Detection Efficiency (NDE) - 5

- Tagging Neutrons (continued):

- ⑨ Cut on neutron missing mass peak $0.9 \text{ GeV} < MM < 1.0 \text{ GeV}$.

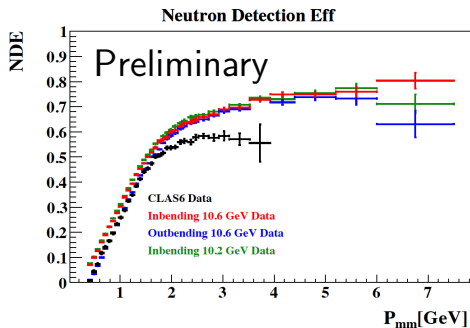


- ⑩ Now studying fits to the missing mass and background subtraction.

Neutron Detection Efficiency (NDE) - 6

- Extracting the NDE

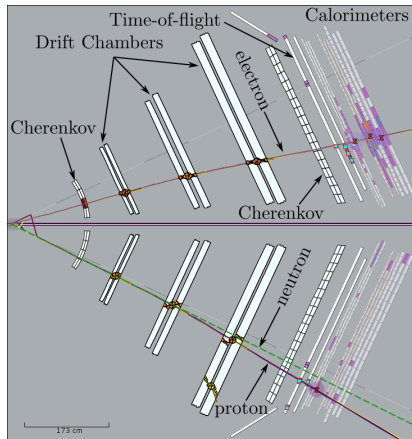
- Expected neutron events ${}^1\text{H}(e, e'\pi^+)X_n$ satisfy electron/pion cuts (from Run Group A) and expected neutron strikes calorimeter.
- Detected neutron events ${}^1\text{H}(e, e'\pi^+n)$ satisfy electron/pion cuts (from Run Group A) and $\Delta C_x \Delta C_y$ and $\Delta\beta$ cuts.
- The NDE ϵ is the ratio of detected to expected neutrons $\epsilon = \frac{N_{\text{detected}}}{N_{\text{expected}}}$.



Acceptance Matching

Since we divide the number of $e - n$ event by the number of $e - p$ ones we must correct for the different acceptances of neutrons and protons in CLAS12.

- 1 Electron passes the selection cuts.
- 2 Using only the electron information, assume elastic scattering, predict the proton momentum, and swim it through CLAS12.
- 3 If the 'swum' proton track strikes the CLAS12 fiducial volume, continue. If it does not, then drop the event.
- 4 Using only the electron information, assume elastic scattering, predict the neutron momentum, and swim the proton track through CLAS12.
- 5 If the 'swum' neutron track strikes the CLAS12 fiducial volume, continue. If it does not, then drop the event.
- 6 If both 'swum' tracks hit CLAS12, begin the nucleon analysis.



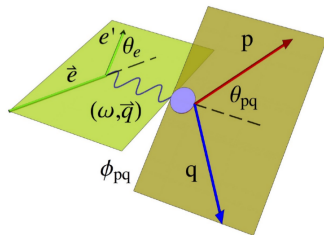
Quasielastic Event (QE) Selection - 1

- 1 The data can be used to calculate the incoming beam energy E_{beam} using different quantities. This feature give us the opportunity for cross-checks and corrections. See CLAS-NOTE-02-008.

$$E_{beam}^{angles} = M \left(\frac{1}{\tan\left(\frac{\theta_e}{2}\right) \tan\theta_N} - 1 \right) \quad E_{beam}^{mom} = p_e \cos\theta_e + p_N \cos\theta_N$$

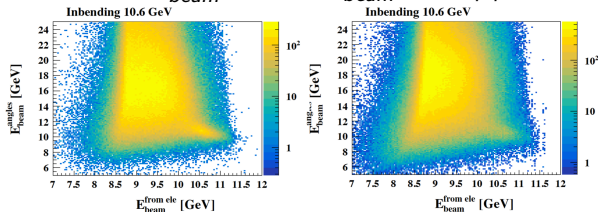
$$E_{beam}^e = \frac{E'}{1 + \frac{2E'}{m_N} \sin^2\left(\frac{\theta_e}{2}\right)}$$

- 2 The angle θ_{pq} is the angle between the 3-momentum transfer and the direction of the struck nucleon. For QE events θ_{pq} should be small while inelastic events are emitted at larger angles.



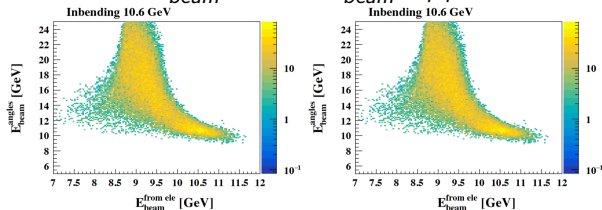
Quasielastic Event (QE) Selection - 2

- Extract the known beam energy from the measured electron and nucleon (see CLAS-NOTE-02-008).
- 2D plots of E_{beam}^{angles} versus E_{beam}^e . No θ_{pq} cut.



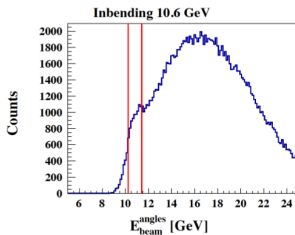
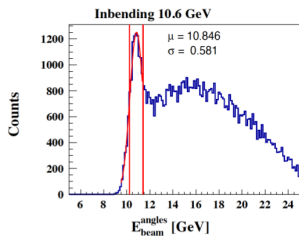
left column - $e-p$
right column - $e-n$

- 2D plots of E_{beam}^{angles} versus E_{beam}^e , $\theta_{pq} < 2.5^\circ$ cut.

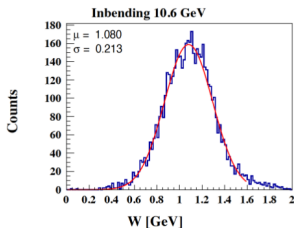
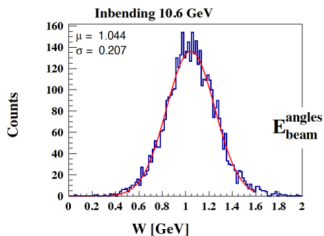


Quasielastic Event (QE) Selection - 3

- Still significant background.



- Shows up in W .

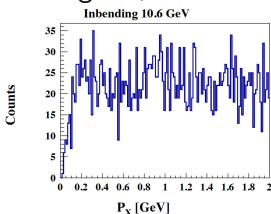
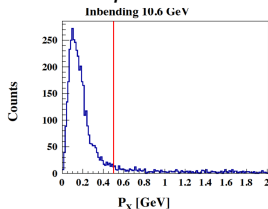


Quasielastic Event (QE) Selection - 4

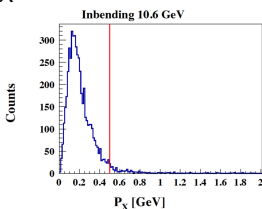
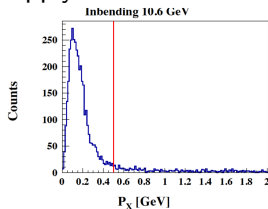
- An additional QE selection cut - missing momentum.

$$|\vec{P}_X| = P_{e'} + P_N - P_{beam}$$

- The $e - p$ distribution looks good, the $e - n$ one not so much.



- Apply corrections to P_X and the distribution improves.



Quasielastic Event (QE) Selection - 5

- Neutron momentum correction

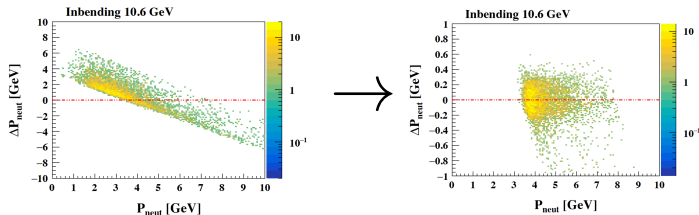
$$\Delta P_{neut} = p_{nfe} - p_n$$

where p_{nfe} is from the electron and p_n is from the neutron timing.

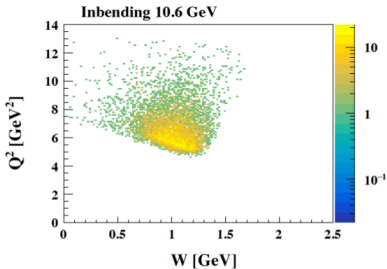
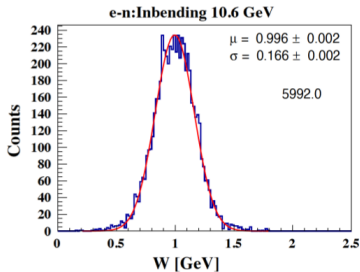
$$p_{nfe} = \sqrt{p_{beam}^2 - 2p_{beam}p_{e'} \cos \theta_e + p_{e'}^2}$$

$$p_{e'} = \frac{p_{beam}}{1 + 2p_{beam} \sin^2 \left(\frac{\theta_e}{2}\right) / m_N} \quad p_n = \frac{m_n \beta_{neutral}}{\sqrt{1 - \beta_{neutral}^2}} \quad \beta_{neutral} = \frac{\ell_n}{c\Delta t}$$

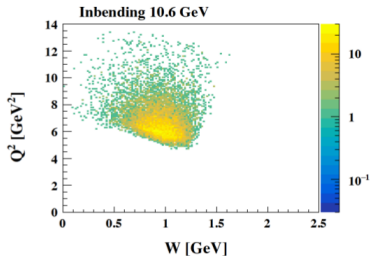
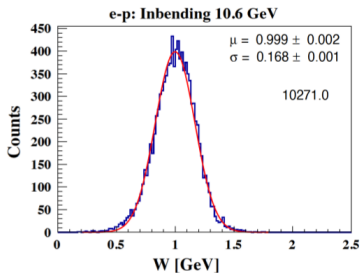
- Results of corrections to P_{neut} and θ_{neut} .



Current Status



e-n events



e-p events

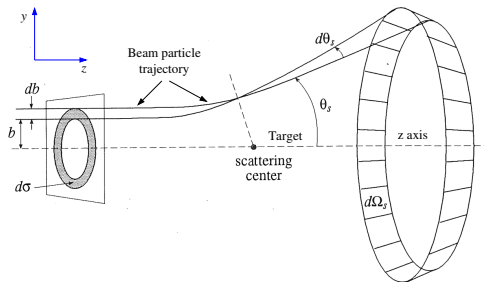
What is a Form Factor?

- Start with the cross section.

$$\frac{d\sigma}{d\Omega} = \frac{\text{scattered rate/solid angle}}{\text{incident rate/surface area}}$$

- For elastic scattering use the Rutherford cross section.

$$\frac{d\sigma}{d\Omega} = \frac{Z_{tgt}^2 Z_{beam}^2 \alpha^2 (\hbar c)^2}{16E^2 \sin^4(\theta/2)}$$



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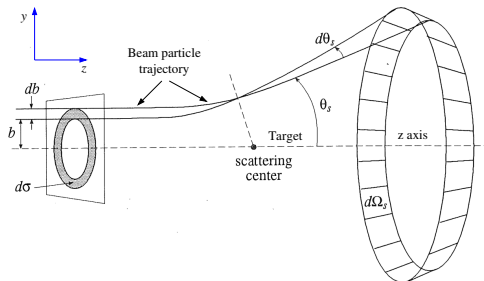
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- Cross section for elastic scattering by point particles with spin.

$$\frac{d\sigma}{d\Omega} = \frac{Z_{tgt}^2 Z_{beam}^2 \alpha^2 (\hbar c)^2}{16E^2 \sin^4(\theta/2)} \left(1 - \beta^2 \sin^2 \frac{\theta}{2}\right) \quad (\text{Mott cross section})$$



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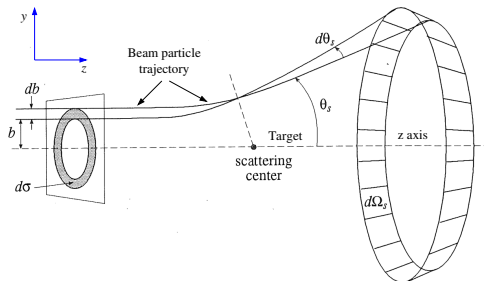
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- What happens when the beam is electrons and the target is not a point?

$$\frac{d\sigma}{d\Omega} = \frac{Z_{tgt}^2 \alpha^2 (\hbar c)^2}{16E^2 \sin^4(\theta/2)} \left(1 - \beta^2 \sin^2 \frac{\theta}{2}\right) |F(Q^2)|^2$$

where Q^2 is the 4-momentum transfer.



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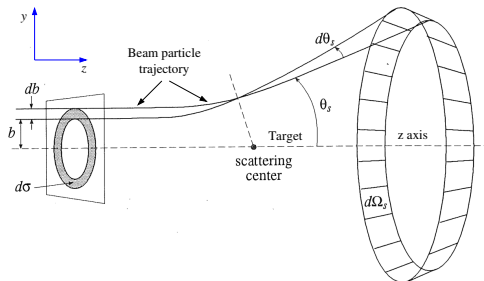
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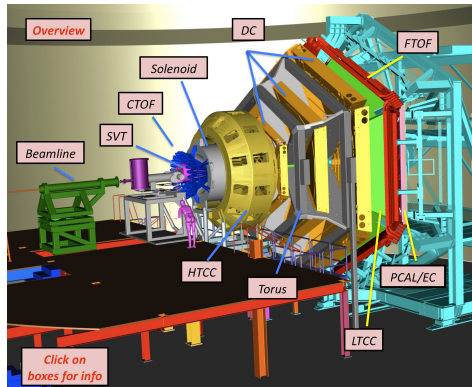
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THE FORM FACTOR!

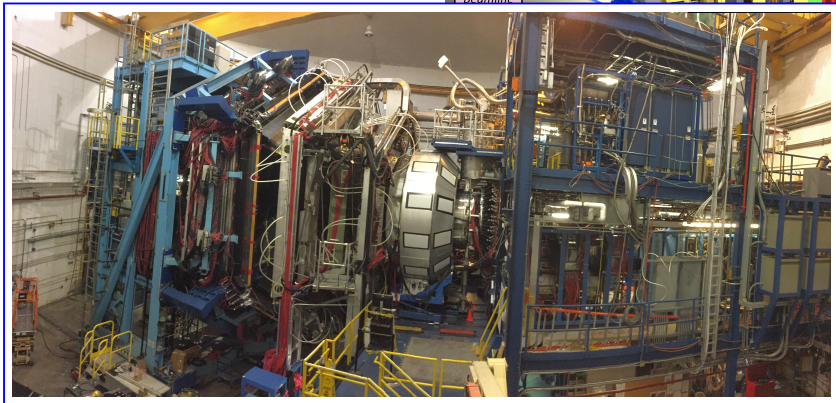
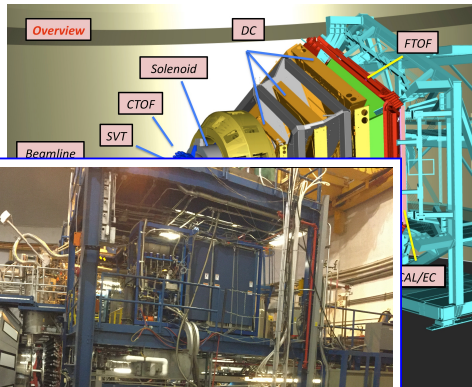
How Do We Measure G_M^n on a Neutron? (Step 2)

- Add one 45-ton, \$80-million radiation detector: the CEBAF Large Acceptance Spectrometer (CLAS12).
- CLAS12 covers a large fraction of the total solid angle at forward angles.
- Has about 62,000 detecting elements in about 40 layers.



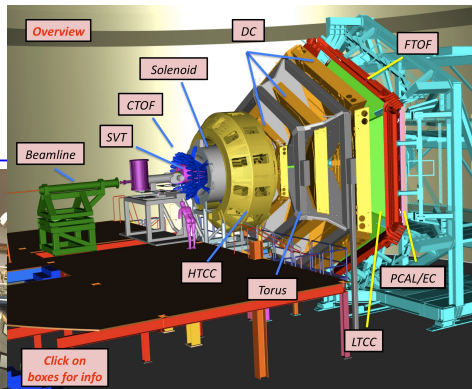
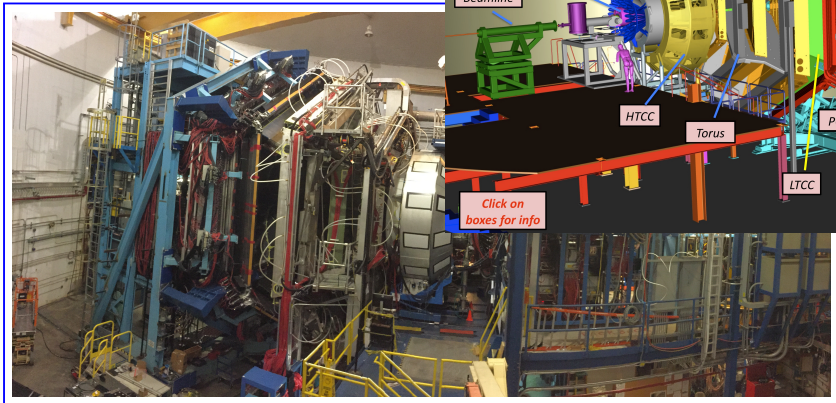
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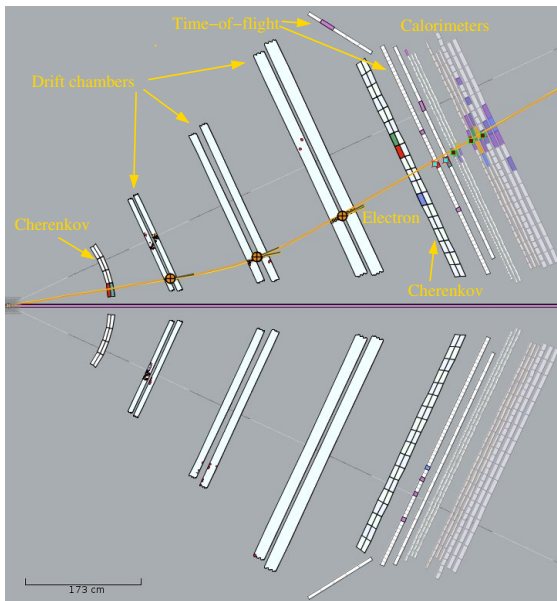


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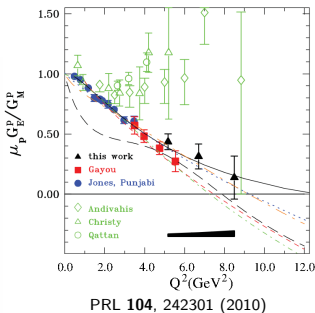
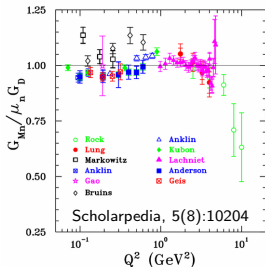
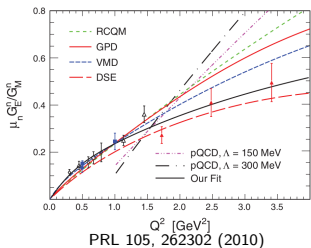
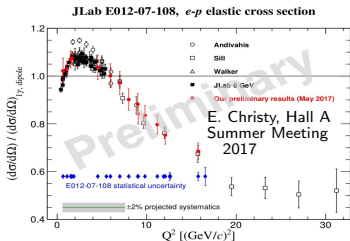


A CLAS12 Event



Where We Are Now.

- G_M^p well known over large Q^2 range.
- The ratio G_E^p/G_M^p from polarization transfer measurements diverged from previous Rosenbluth separations.
 - Two-photon exchange (TPE).
 - Effect of radiative corrections.
- Neutron magnetic FF G_M^n still follows dipole.
- High- Q^2 G_E^n opens up flavor decomposition.



How Do We Measure G_M^n on a Neutron? (Step 3)

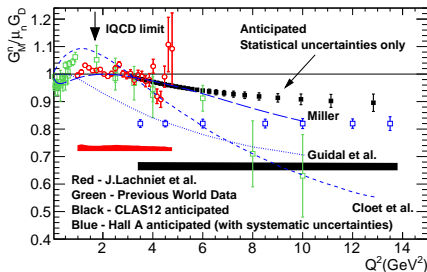
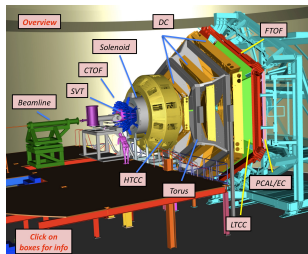
- E12-07-104 in Hall B (Gilfoyle, Hafidi, Brooks).
- Ratio Method on Deuterium:

$$R = \frac{\frac{d\sigma}{d\Omega} [{}^2\text{H}(e, e' n)_{QE}]}{\frac{d\sigma}{d\Omega} [{}^2\text{H}(e, e' p)_{QE}]}$$

$$= a \times \frac{\sigma_{\text{Mott}} \left(\frac{(G_E^n)^2 + \tau (G_M^n)^2}{1 + \tau} + 2\tau \tan^2 \frac{\theta_e}{2} (G_M^n)^2 \right)}{\frac{d\sigma}{d\Omega} [{}^1\text{H}(e, e' p)]}$$

where a is nuclear correction.

- Precise neutron detection efficiency needed to keep systematics low.
 - tagged neutrons from ${}^2\text{H}(e, e' pn)$.
 - LH_2 target.
- Kinematics: $Q^2 = 3.5 - 13.0 \text{ (GeV}/c^2\text{)}$.
- Beamtime: 40 days.
- Systematic uncertainties $< 2.5\%$ across full Q^2 range.
- Half of Run Group B done January, 2020.



How Do We Measure G_M^n on a Neutron? (Step 3)

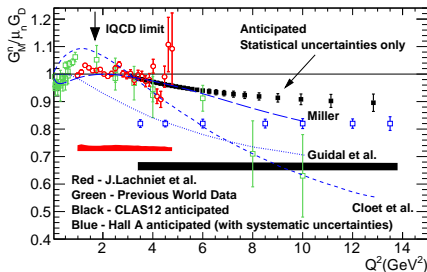
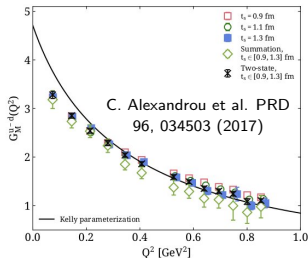
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- Ratio Method on Deuterium:

$$R = \frac{\frac{d\sigma}{d\Omega} [{}^2\text{H}(e, e' n)_{QE}]}{\frac{d\sigma}{d\Omega} [{}^2\text{H}(e, e' p)_{QE}]}$$

$$= a \times \frac{\sigma_{\text{Mott}} \left(\frac{(G_E^n)^2 + \tau (G_M^n)^2}{1 + \tau} + 2\tau \tan^2 \frac{\theta_e}{2} (G_M^n)^2 \right)}{\frac{d\sigma}{d\Omega} [{}^1\text{H}(e, e' p)]}$$

where a is nuclear correction.

- Precise neutron detection efficiency needed to keep systematics low.
 - tagged neutrons from ${}^2\text{H}(e, e' pn)$.
 - LH_2 target.
- Kinematics: $Q^2 = 3.5 - 13.0 \text{ (GeV}/c)^2$.
- Beamtime: 40 days.
- Systematic uncertainties $< 2.5\%$ across full Q^2 range.
- Half of Run Group B done January, 2020.



Anticipated Results

